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PHOTOGRAMMETRY OF THE PARTICLE TRAJECTORIES ON DIPOLE WEST SHOTS 8, 9, 10, AND 11 Volume I - Shot 10

University of Victoria British Columbia Canada V8W 2Y2

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Prepared for

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20. ABSTRACT (Continued)

For each charge configuration, one experiment was carried out over smooth ground and the other over rough ground. In each of the four experiments conducted, photogrammetrical measurements were made of the trajectories of air particle tracers (smoke puffs), which had been placed in a vertical grid at heights ranging from 3 feet to 58 feet above the ground and at radial distances ranging from 25 feet to 140 feet from the vertical axis through the charges. From the measured particle trajectories, calculations were made of the particle velocities, densities, hydrostatic overpressures, and dynamic pressures throughout the blast wave, at times ranging from 3 ms to 100 ms after detonation of the charges. Also determined from the photogrammetrical measurements were the shock front times-of-arrival. These were determined in each experiment for the primary shock front from each of the two charges; for the Mach stems produced above and below the interaction plane midway between the two charges, and for the Mach stem produced at the ground surface. From the shock front times-of-arrival, calculations were made of the shock velocities and, in turn, the peak particle velocities, air densities, and hydrostatic overpressures immediately behind each shock front. Calculations were also made of the variation with time of the particle velocity, density, hydrostatic overpressure and dynamic pressure at several fixed points. Results are presented both graphically and in tables, and are compared to results previously calculated for the same experiment using shock front photogrammetry and refractive image analysis.

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SUMMARY

Owing to the bulk of the material presented, this report is divided into several volumes. Volume 1 introduces the series and presents and discusses the results for Shot 10. Subsequent volumes will present and discuss the results for Shots 8, 9 and 11. The method of analysis is common to all four experiments and is described in detail only in Volume 1.

so that the results from the four experiments may be easily compared, they have been scaled to remove the effects of varying atmospheric conditions. (Results are scaled to a 1 kg charge weight and a standard atmosphere of dry air at 15 °C at sea level.) For the most part, only scaled results are presented. Exceptions include some derived pressure-time histories, which are meant to be compared to actual gauge measurements made in each of the experiments.

Results are presented in SI units, even though the experiments were originally laid out in British units. Only distance and time measurements are affected, however, as velocity, density, and pressure results are presented as dimensionless ratios. A distance units conversion scale is included for convenience, to convert between SI units (meters scaled to a 1 kg charge) and British units (feet scaled to a 1 lb charge), plus a time scale factor and scale factors to convert pressure ratios to both British and SI pressure units. Scale factors

which may be used to compute the distance and time values actually observed under the ambient conditions of each shot are also provided. Real pressure units are used for the results presented at gauge locations.

PREFACE

The authors gratefully acknowledge the opportunity offered by the Defence Research Establishment Suffield and the Defence Nuclear Agency to participate in the experiments described in this report. The analyses described here were carried out under contract with the Canadian General Electric Company, and with additional financial support from a research grant by the National Research Council (A 2952). The advice and assistance of Mr. A.P. Lambert, C.G.E. Project Officer at DRES, and Mr. J. Keefer, of the Ballistic Research Laboratory, is also gratefully acknowledged.

Unit conversion and scaling factors

FEET (SCALING TO 1 LB CHARGE)

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METERS (SCALING TO 1 KG CHARGE)

For feet scaled to a 1000 lb charge, multiply the top scale by 10.

For time scaled to a 1000 lb charge, multiply time scaled to a 1 kg charge by 7.683.

For pressure in kPa, multiply a pressure ratio (in atmospheres) by 101.325. For pressure in psi, multiply the pressure ratio by 14.696. To convert kPa to psi, divide

To obtain distance values actually observed for Shot 10, in meters, multiply scaled values in this report by 8.0718. To obtain the observed distance values in feet, multiply the reported scaled values by 26.482. To obtain observed time values, multiply scaled time values by 8.3742. For observed pressures in kPa, multiply by 94.38; for observed pressures in psi, multiply by 13.69.

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CHAPTER 1, INTRODUCTION

1.1 Blast wave reflection

When a spherical charge is detonated above the ground surface the resulting spherical blast wave reflects from the ground. At a distance from the point on the ground immediately beneath the charge, approximately equal to the height of the charge above the ground, the reflected shock begins to overtake and combine with the primary shock to form a single shock known as the Mach stem. The point at which the primary shock, the reflected shock and the Mach stem meet is called the triple point. As time progresses and the Mach stem shock moves outwards, the triple point moves outwards and upwards in a curved trajectory.

The physical properties of the Mach stem blast wave and the trajectory of the triple point depend primarily on three things: the energy yield of the charge, the height of the charge above the ground, and the nature of the ground surface. When the primary shock reflects from the ground some energy will be absorbed and appear as seismic disturbances, including cratering if the explosion is close enough to the ground. As the blast wave continues to move across the ground surface there will be a continued transfer of energy between the air and the ground,

and also a redistribution of energy within the blast wave. Little is known about the transfer of energy to the ground from an air blast, or about the redistribution of energy within a blast wave as it passes over the ground surface.

In the experiments described in this report, attempts have been made to simulate the "ideal" reflection of spherical blast waves to provide a reference for studies of reflection by real ground surfaces. It is postulated that if two identical charges are simultaneously detonated at a certain distance apart, the two resulting identical spherical blast waves will interact along a plane, and since there will be no energy loss in this interaction it will be possible to observe an ideal spherical blast reflection. The properties of the Mach stems lying above and below the ideal reflecting plane may then be compared with those produced over various real ground surfaces.

1.2 General description of the project

The primary purpose of the project was to obtain information on the interaction of spherical blast waves from explosive sources with real and ideal reflecting surfaces. The results may be used, for example, to evaluate hydrodynamic air blast computer codes. The blast wave interactions were obtained by the simultaneous detonation of two identical spherical charges

placed one above the other, such that the distance between the charges was equal to twice the height of the lower charge above the ground surface.

Two different charge heights were used over two different types of ground surface. The four experiments were titled Dipole West Shots 8, 9, 10, and 11. For Shots 8 and 11 the lower charge was placed at a height of 25 ft above the ground surface and the second charge at an additional height of 50 ft above the first. Shot 8 was carried out over smooth ground and Shot 11 over rough ground. For Shots 9 and 10 the corresponding figures for charge height and charge separation were 15 ft and 30 ft. The ground surface was smooth for Shot 9 and rough for Shot 10.

Two photogrammetrical studies were carried out in each of these experiments. The first study involved the high speed photography, against a suitable background, of the shock fronts produced by the two explosions. This permitted a calculation of the shock front trajectories and velocities, and thus a determination of the physical properties immediately behind the shocks as they moved into ambient air. The results of those calculations have been reported by Dewey et al (1975). The second project involved the high speed photography of an array of smoke puffs which acted as particle flow tracers, to determine the particle trajectories within the blast waves.

These trajectories have been analyzed to provide the space and time variations of particle velocity, density, hydrostatic overpressure, and dynamic pressure within the waves.

This volume describes the photogrammetry of the particle trajectories for Shot 10. The particle trajectory photogrammetry for the other three experiments will be reported in subsequent volumes. Shot 10 was analyzed first because the data collected were fewer than those collected for Shots 8 and 11 (fewer smoke puffs were used in Shot 10), and of better quality than those collected for Shot 9 (all smoke puffs detonated in Shot 10 and most remained visible throughout the time interval of interest). Thus the analysis of Shot 10 was carried out with a minimum of data processing and maximum expectation of useful results.

1.3 Description of Shot 10

Dipole West Shot 10 was fired on November 2nd, 1973 by the Ballistics Research Laboratories at the Defence Research Establishment Suffield, in Alberta, Canada. Two 1080-lb. spheres of Pentolite were detonated simultaneously, to within 5 microseconds, at nominal charge heights of 15 and 45 feet over a relatively rough (ploughed) ground surface.

Particle trajectory data were gathered by photographing the movement of smoke puffs formed in a vertical plane running

out from ground zero at 6.7° south of west. A WF5 camera operating at about 3400 frames per second was positioned 30 ft above the ground level at the main camera position 600 ft due south of ground zero.

Figure 1 shows a plan view of the field layout. The dashed line represents the approximate line of sight of the WF5 camera. Figure 2 shows the field of view of this camera.

The smoke puff grid was made up of 9 columns of 12 puffs each, hung vertically on strings. The vertical spacing of puffs was 5 ft, beginning 3 ft above ground level and ending at a height of 58 ft. The horizontal spacing of the columns of puffs was 10, 7 or 5 ft, depending on the distance from ground zero, beginning at about 25 ft and ending at about 85 ft. All 108 smoke puffs detonated successfully and a good film record was obtained.

This report describes the analysis of the smoke puff data collected for Shot 10, and presents and discusses some of the results of that analysis.

CHAPTER 2, METHOD OF ANALYSIS

2.1 Camera calibration

The exact orientation and lens position of the camera were determined using the measured positions in the film image of the photomarkers and other surveyed objects, such as pressure gauges. This calibration procedure has been described in detail by Dewey et al (1975). The positions of all fiducial markers, as they appeared in a chosen calibration frame, were transformed to a camera object plane in real space using the camera position, two camera orientation angles, and two "optical" parameters: a magnification factor and an axial rotation angle. The orientation angles and optical parameters were calculated so that the two reference point images in the calibration frame, Pl and P2, transformed to coincide with their positions in the object plane calculated directly from the field survey data. Since an exactly surveyed camera position was not provided, the three camera position coordinates were also calculated, together with the two orientation angles and the two optical parameters. The camera position "optimization" was accomplished by requiring that a third reference point image, P3, transformed to its object plane position as calculated from the survey data, and that

the distance between a pair of reference point images, P4 and P5, transformed to the corresponding distance in the object plane calculated from the survey data. The calibration procedure was iterative, and ended when the results of successive calculations differed by less than a pre-set tolerance.

The field survey data used in the calibration procedure for Shot 10 are listed in Table 1. The calculated camera position (\pm 0.1 feet) and orientation angles (\pm 0.001 °) are listed in Table 2, together with the positions of all the calibration points in the object plane and their "shifts", that is, the differences between the transformed calibration points and their object plane positions calculated directly from the survey data.

The object plane positions and shifts of the calibration points are also shown in Figure 3. Circles represent the transformed calibration points and squares represent positions calculated using the survey data. The larger circles indicate the three points that were made to coincide exactly, viz. P1 = W1, P2 = W3, and P3 = 300W1. The large half-circle indicates the final reference point (P5 = 300W2) used in the camera position optimization (P4 was set to equal P3). C1 and C2 denote the two charges. Frame centre is near the camera centering photomarker, VP2A. The photomarkers VP1A and VP1B show large shifts probably because they were not re-surveyed

after the pole on which they were mounted was lowered to replace the marker missing on Shot 8. No similar explanation could be found for the vertical shift of the foreground photomarker 300W2.

Consistency in calculating the optimum camera position from experiment to experiment was one of the criteria used in judging whether or not the optimum position was acceptable. The optimum camera positions used were the ones showing the least movement between experiments, on the assumption that the actual camera position was the same for each experiment. The calculated camera positions are shown below for the four experiments. Ground zero was given coordinates 2000 ft east, 2000 ft north, and 2316.32 ft elevation.

CALCULATED CAMERA POSITIONS

Shot	East	North	Elevation
8	2003.38	1385.09	2341.84 ft
9	2002.76	1385.15	2341.66 ft
10	2002.62	1388.29	2341.72 ft
11	2003.25	1386.89	2340.64 ft
	-		
Range	0.76	3.20	1.20 ft

The greater range in the north coordinates is probably due to the method of optimization used for this coordinate direction. The method was based on the apparent distance

between two points in the object plane rather than on the apparent position of one point in the object plane.

The position and orientation of a camera, once calculated, were assumed to remain constant throughout an entire film.

The magnification factor and axial rotation angle were re-calculated for each frame in the data reduction part of the analysis described below, to account for frame-to-frame variations in the film projection optics.

2.2 Data reduction

The positions of the smoke puffs and two reference photomarkers as they appeared in the projected film image were digitized for each frame over a section of film that corresponded approximately to the duration of the positive phase of the blast waves. For the Shot 10 film, this time interval corresponded to frames 10 through 200 (frame 0 corresponded approximately to detonation time). Before each digitizing session, the projection system optics were aligned using a precision grid in the film gate of the projector.

The two reference markers used for the reduction of Shot 10 were P6 = VP3B and P7 = 300W2. The digitized data were transformed mathematically to the camera object plane using the image positions of these two reference markers in each frame and their positions in the calibration frame, to correct for frame-to-frame variations in image magnification and axial

rotation in the film projection optics and to correct for frame-to-frame effects of camera vibration.

The camera object plane was defined to coincide as nearly as possible to the actual plane of smoke puffs. Smoke puff positions in the object plane were then transformed to the actual smoke puff plane, which was assumed to be a vertical plane running from ground zero at an angle of 6.7° south of west. The coordinate system in this plane was defined to have its x-axis horizontal, its y-axis vertical, and its origin at ground zero. Thus the entire grid of smoke puffs was positioned in real space for each frame.

A time (after charge detonation) was assigned to each frame by computing the film speed at selected frames. Film speeds were computed using the 1 ms timing marks placed on the film at the time of its exposure. The reciprocal of a simple polynomial function fitted to the computed film speeds was then integrated to determine the successive frame times. Zero time was determined using the detonation zero timing mark placed on the film. This method of film timing has been described in detail by Dewey et al (1975). A complete set of timing data for the Shot 10 film is provided in Table 3.

2.3 The output data plane

The coordinate origin in the smoke puff plane was nominally ground zero. However, the coordinate origin was not located at the ground zero position which was listed with

the field survey data. The location used was instead a value corrected to account for the fact that the charges were not positioned in any experiment exactly above the surveyed ground zero. The smoke puff plane origin was defined to be that point which had the same elevation as the surveyed ground zero, but which lay directly below a point halfway along the line joining the charge centers. This corrected ground zero for Shot 10 was displaced horizontally by 1.2 ft from the surveyed ground zero, in a direction which was approximately 30° south of west.

Particle trajectory data were output in the smoke puff plane using an x-coordinate axis which was defined to run horizontally outwards from the corrected ground zero and a y-coordinate axis which ran vertically upwards from this same point, so that increasing x-coordinate values denoted increasing distance from an idealized vertical charge axis and increasing y-coordinate values denoted increasing distance above an idealized horizontal ground surface. All data in the smoke puff plane as they are presented in the figures of this report have x-coordinate values increasing to the right, so that the smoke puffs appear to run to the right of the charges rather than to the left as they appeared in the original photographic images.

Figure 4 shows the positions of the 108 detonated smoke puffs comprising the grid for Shot 10 at a time prior to the

detonation of the two charges. These positions are in the plane of the smoke puff grid and the charges, as described above. The smoke puff plane is not exactly parallel to the camera image and object planes (Figures 2 and 3), and various geometrical corrections were applied to make the transformation between them. The puffs enclosed in parentheses were not visible in the earlier film frames, but were seen later when they were illuminated by the light of the fireball. Charge positions in the figures are plotted as if they were positioned exactly above the corrected ground zero origin. The data shown in Figure 4 have not been scaled.

2.4 Scaling and units

The position-time histories of individual smoke puffs were extracted from the frame-by-frame positions of the smoke puff grid. This set of particle trajectories was then scaled to standard atmospheric conditions and charge weight to remove experiment-to-experiment differences due to variations in atmospheric conditions and charge weight. A change from British units to SI units was made at this stage of the analysis.

Particle trajectory data were scaled by dividing all distances by Sachs scaling factor $S = \sqrt[3]{(WP_O)/(W_OP)}$ and multiplying all times by the factor $C/(C_OS)$, where C is

the ambient sound speed computed for Shot 10. Data used to compute C and S, and define the scaled event, are listed below with the computed values of C and S.

Ambient temperature,	T = -5.94 °C	(21.3 °F)
Ambient pressure,	P = 94.38 kPa	(13.689 PSI)
Relative humidity,	RH = 81.0 %	
Computed vapour pressure,	VP = 0.32 kPa	(2.4 mm Hg)
Ambient sound speed,	C = 328.003 m/s	(1076 ft/s)
Charge weight,	W = 489.9 kg	(1080 lbs)
Sachs scaling factor,	S = 8.0718	
Standard charge weight,	$W_O = 1.0 \text{ kg}$	(2.2 lbs)
Standard pressure,	$P_{O} = 101.325 \text{ kPa}$	(14.7 PSI)
Standard temperature,	$T_O = 15 ^{\circ}C$	(59 °F)
Standard sound speed, (dry air)	$C_{O} = 340.292 \text{ m/s}$	(1116 ft/s)

The results presented in this report therefore apply to a scaled event which is the detonation of two 1 kg charges in a standard atmosphere. The scaled heights of burst for Shot 10 were 0.563 meters and 1.713 meters.

2.5 Trajectory fitting

Figure 5 shows the scaled particle trajectory data for Shot 10 in the smoke puff plane with positions measured horizontally and vertically from corrected ground zero. Approximately 9600 puff positions are represented. As represented, the raw trajectory data have not been edited, and a number of obvious film reading errors can be seen.

The raw particle trajectory data were edited to remove obvious data processing errors, such as a single point widely displaced from its trajectory for one or two frames. The trajectories were then smoothed by fitting, for each puff in turn, simple polynomial expressions separately to both the x-and y-coordinate data, these being discrete functions of frame time. A least squares fitting technique was used. The adequacy of each fit was determined by examining on the same graphical output plots of both the raw trajectory data and the fitted curve. For Shot 10 this meant examining and adjusting 216 such plots, at least two or three times each.

For a given puff, the first step in fitting the raw trajectory data was to set the time of arrival of the shock front first hitting the puff. This was done subjectively, using the plotted raw data. The obvious constraint was used that this time for a single puff should be the same for the x- and the y-coordinates. Prior to this time of arrival, the measured positions of the puff were averaged to establish the initial puff position.

The raw trajectory data at times subsequent to the time of arrival were fitted with polynomial functions as described above. The fits were individually adjusted by varying the order of the polynomial and/or the relative weight given to each of the raw data points. The lowest order of polynomial which appeared to describe the particle trajectory was used. In all cases polynomials of fifth order or less were found to be adequate. Data points in obvious error were edited out by weighting them zero. For some puffs undergoing relatively more complex motions (i.e. those nearer to the charges), trajectories were fitted in two sections. If it appeared that the change in motion of the puff was the effect of a second shock wave, then it was not required that the first derivative of the fitting function (the velocity of the puff) be continuous. Otherwise it was required that this derivative at least appear to be continuous. In judging the motion of a particular puff, the motions of its near neighbours were also taken into account.

Using the fitted functions, the positions of the smoke puffs were interpolated at a series of discrete times, which were not necessarily the frame times. The first derivatives of the fitted functions were also calculated at the interpolated times for use in later calculations of particle velocity.

2.6 Regionalization and shock strength calculations

Knowing the time of arrival of a shock front at each puff, and the position of each puff, calculations were made of the shock trajectory, the shock velocity, and various peak parameters associated with shock front strength. These shock trajectories and associated peak parameters were then compared to the corresponding results calculated as part of the refractive image analysis already reported (Dewey et al, 1975).

To associate a shock radius with a given puff position, the smoke puff plane was divided into five regions. All puffs in a given region were assumed to have been first hit by the same shock front. The region boundaries were defined using triple point path measurements obtained from the refractive image analysis mentioned above.

The regions which were defined for Shot 10 are shown in Figure 6. In regions 1 and 2 the puffs were assumed to have been first struck by primary shock fronts, and the shock radii were measured from the appropriate charge center. Puffs in regions 3, 4 and 5 were assumed to have been first hit by a Mach stem: along the ground, and below and above the interaction plane, respectively. The interaction surface, along which the upper and lower primary shock fronts interact without energy loss, was assumed to be the plane of symmetry separating the charge centers.

In the two Mach stem regions adjacent to the interaction plane, radius values were measured from the point halfway between the charges, and in the Mach stem region adjacent to the ground surface they were measured from corrected ground zero. This was done in accordance with the assumptions that the Mach stem shock fronts were spherical and that the sphere centers remained fixed with time at the points where the vertical coordinate axis intersected the interaction and ground planes. Shock strength results computed in this way for Shot 10 represented an improvement over similar results computed using the assumption that the Mach stem fronts were cylindrically shaped and centered about the vertical coordinate axis. This question of Mach stem shape is discussed further in a later section of this report.

In each of the five regions, the shock trajectory data obtained from the first movement of the smoke puffs were fitted to a function of the form

$$r(t) = A + Bt + C \log(1 + t) ,$$

where r is the shock radius, t is the time after detonation, and A, B, and C are the fitted coefficients. The shock velocities were calculated by differentiating this function. The peak particle velocity, $V_{\rm S}$, peak density, $D_{\rm S}$, and peak hydrostatic overpressure, $P_{\rm S}$, in each of the five regions were

calculated as functions of shock radius, using extensions of the Rankine-Hugoniot equation, viz:

$$\frac{V_s}{C_a} = \frac{2}{\gamma + 1} \frac{M^2 - 1}{M^2}$$

$$\frac{D_{s}}{D_{a}} = \frac{(\gamma+1)M^{2}}{(\gamma-1)M^{2}+2}$$

$$\frac{P_s}{P_a} = \frac{2\gamma}{\gamma + 1} (M^2 - 1)$$

where C_a , D_a and P_a are the ambient values of sound speed, density and pressure, respectively; γ is the ratio of specific heats and M the Mach speed of the shock front. Peak particle velocity will be used in this report as a measure of shock strength because other independent calculations of this parameter were also made. Peak values of density and pressure were calculated using the above method only.

2.7 Particle velocity calculations

Assuming the smoke puffs to be perfect tracers of the gas flow within the blast waves, the gas particle velocities throughout the field of smoke puffs were calculated for a sequence of times. This was done by calculating the derivative of the polynomials used to interpolate the fitted trajectories. The derivatives of these polynomials were also calculated at

the times-of-arrival of the shock at each puff to provide a second measure of $V_{\rm S}$, the peak particle velocity immediately behind the shock.

The peak particle velocities V_s were also calculated by a third method. At specific times the particle velocities, V_s , and the radial positions, r_s , of all the puffs within each of the five regions were fitted to a function of the form

$$V(r) = A + B^2 r ,$$

that is, a straight line with positive slope. (The data scatter did not justify the fitting of a more flexible function.) This fitted function, at each time and in each region, was extrapolated to the appropriate shock radius at that time to determine peak particle velocity.

2.8 Density and hydrostatic overpressure calculations

The particle trajectory data were also used to compute time varying air densities throughout the field of the smoke puff grid. The calculated densities were in turn used to compute hydrostatic overpressures in the same field.

The smoke puff grid of 9 columns of 12 puffs each was treated as a grid of 88 cells, each cell being defined by 4 adjacent smoke puffs. Each such quadrilateral cell was

viewed as representing the cross-section of an elemental volume defined by rotating the area of the cell about the axis of cylindrical symmetry rising vertically through the two charges. Conservation of mass within this volume demands that the product RAD is constant at all times, where R is the radial coordinate of the centroid of the cell defining the volume, A is the area of the cell, and D the average density of air in the volume. In particular, the product RAD must always equal $(RAD)_a$, its ambient value computed before the arrival of any shock fronts at the cell. This equality gives the relative density D/D_a at any time, as

$$\frac{D}{D_a} = \frac{(AR)_a}{A R}$$

Pressure values were calculated using the fact that P/D^{γ} remains constant as long as the entropy of the gas in a cell remains constant, where P is absolute pressure at a point, D density, and γ the ratio of specific heats of air. In particular, the value of P/D^{γ} equals the value $P_S/(D_S)^{\gamma}$ measured immediately behind the leading shock front, and remains constant until the entropy is changed by the arrival of a second shock. The relative overpressure at a point can be computed using

$$\frac{P}{P_{a}} - 1 = \frac{P_{s} (D/D_{a})^{\gamma}}{P_{a} (D_{s}/D_{a})^{\gamma}} - 1$$

The values P_s/P_a and D_s/D_a were computed at each puff's position immediately after it was first struck by a shock front, taking P_s and D_s from the shock velocity data computed according to the method described earlier (section 2.6). The values computed for P_s/P_a and D_s/D_a at the cell corners were averaged. The overpressure P/P_a - 1, subsequently calculated, therefore also represents an average value over a given cell. The arrival of shock fronts subsequent to the first could not be detected accurately, so that no attempt was made to determine exactly those times after which the calculated overpressures in given cells became invalid.

2.9 Summary

Previous sections describe the method of particle trajectory analysis by which data of four basic types were obtained: times of shock front arrival, particle velocities, densities and hydrostatic overpressures. To obtain a measure of shock strength the particle trajectory data were analyzed in separate groups according to the shock front which first arrived at the smoke puffs in each group. It was assumed that the shock fronts had a specific regular shape. In this way the smoke puffs at the observed times of shock front arrival were assigned radial position coordinates to supplement their measured Cartesian coordinates. Radial shock velocities

throughout the smoke puff grid were computed and peak particle velocities were derived from the shock velocities. These peak velocity values were used as a measure of shock front strength and were compared with values obtained using smoke puff data in two other ways, and with values obtained previously using refractive image analysis. The particle trajectory analysis depended on the refractive image analysis only to define the five regions in the smoke puff plane. These regions were used only in the analysis of the shock fronts and were of no particular significance in the subsequent analysis of the smoke puff trajectories behind the shocks.

The basic data obtained from the smoke puff trajectories after shock front arrival were derived in Lagrangian coordinates, that is, for a grid of points which moved with the air in the blast wave and did not remain fixed with respect to the ground surface. The method used to transform these data to Eulerian coordinates, that is, fixed in space, and which also permitted the calculation of dynamic pressure fields, is the subject of the next chapter.

CHAPTER 3. SURFACE REPRESENTATION

3.1 Introduction

Measurement of the smoke puff positions permitted the determination of several physical properties of the blast waves in a vertical x-y plane passing through the corrected ground zero. Letting z denote values of a particular physical property at a specific time, the results of the particle trajectory analysis may be viewed as sets of discrete points on three-dimensional surfaces, f(x,y,z) = 0. If such surface representations are made explicit, they may be used to interpolate z-coordinate values, i.e. values of the physical property, at x-y positions other than the actual smoke puff positions, such as gauge locations. Such surface representations also may be used to display graphically an entire set of physical properties in the form of contour lines joining points of equal z-value. The following sections describe how the various surfaces were fitted.

3.2 Preliminary interpolations

Source data (z-coordinates) for surface fitting were the times-of-arrival of the first shock front at the individual smoke puffs, the measured particle velocities at puff positions at a sequence of times (velocity fields), and average densities

at cell centroids at the same sequence of times (density fields). Values of hydrostatic overpressure at the cell centroids were also fitted, these values being derived from the density data, as described previously. The x-y coordinates were the initial puff positions, the puff positions at a time, t , and the cell centroid positions at a time, t .

The first step in the surface fitting procedure was the completion of the data grid. This meant generating x-y coordinates to represent "missing" puffs and cell centroids, and interpolating z-coordinates at these new positions. This was done in a very conservative manner, using linear interpolation in a domain bounded by the completely convex set of straight line sections which connected the outermost of the originally existing puffs. No extrapolation beyond this domain was attempted.

The positions of missing puffs were assigned by interpolating x- and y-coordinates in turn, between those of the adjacent puffs in each missing puff's row and column. The z-data assigned to the new positions were then interpolated along the row and along the column, and the two results were averaged. New z-coordinates were assigned only to new puff locations that fell inside the original domain boundaries. A new puff position on the domain boundary, that is a puff with only one nearest neighbour in its row or column, was given a position coordinate and a z-coordinate equal to

those of its single neighbour in the appropriate row or column calculation.

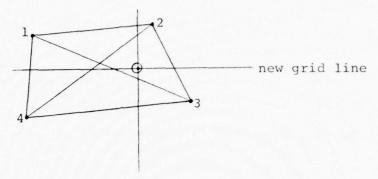
In Shot 10 there were no missing puffs in the original grid and it was not necessary therefore to apply the above technique when fitting the time of arrival surface. In all other cases, missing puff data were generated only to expedite subsequent interpolations.

3.3 Interpolation onto an Eulerian grid

The second step in the surface fitting procedure was the definition of a rectangular grid and the subsequent interpolation of a z-coordinate value at each new grid point inside the original domain. In all cases the new grid was a regular square grid with a constant mesh size of 0.1 m, which was approximately one-half the mesh size of the original smoke puff grid after scaling. The new grid was in the same plane as the smoke puffs but, unlike the smoke puff grid, it remained stationary relative to the coordinate origin.

Interpolation at the new grid points was done by fitting a polyhedral surface to the original, completed smoke puff grid. This was done by fitting plane surfaces to triplets of data at the corners of the original smoke puff cells. For example, when a new grid point fell inside an

original cell of smoke puffs (1,2,3, and 4) as shown:



a z-coordinate value was assigned to the new grid position "\phi" by fitting planes to the puff data at 1, 2, and 3, and at 2, 3, and 4, and interpolating at the new grid point using each plane in turn, and averaging the two results. If, at the edge of the domain, a cell was defined by only three puffs, only one interpolation plane was used. When two or more cell corners were beyond the domain, no data were interpolated to the new grid. In this way the original domain boundaries were shrunk, and made concave, on occasion.

The surface represented by the resulting grid of interpolated data was therefore an array of connected triangular plane elements, with four such elements for each complete smoke puff cell. The connected array was everywhere continuous in z, albeit discontinuous in the derivatives of z across the connections, and it passed through all of the original smoke puff data. There was therefore no smoothing of the

original data, and no chance of excessive values of z between original data due to uncontrolled fitting functions.

3.4 Dynamic pressure calculations

Surfaces were fitted to the times of shock arrival at each smoke puff and to the particle velocities, densities and hydrostatic overpressures at a sequence of times. The grids were Eulerian in that their position in the x-y plane did not change with time. The z-coordinates were, specifically, the velocity vector magnitudes expressed as Mach numbers, density ratios relative to ambient density, and overpressure ratios relative to ambient pressure.

Since the same interpolated grids were used for velocity and density, the dynamic pressure at each grid point also could be computed using

$$\text{dynamic pressure ratio} = \frac{\frac{1}{2}D\left|V\right|^{2}}{P_{O}} = \frac{\gamma}{2}\left(\frac{D}{D_{O}}\right)\left(\frac{\left|V\right|}{C_{O}}\right)^{2}$$

where $\gamma=1.4$, (D/D_O) is the measured relative density, $(|V|/C_O)$ the measured particle velocity Mach number, P_O the standard pressure, and C_O the standard sound speed. Dynamic pressures could not be calculated prior to this point in the analysis because the original, uninterpolated values of relative density and particle velocity were not measured at the same set of points in the x-y plane.

3.5 Contours and time histories

The fitted surfaces described above were used only for interpolation; the data were not smoothed. Interpolation onto an Eulerian grid made possible the calculation of dynamic pressure fields. Further interpolation to points inside the Eulerian grid made it possible to compute and plot contour lines over the various data fields, and also time history curves for data at specific fixed locations within those fields.

The contours calculated were those of equal time of shock arrival (isochrones), particle velocity (isotachs), density (isopyknics), hydrostatic overpressure (static isobars), and dynamic pressure (dynamic isobars). The isochrones represent the shock front shapes at different times. The contouring technique therefore provided a complete mapping of all the physical properties of the blast wave within the smoke puff domain without reference to other measurements. The time histories of these physical properties at several fixed locations were also calculated. Some of these were calculated at pressure gauge positions in unscaled units, so they may be compared to actual gauge measurements.

Contours of specific z-values were calculated using linear interpolation, searching between pairs of grid points along the Eulerian-Cartesian grid lines and along the mesh diagonals for possible contour line intersection points, and

joining the points which were found. Time history curves were calculated by interpolating z-values at a fixed location in a time-series of data fields which had been interpolated onto the same Eulerian grid. The interpolations to points inside this common grid were of the polyhedral surface type describe previously. Because all fitted surfaces were continuous in z and because the mesh size could be made relatively small, satisfactory contours and time histories were obtained. Because the fitted surfaces were smooth (comparitively speaking) and, in the limit of an infinitely small mesh, they would pass through all the control points in the measured data fields, the contours and time histories were smooth and represented the measured data faithfully in all cases.

CHAPTER 4. SHOT 10 RESULTS

4.1 Times of shock front arrival

The measured initial puff positions, the times of first shock arrival, and the peak particle velocities obtained by differentiating the functions fitted to the particle trajectories are presented in Table 4. Puff position is given relative to corrected ground zero as origin with horizontal and vertical axes. Puff position and the time of arrival of the first shock are given both as observed and scaled. Particle velocities listed are derivatives of the fitted puff trajectories at the times of shock arrival, and are expressed in Mach units. Expressed this way, the particle velocities are the same scaled as unscaled. Also listed are the initial radial puff positions (scaled) and region codes.

Shock front data determined from the first movement of the smoke puffs, i.e. calculated from the time-of-arrival data in Table 4, are listed in Tables 5.1 - 5.5. Each table corresponds to one of the 5 regions used. Listed are the observed and fitted unscaled shock trajectory data, the scaled fitted shock trajectory data, and the computed shock velocities and peak parameters associated with shock front

strength: peak hydrostatic overpressures in atmospheres and in kilopascals, peak particle velocities in Mach units, and peak density ratios. Given as ratios, these peak parameters are the same scaled as unscaled. Pressure given in kilopascals in the tables refers to the unscaled (observed) case only.

The shock front radius versus time data derived using particle trajectory analysis (PTA) are also shown in Figures 7.1 - 7.3 for the two primary fronts, the two Mach stems at the interaction plane, and the ground Mach stem, respectively. They are compared to corresponding data derived from refractive image analysis (RIA) reported by Dewey et al (1975). The refractive image analysis results were obtained using photogrammetry against a striped canvas backdrop and they describe the shock as it travelled in a direction almost diametrically opposite to the direction of the smoke puff grid.

4.2 Shock strengths

Peak particle velocities calculated from shock front velocities are shown in Figures 8.1 - 8.3 for the primary fronts, interaction Mach stems, and the ground Mach stem. This method of determining peak particle velocities has been labelled method 1, and the data plotted correspond to

those listed in Tables 5.1 - 5.5. The results in the figures are compared with those previously obtained using refractive image analysis (RIA). In the case of the primary shock fronts, results are also compared to those of Brode (1957) for TNT.

Peak particle velocities calculated as the derivatives of the fitted particle trajectory functions at the time of shock arrival are shown in Figures 9.1 - 9.3, labelled method 2. Corresponding numerical data are listed in Table 4. Results are compared to Brode in the primary front case, and with the results of method 1 in the Mach stem cases.

Peak particle velocities calculated by fitting particle velocity as a function of radial position in a region at a fixed time, and extrapolating to the shock radius at that fixed time (method 3), are shown in Figures 10.1 - 10.3.

The corresponding numerical data are listed in Tables 6.1 - 6.5. Peak particle velocities are compared to Brode in the primary front case, and to method 1 results in the Mach stem cases.

Some of the data from which the method 3 results were extrapolated are shown in Figure 11.1. Shown are the magnitudes of the particle velocity vectors throughout the blast wave in the regions of the Mach stem regions above and below the interaction plane. Results are plotted versus a radial coordinate computed in the same manner as for the time-of-arrival results. Also shown are the straight line fits to the data, extrapolated to the shock front radius at each of

the times shown. Figure 11.2 shows similar results for the primary shock front regions, compared to the results of Brode for TNT.

4.3 Particle velocity fields

The calculated particle velocities in the plane of the smoke puffs are shown as vectors in Figures 12.1 through 12.23 for times between 0.4 and 1.0 ms in 0.1 ms increments, and for times between 1.0 and 6.5 ms in 0.5 ms increments. The particle velocity field is shown at several other intermediate times to illustrate certain columns of smoke puffs being struck by the advancing shock fronts. All times and positions are scaled to a 1 kg charge in a standard atmosphere. Position coordinate axes extend horizontally and vertically from the corrected ground zero. The charge positions are also shown. The particle velocity vectors represent the derivatives of the smoothed particle trajectories, and their magnitudes may be judged using the standard vector shown on each figure. All velocities are measured in Mach units, relative to the standard sound speed. Puffs not yet struck by a shock wave are represented by small circles (zero velocity).

Numerical data corresponding to Figures 12.1 - 12.23 are listed in Tables 7.1 through 7.16, along with scaled radial positions of the puffs, and region codes as defined in Figure 6.

Conversion factors are given at the foot of each table, which may be used to convert the scaled data in the tables and figures back to their original unscaled values.

4.4 Density and hydrostatic overpressure fields

Calculated average relative densities throughout the smoke puff plane are depicted graphically in Figures 13.1 - 13.23, for times between 0.4 and 1.0 ms in 0.1 ms increments, and between 1.0 and 6.5 ms in 0.5 ms increments. The average density field is shown at several other intermediate times to illustrate the transition of certain cell columns through their periods of peak density. All time values are scaled. Cell positions are scaled and are given relative to the corrected ground zero as origin with horizontal and vertical axes. Charge positions are also shown. The calculated densities may be judged using the density shading scale shown on each figure. Density is given as a ratio, relative to ambient density. Cells not yet struck by a shock wave and cells in which the density has dropped to a value less than ambient density are shown blank.

Corresponding numerical data are listed in Tables 8.1 - 8.8 along with radial cell positions computed according to the regions defined previously. Numerical data describing the fields of hydrostatic overpressure are similarly listed in Tables 9.1 - 9.8. The pressure results for a given cell were obtained by multiplying the density results for that cell by

a factor determined by the strength of the shock which first traversed the cell, but which did not vary with time.

4.5 Times-of-arrival surface

Figure 14 shows a perspective view of the surface fitted to the original smoke puff positions and the observed times of first shock front arrival, i.e. to the data listed in Table 4. The grid mesh size is 0.1 by 0.1 meters (scaled), about 2.5 feet square (unscaled), or about 1/2 that of the original smoke puff grid. The charge positions are indicated on the vertical distance axis.

The times-of-arrival surface is smooth enough to permit contouring and the contours in this case (isochrones) represent shock front shapes at different times. These contours are plotted in Figure 15. The times-of-arrival surface was not smooth enough to permit the calculation of gradient vectors which could be used to compute shock velocity vectors and shock strengths over the new grid.

Two attempts were made to obtain contours of shock strength. In the first, the times-of-arrival surface was smoothed by least-squares fitting low-order, one-dimensional polynomial functions to the time-of-arrival data along each grid row and column separately, and computing the derivatives of the fitted functions to obtain the associated components of the surface

gradient vectors. Shock velocity vectors were obtained from the time-of-arrival gradients, and from these peak particle velocities were computed. The peak particle velocity (shock strength) surface is shown in Figure 16. The contours of this surface (not shown) did not exhibit any discontinuities across the boundaries of the shock front regions, as they would if surfaces were fitted to the times of arrival in each region separately.

The results of the second method used to compute shock strength contours are shown in Figure 17. These were obtained by interpolating an average shock radius at each value of peak particle velocity shown, for each shock front region in turn, using a composite of the various peak particle velocity versus radius curves shown in Figures 8 through 10. Arcs of circles with these radii, centered on the appropriate points along the vertical charge axis, were then drawn in the regions to represent shock strength contours. These peak value contours are discontinuous across triple point locii and other region boundaries. As.a result, some horizontal lines are crossed twice by the same contour or, in other words, identical shock strengths can be found at two locations the same vertical distance from a reflecting surface, but at different radial distances from the vertical charge axis. Examples are illustrated in the figure.

4.6 Field surface contours

Contours of equal particle velocity, density, hydrostatic overpressure, and dynamic pressure in the blast waves were determined for a series of times, using surfaces fitted to the various measured data fields at those times. Sample results are shown in Figures 18 through 21 at scaled times of 2.5 ms and 4.0 ms. The shock fronts shown in these figures are obtained from the time-of-arrival surface (as were those in Figure 15). Field contours such as those shown can be drawn for any scaled time between 0.5 ms and 6.5 ms. It should be re-stated that all of these results were obtained from the photography of the smoke puffs only and do not rely on the results obtained using the refractive image analysis (Dewey et al, 1975).

4.7 Time histories

By mapping the physical properties of the blast waves at short time intervals it was possible to determine the time histories of these properties at any selected fixed position within the smoke puff grid. This was done at 12 fixed locations, three in the two primary regions and three in each of the three Mach stem regions, as shown in Figure 22. At each distance from the axis of the charges in the Mach stem regions, each of the time history stations is the same distance from either the interaction plane or the ground plane. Particle velocity time histories could be interpolated closest to the ground level

because these were measured at puff locations, whereas the density and pressure data were measured at cell centers.

Time histories of particle velocity, density, hydrostatic and dynamic overpressure at these locations are given in Figures 23 to 26.

Also plotted with the time histories are the interpolated values of the time of arrival of the first shock front at the stations. The height of this time-of-arrival line represents a peak value derived from the shock velocity analysis.

Time histories for hydrostatic and dynamic pressure are also plotted in Figure 27.1 to 27.7 for stations at the nominal positions of field-mounted pressure gauges on the "60 foot gun barrel". The gauges on this gun barrel were mounted at nominal elevations of 10, 15, 20, 27, 30, 33, and 40 feet. The time histories at these locations are given in unscaled units in order to facilitate comparisons with the gauge measurements.

The dynamic pressures plotted in figures 26 and 27 are maximum values, computed using both the x and y components of particle velocity. Similar plots were made of the horizontal components of dynamic pressure, but the differences were not significant since the y components of particle velocity at these locations were small. Other locations could have been chosen at which the y components would not have been insignificant.

CHAPTER 5. DISCUSSION

5.1 Choice of Dipole West Shot 10

The results presented in this report were obtained from an original attempt to analyze an array of discrete two-dimensional particle trajectories within a blast wave.

(During the preparation of this report Suffield Special Publication No. 71, by John Anderson was published, which also describes a blast wave two-dimensional particle trajectory analysis using somewhat different techniques to the ones presented here.) Previous analyses of blast wave particle trajectories were limited to those cases in which it could be assumed that the motion of the particles was rectilinear (Dewey, 1964; 1971).

Particle trajectory photogrammetrical measurements were made on Dipole West Shots 8, 9, 10 and 11, and it was decided to develop the two-dimensional analytical techniques using the measurements from Shot 10 because in this experiment a smaller number of smoke puffs had been used, and these puffs had formed with a high degree of success so as to produce an almost always complete rectangular array of particle tracers.

5.2 Particle trajectory analysis technique

The simplest possible techniques were used to analyze the particle trajectories obtained for Shot 10. The trajectory

of each puff was described by least squares fitting a low order polynomial to the time variation of both the x and y coordinates of the puff. The two polynomials obtained in this way could be differentiated at any point to give the particle velocity components, and interpolated to give the particle position at any time. The gas density within the blast wave was determined by using the positions of four adjacent puffs to define a quadrilateral cell. The resulting density is thus an estimate of the average density within the cell, the position of which is described by its centroid.

This analysis system considers each smoke puff in isolation, and although some smoothing of the data is achieved by the least squares fitting, no account is taken of the positions of adjacent puffs. Subsequently linear interpolation is used between the smoke puff positions so as to obtain a complete mapping of the blast wave properties throughout the domain of the smoke puff array.

5.3 Reliability and accuracy of the measurements

An attempt has been made in this report to evaluate the reliability of the particle velocity and density measurements obtained from the analysis of the time resolved smoke puff trajectories. This is not easy to do since there are no similar measurements with which the results can be compared. However, the particle trajectory results have been compared

with the peak shock front values obtained from the refractive image shock velocity measurements (Dewey et al, 1975). Also, hydrostatic and dynamic pressure results obtained from the particle trajectory analysis are presented for positions at which pressure gauges were mounted so that a direct comparison of the pressure-time histories will be possible.

comparisons with the peak values at the shock front are not completely valid because the particle trajectory analysis method does not give results at the shock itself but only in the blast flow behind the shock. This necessitates extrapolating the particle trajectory data to the expected position and time of arrival of the shock. Extrapolation is rarely reliable and in the technique described as method 3, extrapolation was applied to data gathered from a large region of the blast wave and which thus showed significant scatter.

In order to calculate shock velocities assumptions were made about the shapes of the shock fronts. In the case of the primary shocks the assumption of sphericity appears to be valid. The shape of the Mach stem shock is still an open question. In the refractive image analysis (Dewey et al, 1975) the Mach stem shock positions were measured as close as possible to the reflecting surfaces and the assumption of cylindrical symmetry appears to have been valid. However, in the analysis of the smoke puff trajectories, measurements

could not be made so very close to these reflecting surfaces. It was clearly seen that the Mach stems were not cylindrical and as a first approximation it was assumed that these shocks were spherical with the centre at ground zero, or at the equivalent point on the ideal reflecting surface. This assumption gave results which were internally consistent, but it is clear that an investigation of the shape of the Mach stem shocks produced by a spherical shock reflecting from a plane surface, is a matter of urgency.

Three methods of calculating shock strength are described. The first uses the time of arrival of the shock front at the smoke puffs within a given region and analyzes these data in the same way as for the shock front refractive image photogrammetry described by Dewey et al, 1975. These results (PTA) are compared with the refractive image (RIA) results in Figures 8.1 to 8.3. There seem to be some significant differences between the results of the two measurement techniques: in particular the lowering of the Mach stem shock strength over the rough ground compared with the ideal reflecting surface is not so apparent when measured using the particle trajectory technique. This is probably due to the fact that the refractive image measurements were made within 0.3 m of the rough ground surface whereas the particle trajectory time-of-arrival measurements were made throughout the Mach stem region up to a distance of 4 m from the ground. Also,

the two sets of results describe the shock fronts travelling in opposite directions from the charges.

The second method of calculating peak particle velocities (shock strengths) used the derivatives of the two polynomials fitted to each particle trajectory, evaluated at the estimated time of shock arrival. Differentiation of a fitted polynomial at the extreme end of the observed data range is of questionable validity, and as might be expected the results plotted in Figures 9.1 to 9.3 show a considerable amount of scatter. This scatter is, however, uniformly distributed about the expected curve.

The peak particle velocities were also calculated by fitting all particle velocities in a region as a function of radial position in the region at a fixed time, and extrapolating to the shock radius at that time. Because of the scatter in the data obtained throughout the region, fitting the results to a curve of higher than first order did not seem justified. The results of these extrapolations are shown in Figures 10.1 to 10.3.

Although the extrapolation of the particle trajectory data beyond its domain of validity has produced a large degree of scatter, the peak particle velocity results are generally consistent with the results calculated using the method of refractive image analysis.

An attempt has been made to estimate the inherent accuracy and reliability of the particle trajectory measurements. A smoke puff typically has a diameter of about 0.3 m, and it is estimated that the centre of such a puff is identified with an accuracy of approximately ± 0.1 m. This leads to an estimated maximum error in particle velocity of ± 0.008 m/ms ($\pm \text{Mach } 0.035$), and an average density ratio of ± 0.35 . The mean peak particle velocity and density ratio measured in this experiment were approximately Mach 1 and 2.5 respectively, and so it is estimated that the accuracy of the reported results is approximately 5% for particle velocity and 15% for density ratio.

Dynamic pressure ratios were computed from the velocity and density results and are estimated to be accurate to within about 20%. Hydrostatic pressure ratios were computed using the density results and the peak overpressures computed from the shock velocities. Because the shock velocity data depended on region definitions and, in the case of Mach stems, on assumed shock front shapes, the accuracy of the hydrostatic overpressures computed behind the shock fronts is the most difficult to judge. Hydrostatic pressure ratios computed using cylindrical and spherical Mach stem shapes differed by as much as 10%. Combining this maximum figure with the estimated uncertainty given above for densities, it is estimated that the inaccuracy of the hydrostatic overpressures reported may be as high as 25%.

5.4 Surface mapping and contouring

The initial analysis of the smoke puff positions measured from the films provided a measure of the particle velocity of each smoke puff and the average density in each cell of four adjacent puffs. Chapter 3 described the use of linear and planar interpolation between the discrete values of particle velocity and density so as to provide a complete mapping of the parameters in the x-y plane, at any specified time. From such mappings, various contours were drawn to connect points having the same parameter value. In addition, from sequences of such mappings it was possible to determine the time histories of the parameters at several fixed positions within the smoke puff array.

A considerable amount of information is obtained from the analysis of the particle trajectories, and it seems clear that contour mappings and time histories best display this information. The contours can be compared easily with similar results from computer code predictions of the blast properties, and the time histories of hydrostatic and dynamic pressure can be readily compared with gauge measurements. To facilitate future comparisons the contours can be mapped at any specified time during the passage of the blast wave, and the time histories can be calculated for any specified position within the smoke puff array. Such outputs will be provided on request.

5.5 Dynamic pressure corrections

The dynamic pressure presented in this report is defined as one half the local gas density times the square of the maximum velocity component. In comparing these results with total head gauge measurements, compressibility corrections should be applied in the higher velocity regions.

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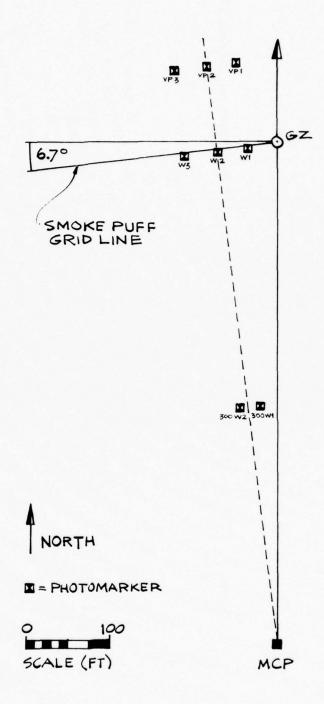


Fig. 1 Plan view of test site, Dipole West/10

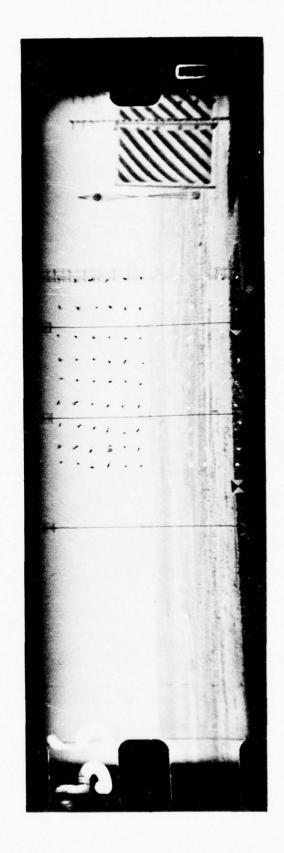
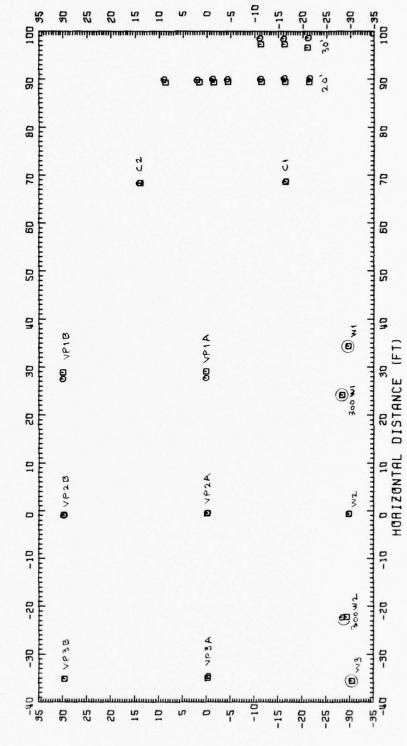


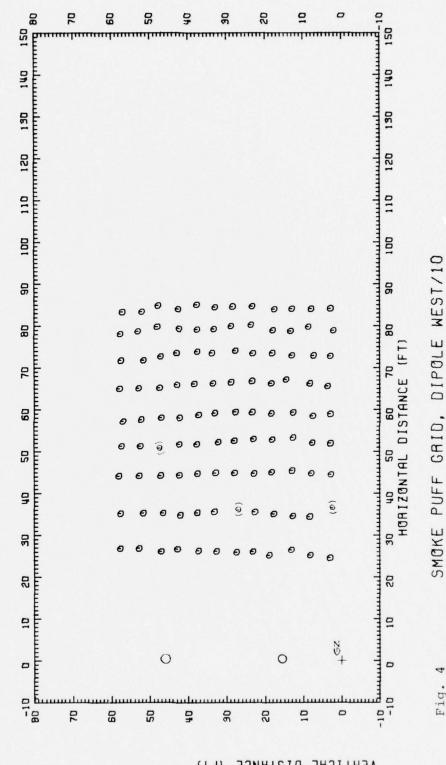
Fig. 2 Field of view of camera, Dipole West/10

D = PHOTOMARKER POSITION IN OBJECT PLANE CALCULATED FROM SURVEY DATA 0 = PHOTOMARKER POSITION IN OBJECT PLANE TRANSFORMED FROM FILM IMAGE

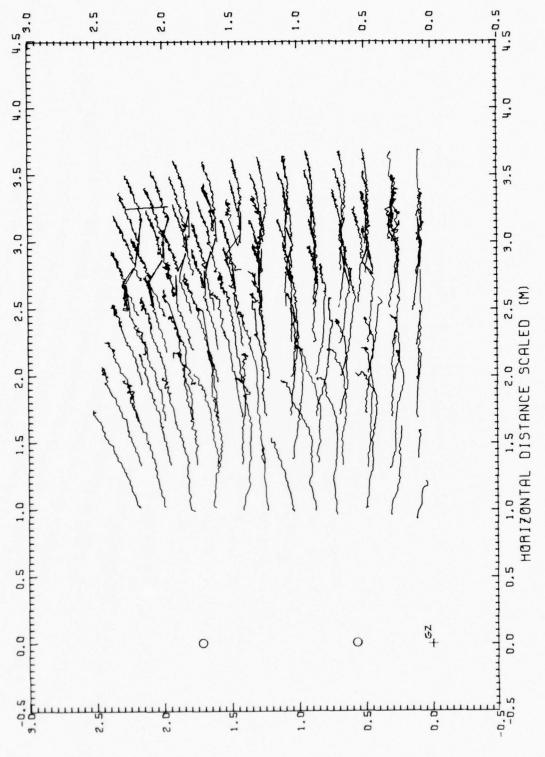


CAMERA CALIBRATION, DIPOLE WEST/10

Fig. 3



VERTICAL DISTRNCE (FT)

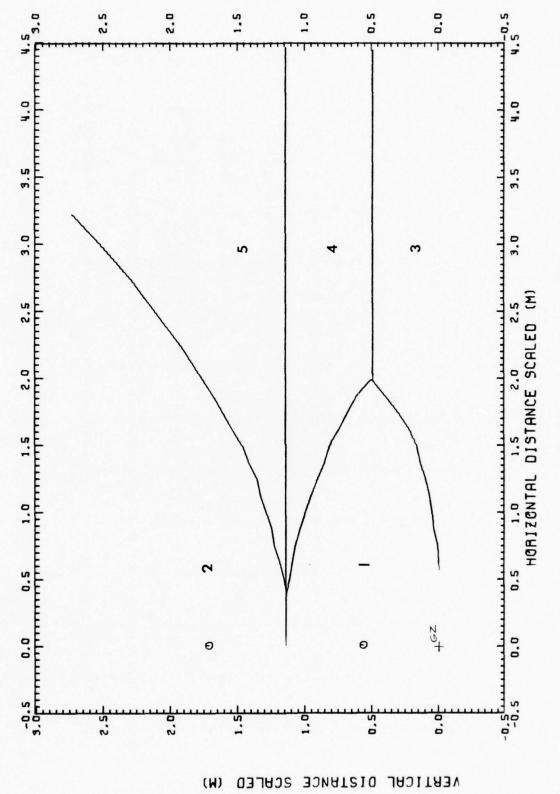


PARTICLE TRAJECTORIES, DIPOLE WEST/10

5

Fig.

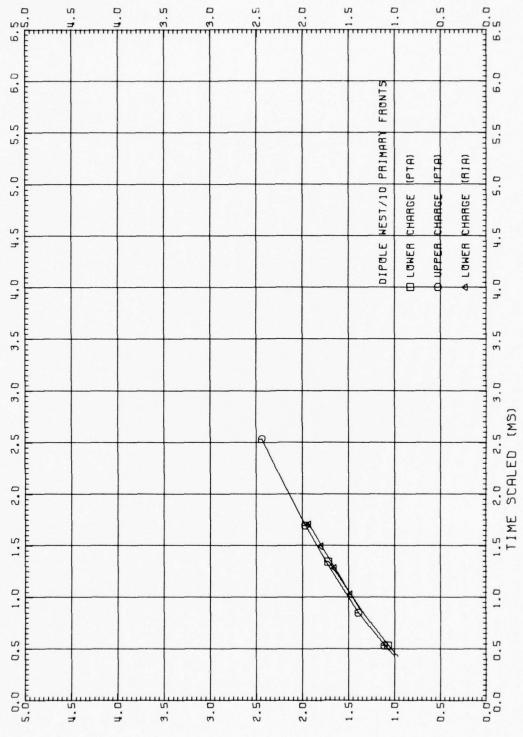
VERTICAL DISTANCE SCALED (M)



REGIONS DEFINITION, DIPOLE WEST/10

9

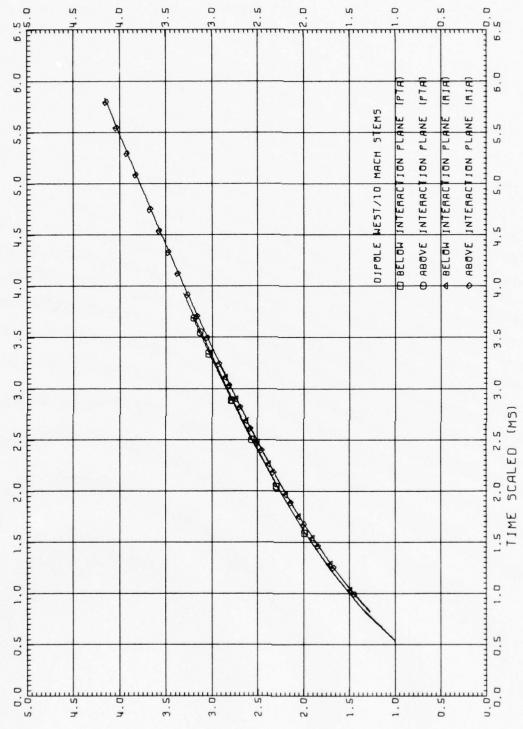
Fig.



RADIUS SCALED (M)

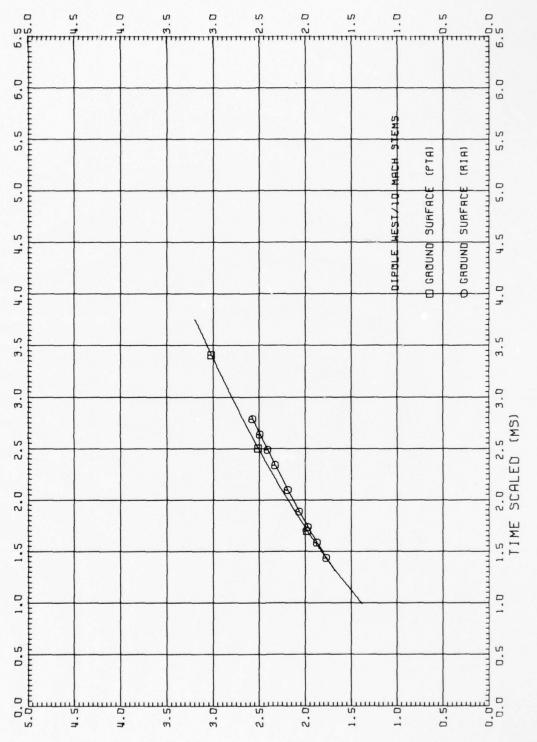
SHOCK TRAJECTORIES, DIPOLE WEST/10

Fig. 7.1



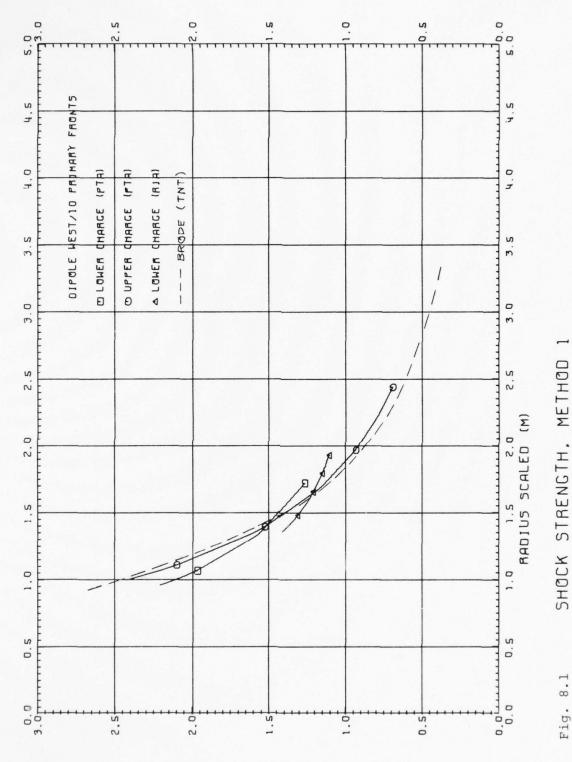
DIPOLE WEST/10

SHOCK TRAJECTORIES,

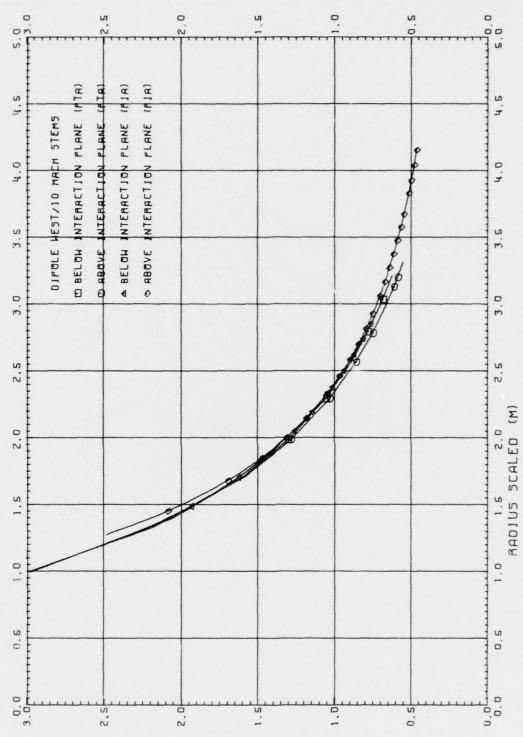


BADINS SCALED (M)

SHOCK TRAJECTORIES, DIPOLE WEST/10



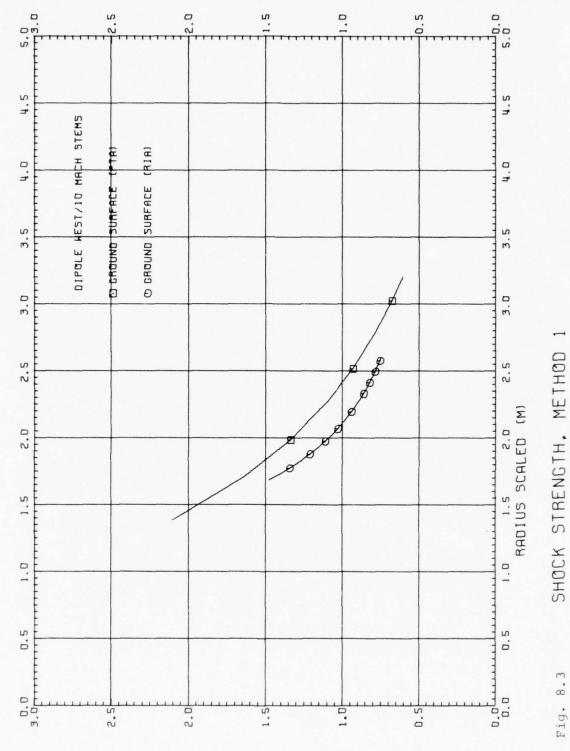
PEAK PARTICLE VELOCITY (MACH UNITS)



SHOCK STRENGTH, METHOD

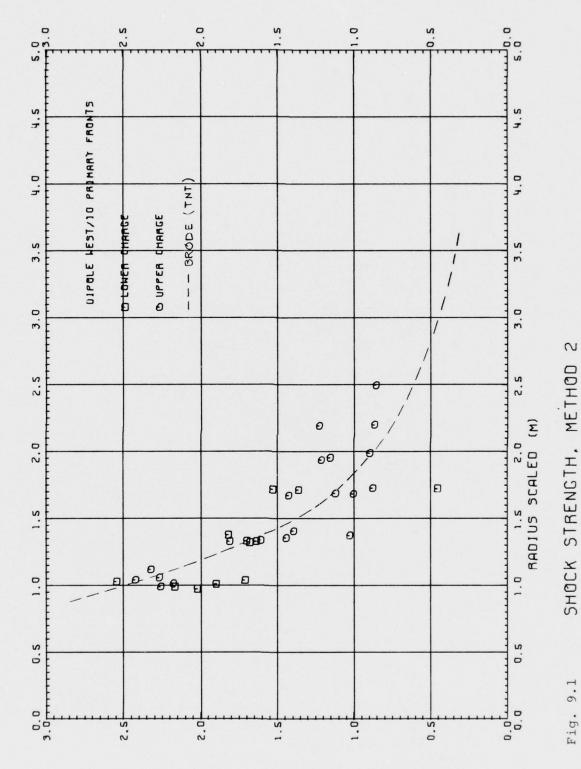
Fig. 8.2

PERK PRRIICLE VELOCITY (MRCH UNITS)

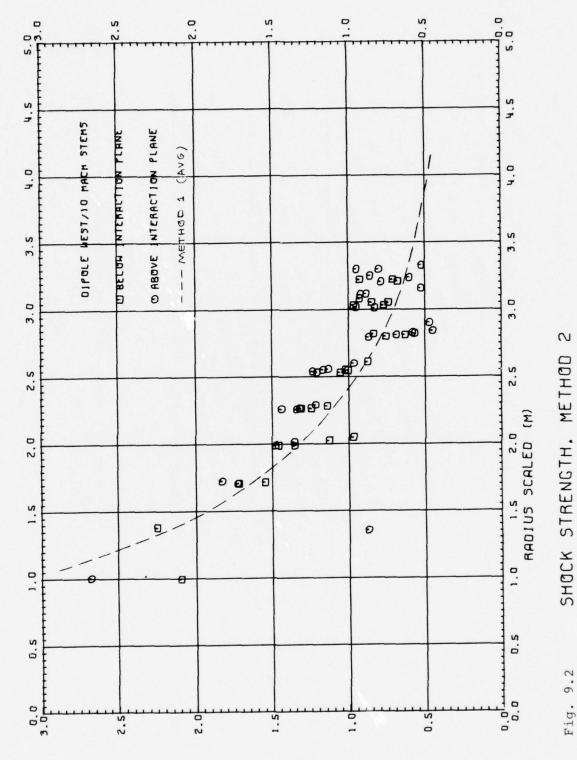


PEAK PARTICLE VELOCITY (MACH UNITS)

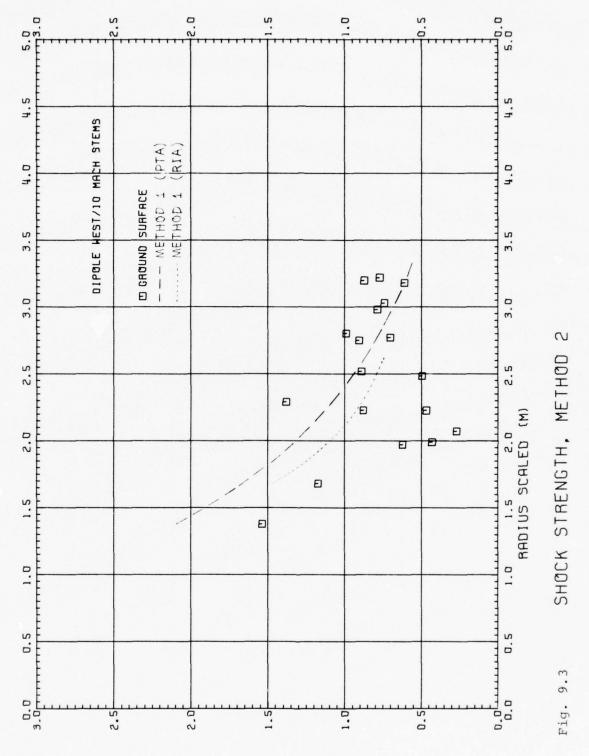
Fig. 8.3



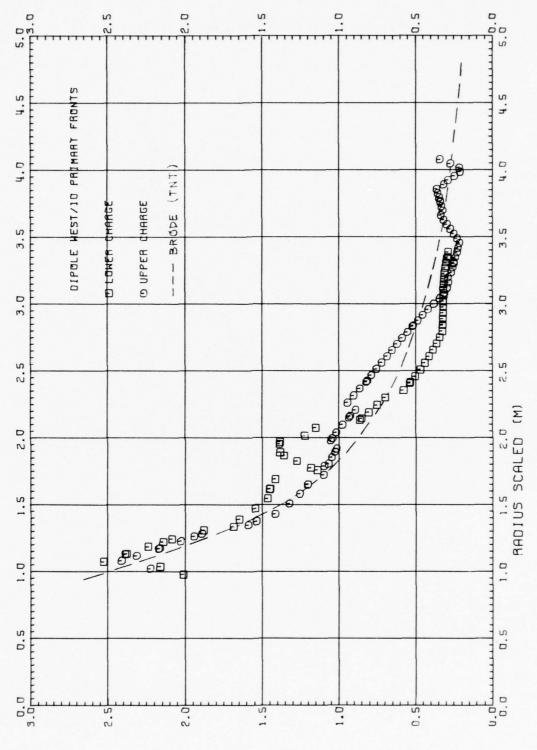
PEAK PARTICLE VELOCITY (MACH UNITS)



PEAK PARTICLE VELOCITY (MACH UNITS)



PEAK PARTICLE VELOCITY (MACH UNITS)

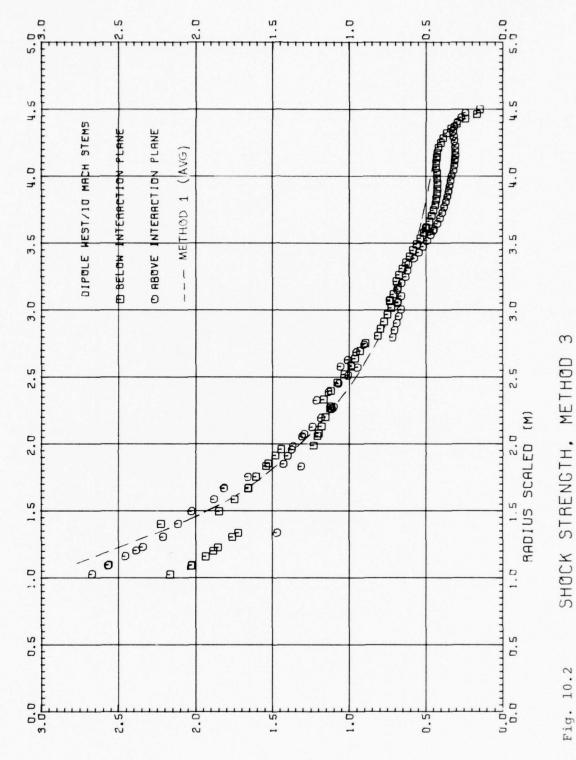


3

SHOCK STRENGTH, METHOD

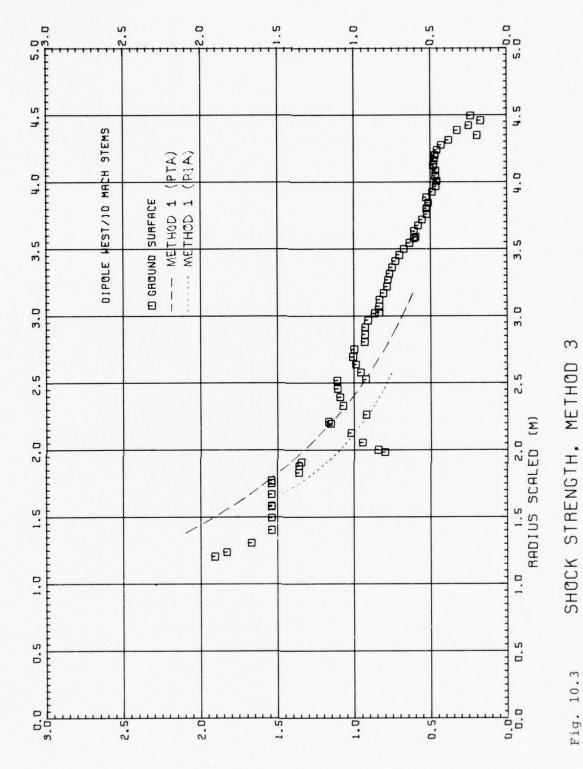
Fig. 10.1

PEAK PARTICLE VELOCITY (MACH UNITS)

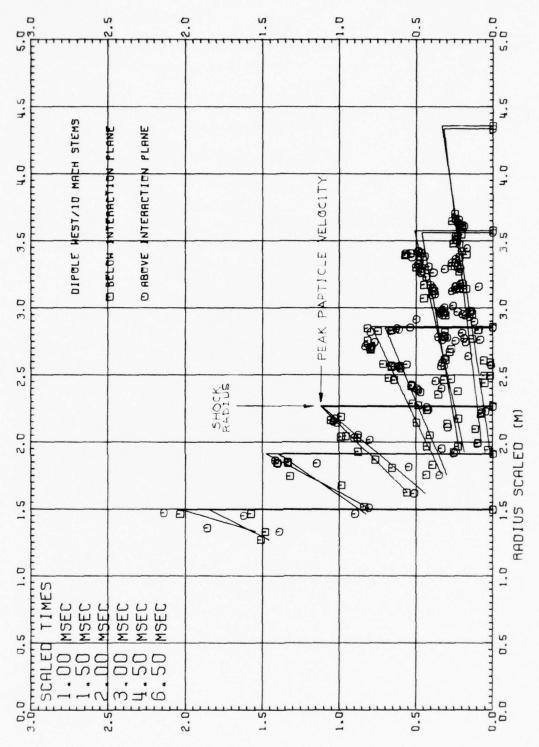


PEAK PARTICLE VELOCITY (MACH UNITS)

Fig. 10.2



PEAK PARTICLE VELOCITY (MACH UNITS)

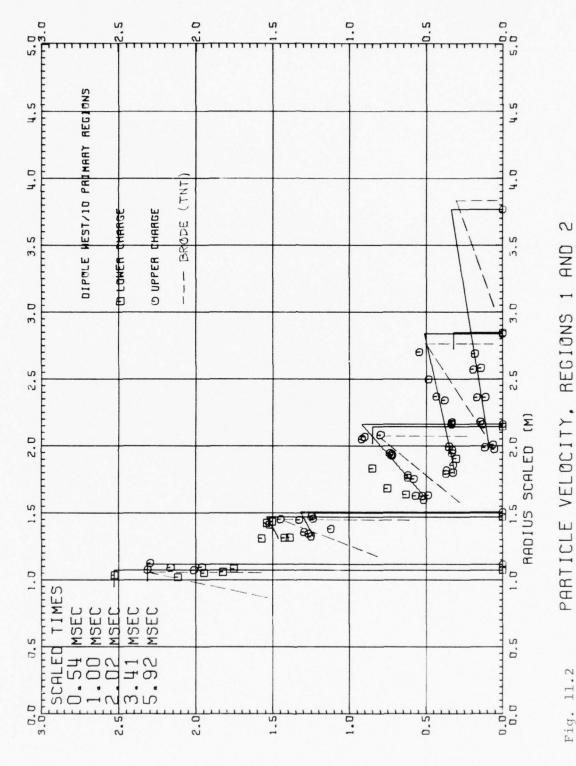


S

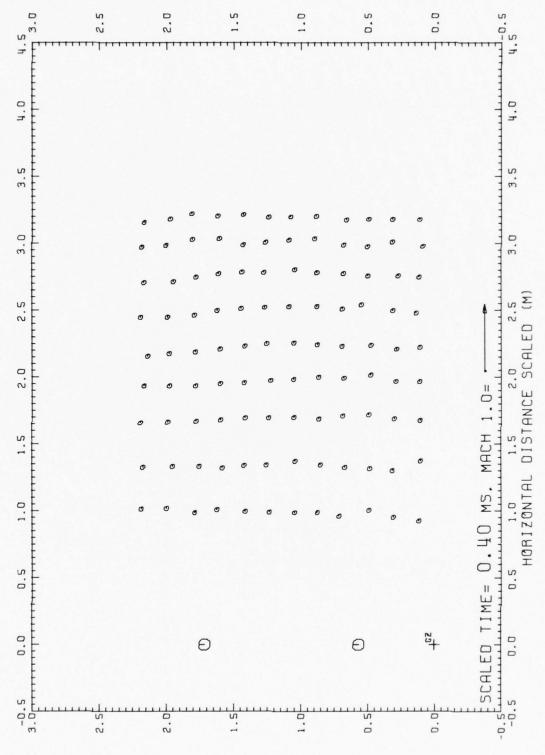
PARTICLE VELOCITY, REGIONS 4 AND

Fig. 11.1

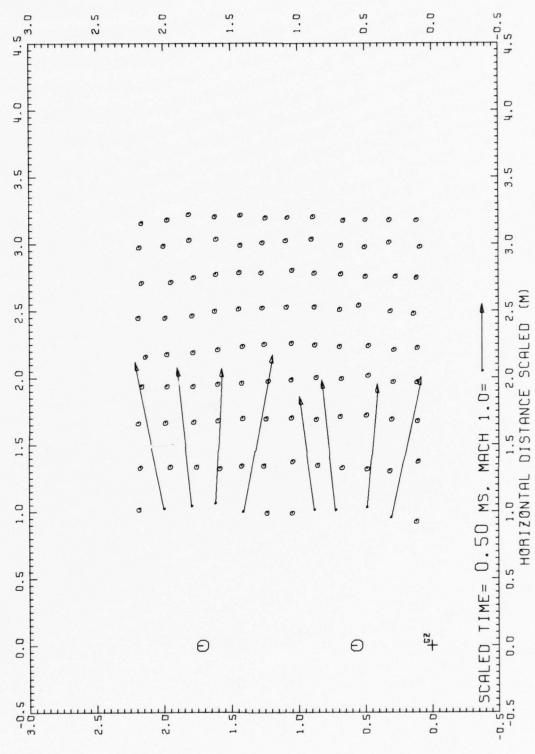
PARTICLE VELOCITY (MACH UNITS)



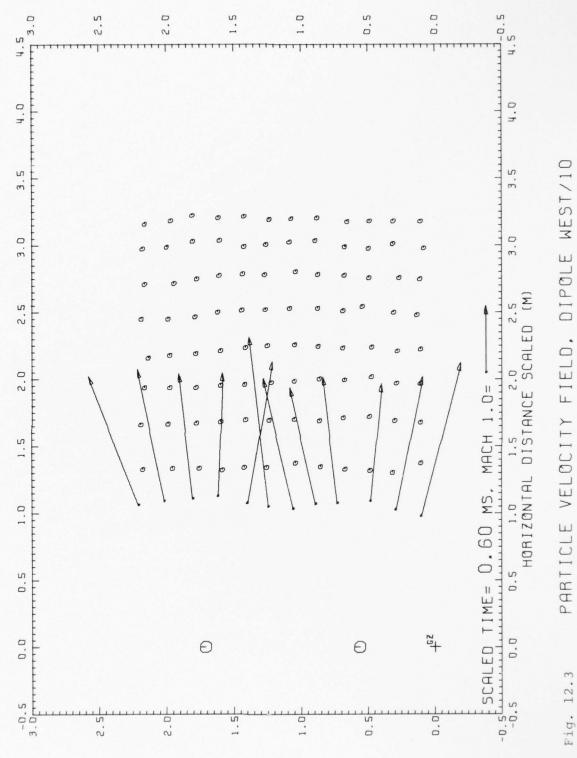
PARTICLE VELOCITY (MACH UNITS)



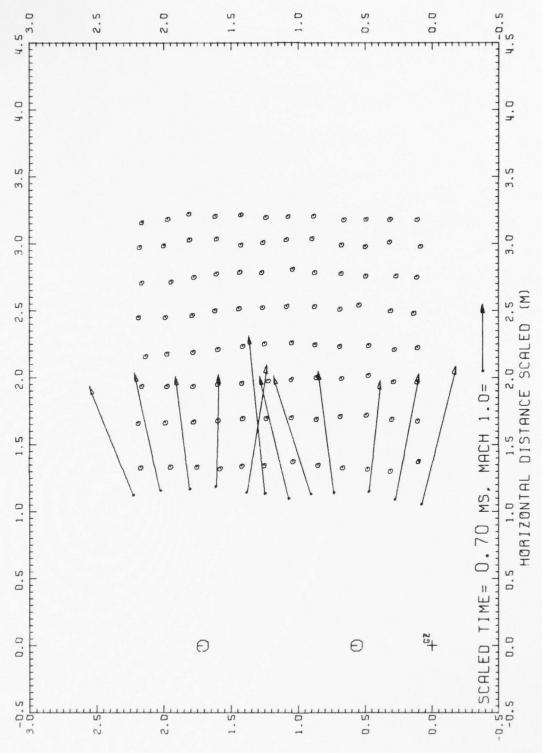
VERTICAL DISTANCE SCALED (M)



VERTICAL DISTANCE SCALED (M



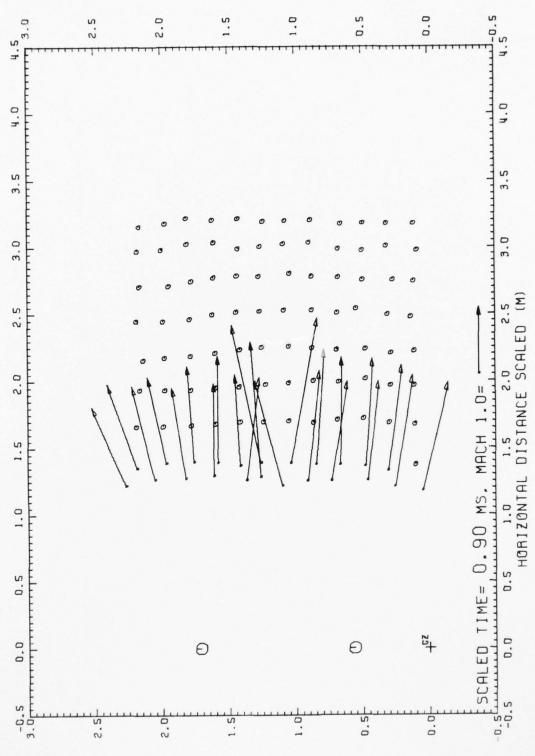
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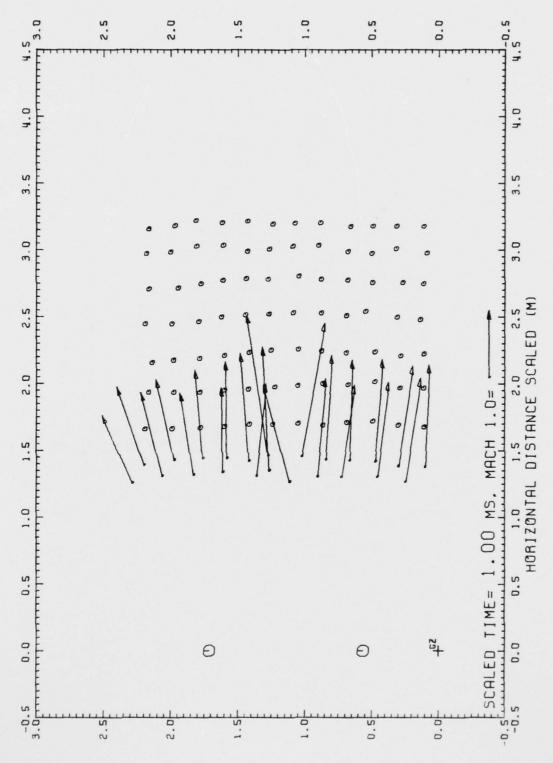
VERTICAL DISTANCE SCALED (M

PARTICLE VELOCITY FIELD, DIPOLE WEST/10 Fig. 12.5

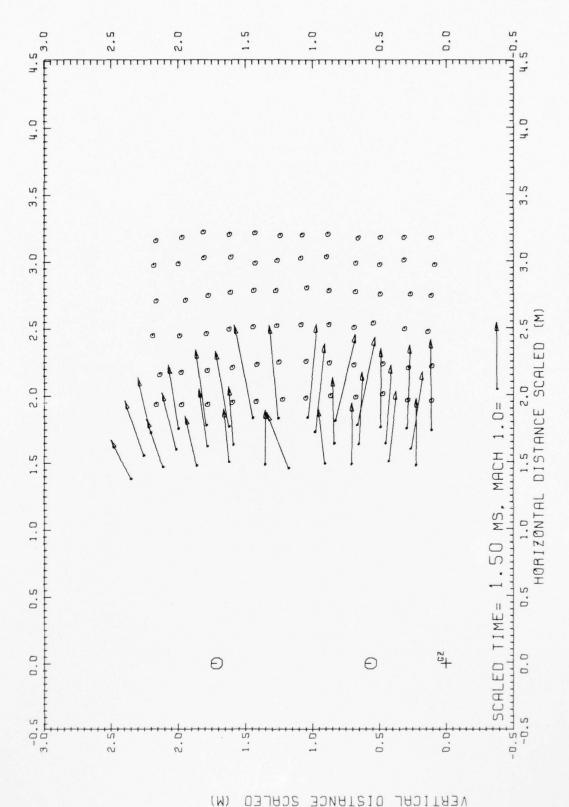
VERTICAL DISTANCE SCALED



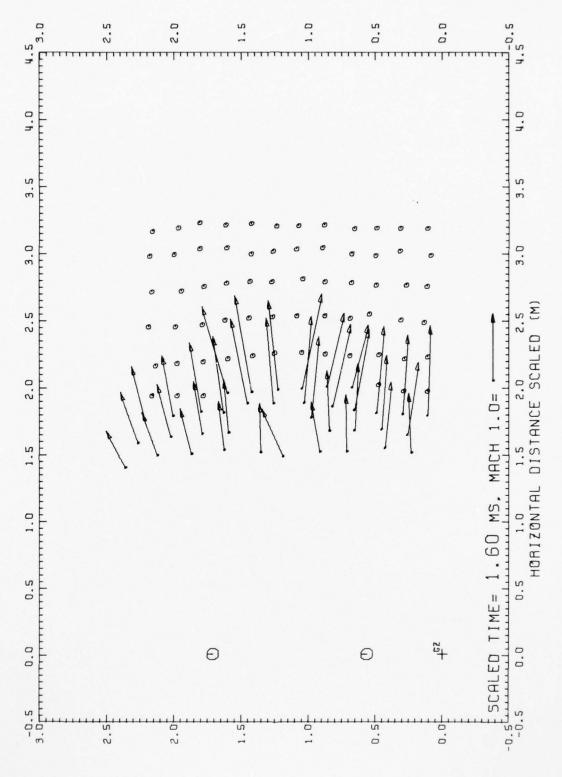
VERTICAL DISTANCE SCALED (M



VERTICAL DISTANCE SCALED (M)

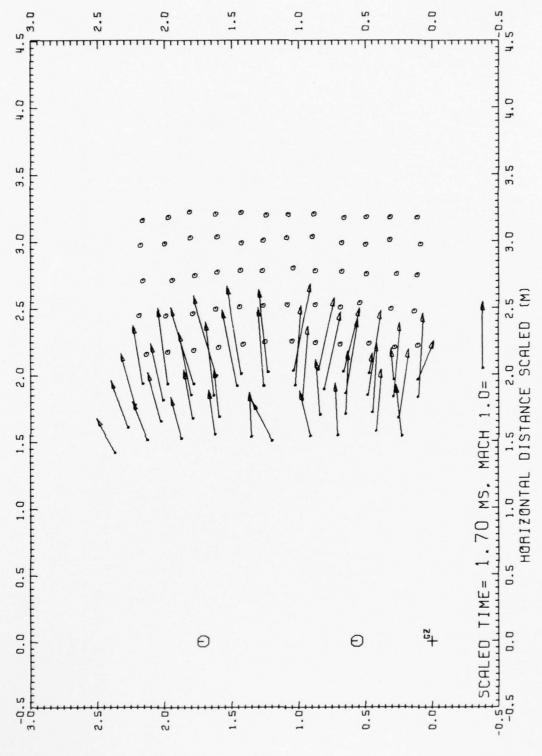


PARTICLE VELOCITY FIELD, DIPOLE WEST/10 Fig. 12.8

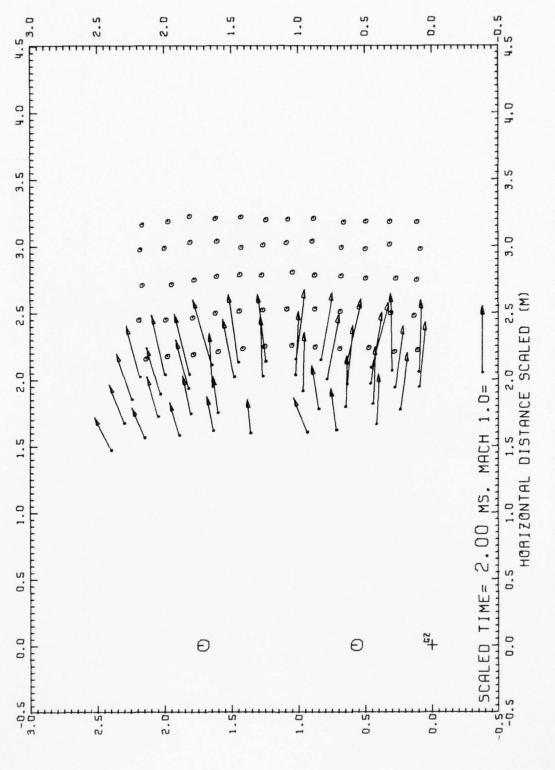


PARTICLE VELOCITY FIELD, DIPOLE WEST/10 Fig. 12.9

VERTICAL DISTANCE SCALED

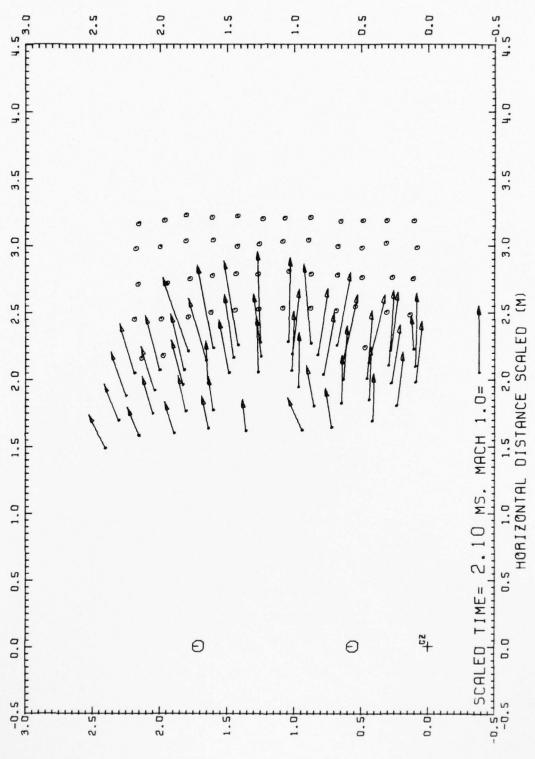


VERTICAL DISTANCE SCALED (M)

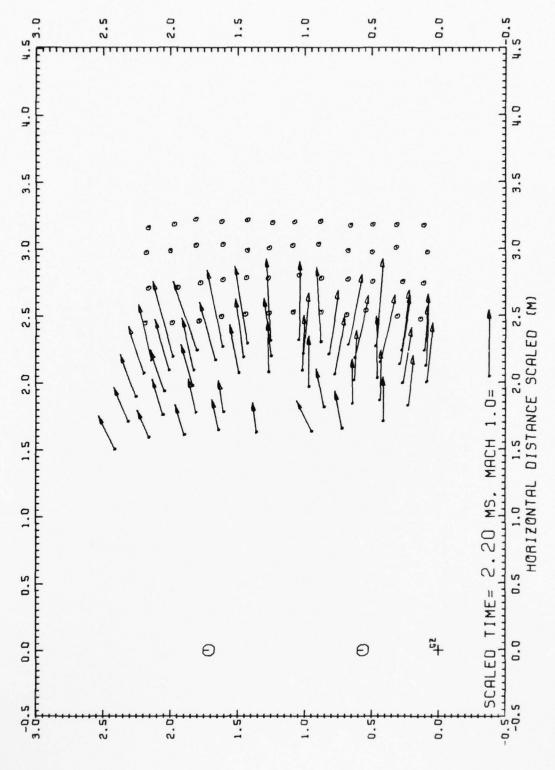


PARTICLE VELOCITY FIELD, DIPOLE WEST/10 Fig. 12.11

VERTICAL DISTANCE SCALED

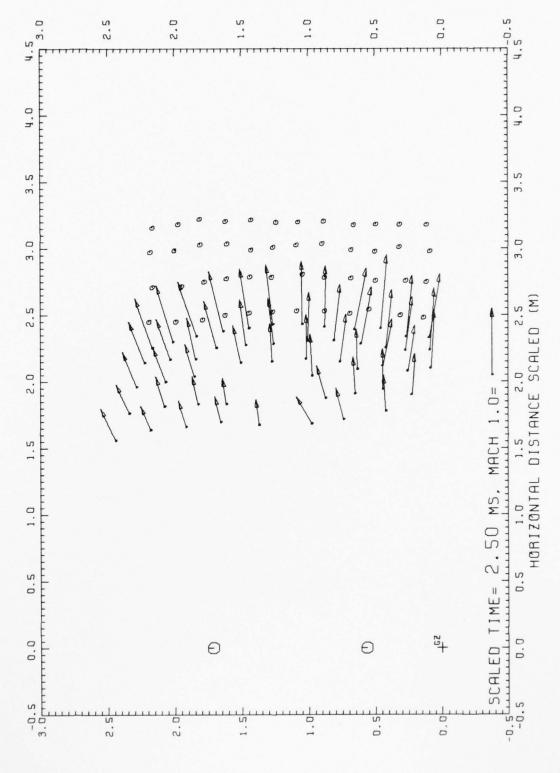


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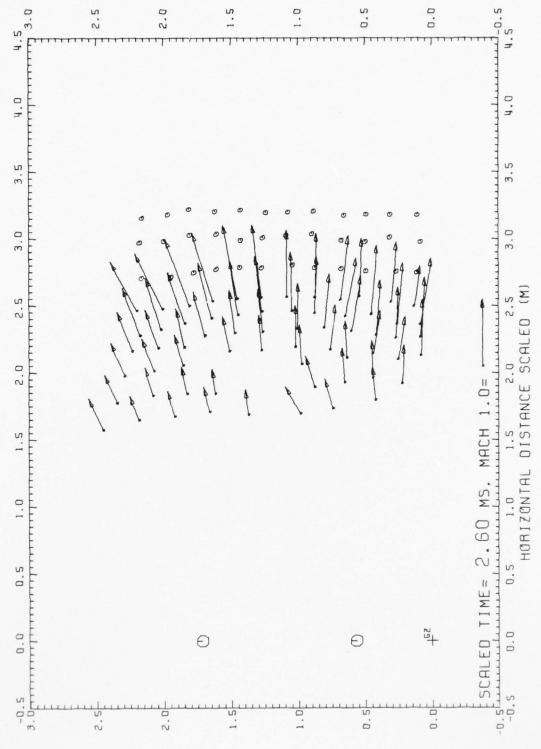
PARTICLE VELOCITY FIELD, DIPOLE WEST/10 Fig. 12.13

VERTICAL DISTANCE SCALED

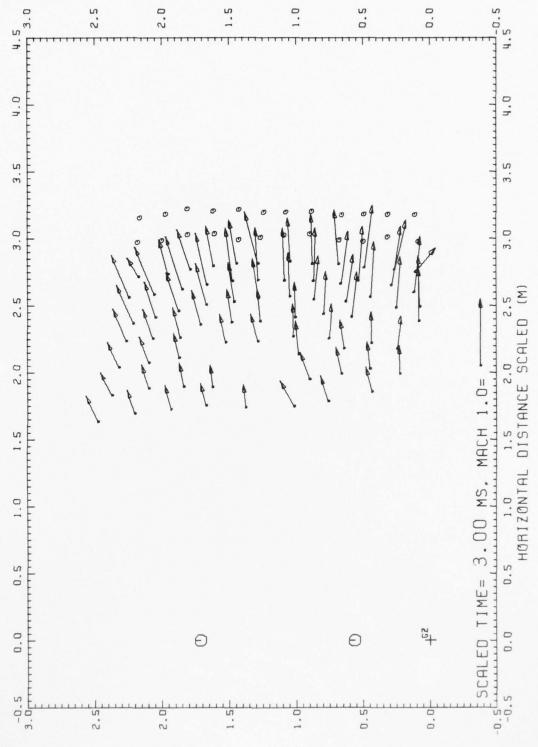


PARTICLE VELOCITY FIELD, DIPOLE WEST/10 Fig. 12.14

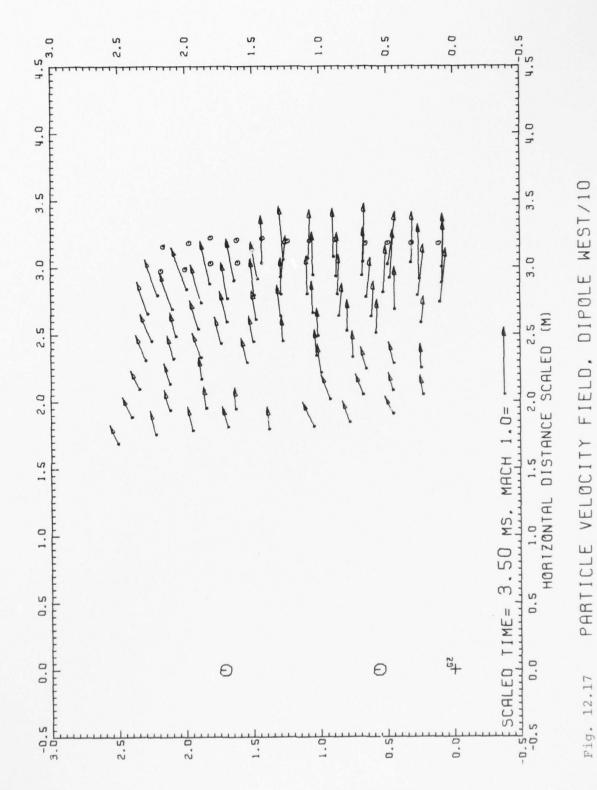
VERTICAL DISTANCE SCALED (M)



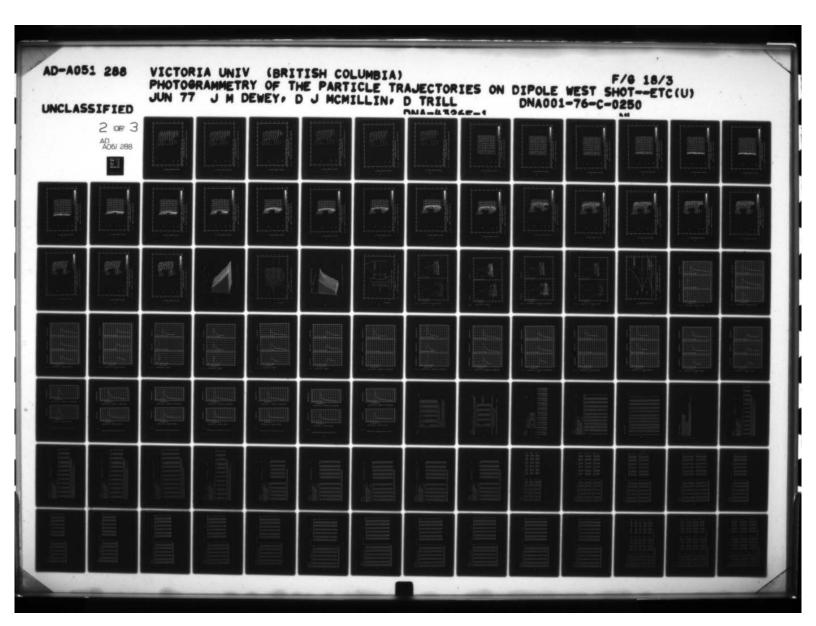
VERTICAL DISTANCE SCALED (M)

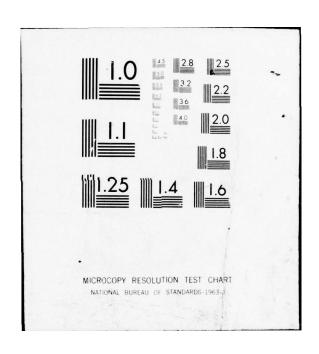


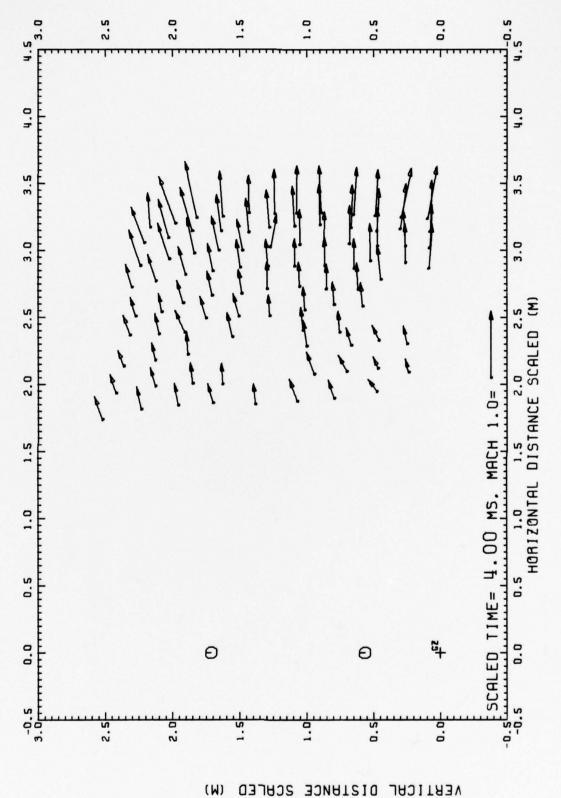
VERTICAL DISTANCE SCALED



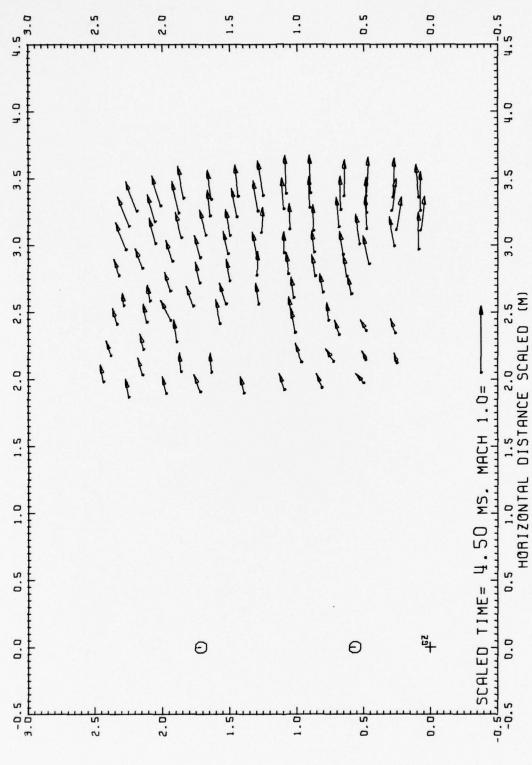
VERTICAL DISTANCE SCALED (M)



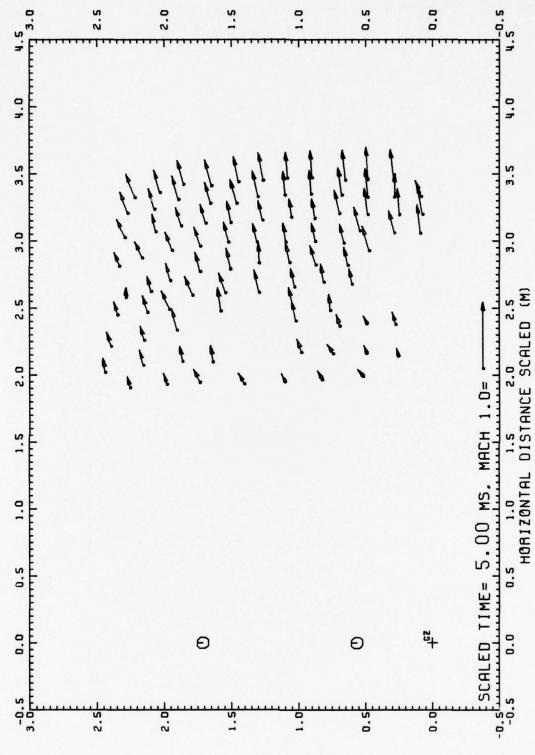




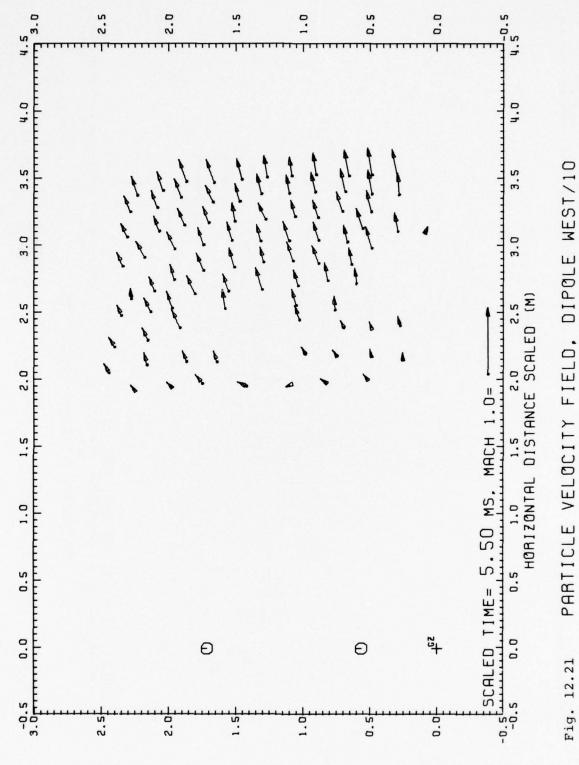
PARTICLE VELOCITY FIELD, DIPOLE WEST/10 Fig. 12.18



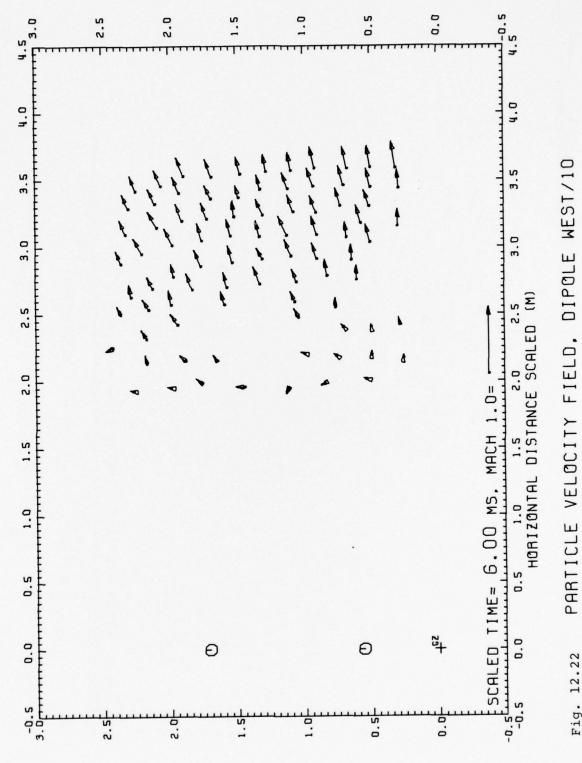
VERTICAL DISTANCE SCALED (M)



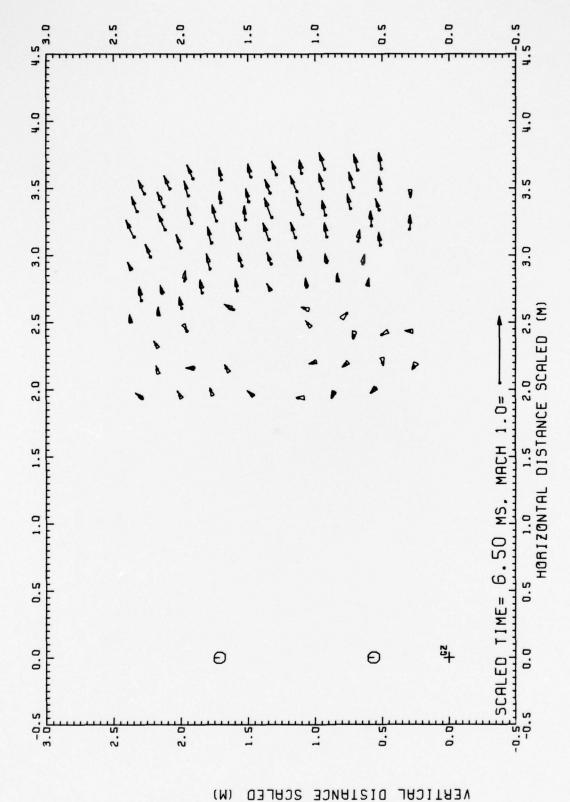
VERTICAL DISTANCE SCALED (M.



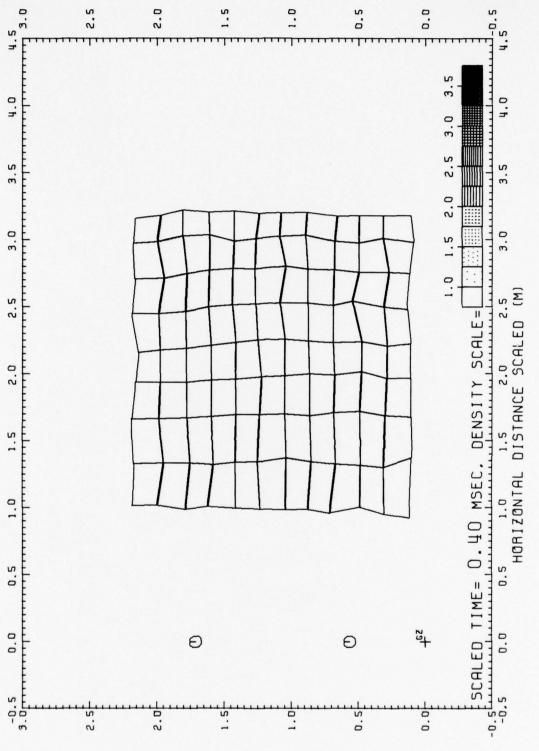
VERTICAL DISTANCE SCALED (M



VERTICAL DISTANCE SCALED (M)

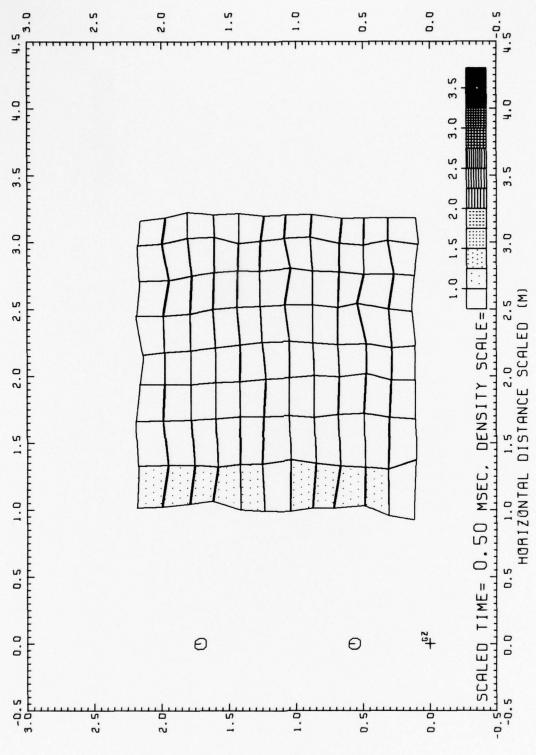


PARTICLE VELOCITY FIELD, DIPOLE WEST/10 Fig. 12.23



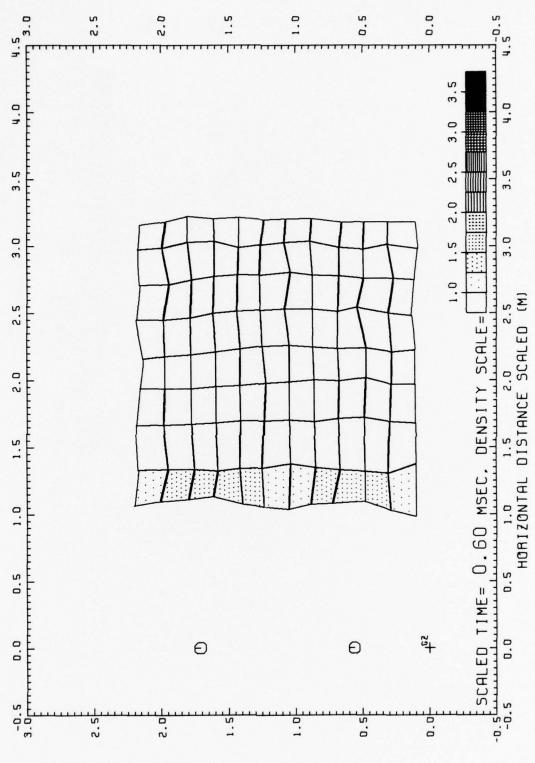
VERTICAL DISTANCE SCALED (M)

DENSITY FIELD, DIPOLE WEST/10



VERTICAL DISTANCE SCALED (M)

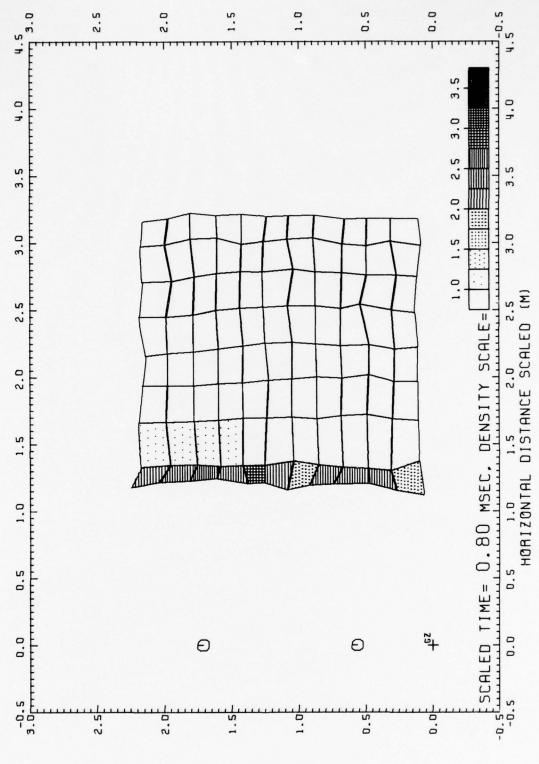
DENSITY FIELD, DIPOLE WEST/10



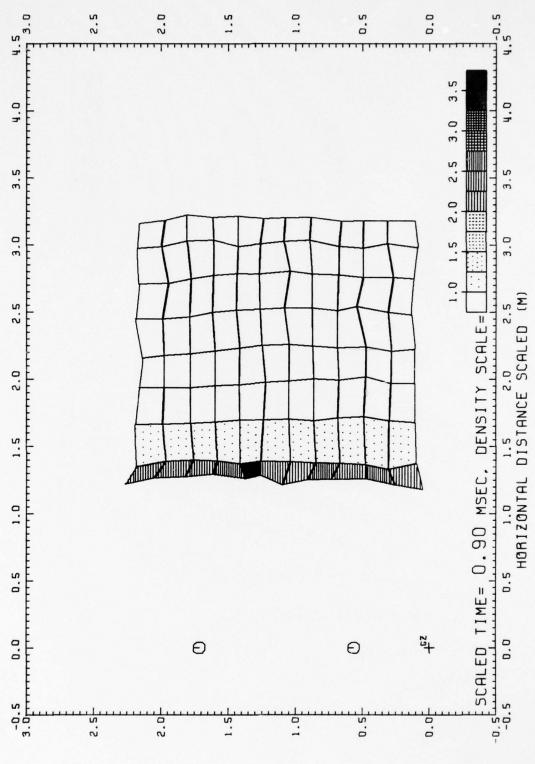
VERTICAL DISTANCE SCALED (M

4 DENSITY FIELD, DIPOLE WEST/10

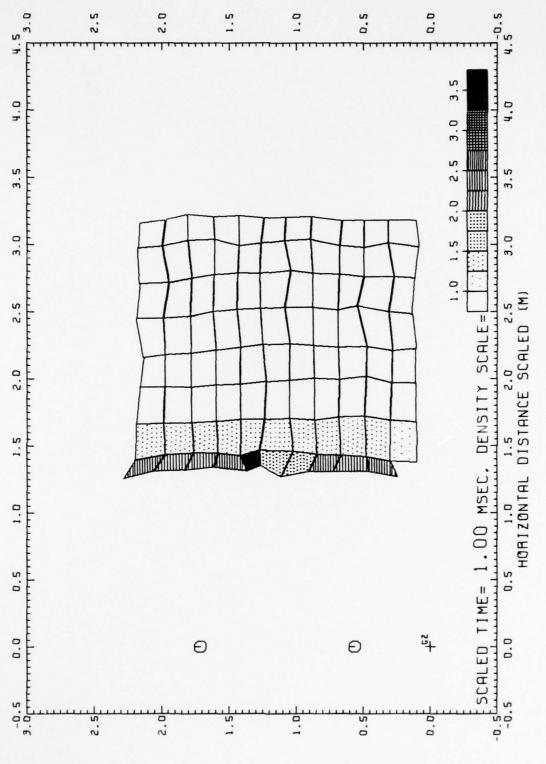
VERTICAL DISTANCE SCALED



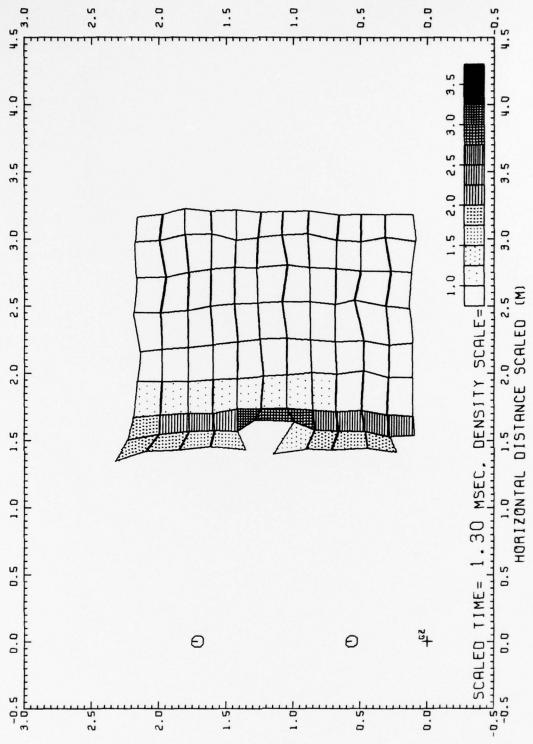
VERTICAL DISTANCE SCALED (M



VERTICAL DISTANCE SCALED (M)

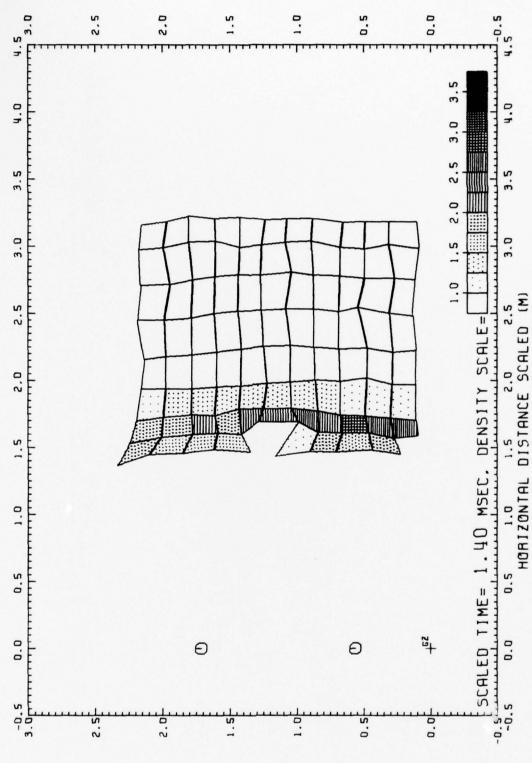


VERTICAL DISTANCE SCALED (M)

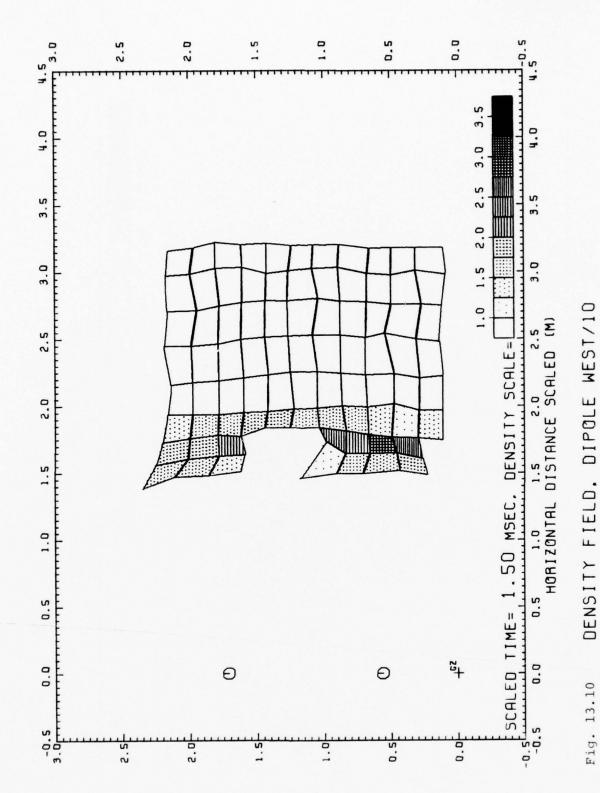


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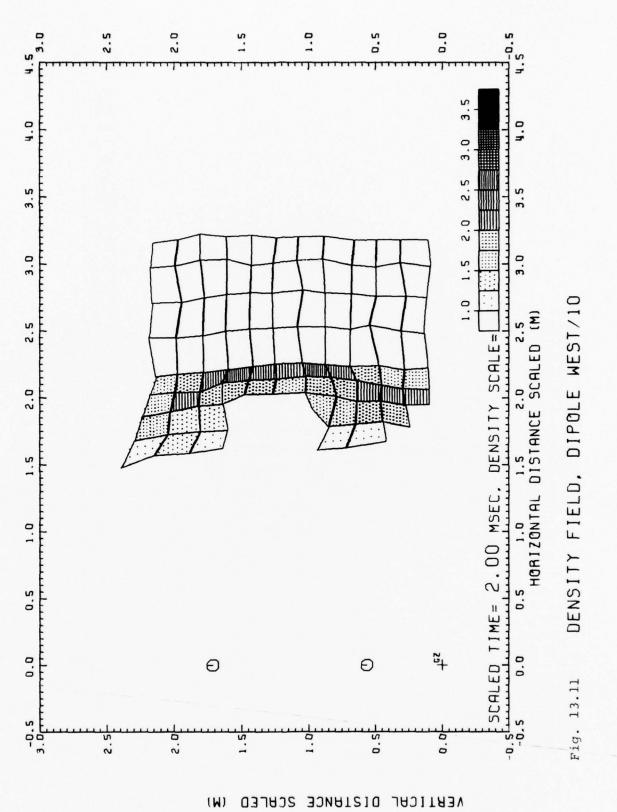
VERTICAL DISTANCE SCALED (M)



VERTICAL DISTANCE SCALED (M)



VERTICAL DISTANCE SCALED (M)



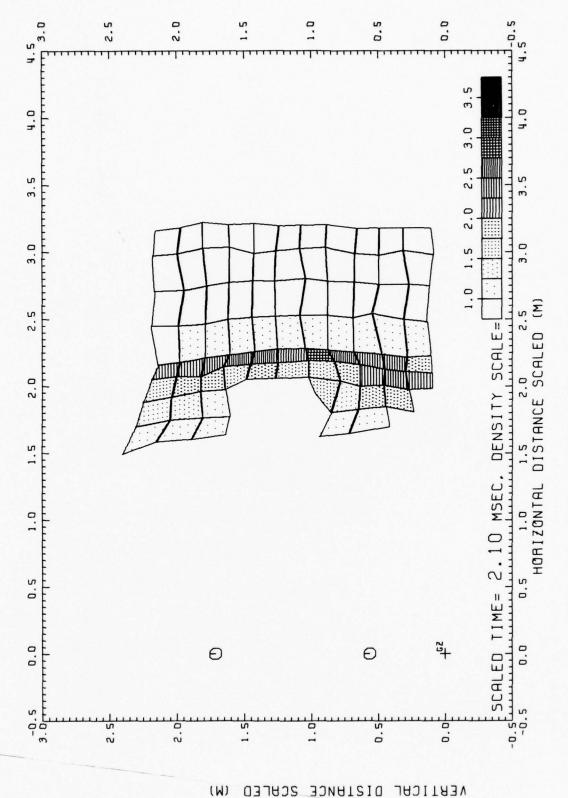
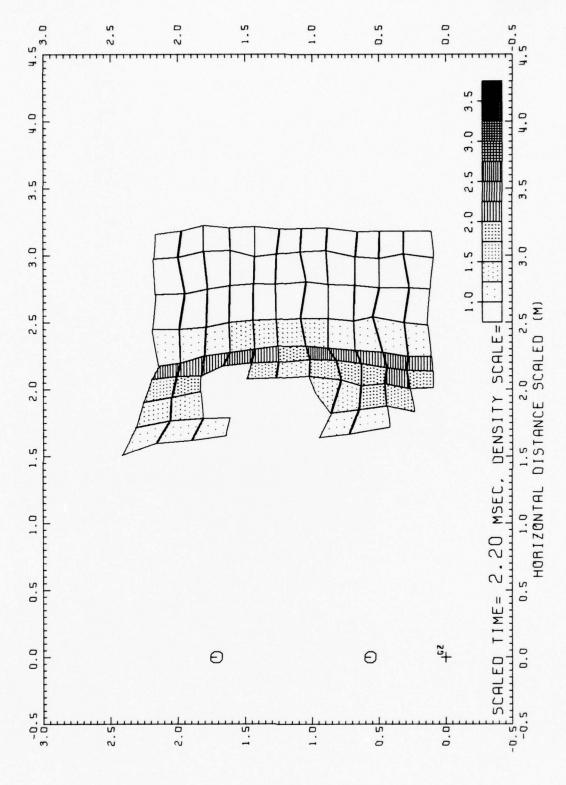
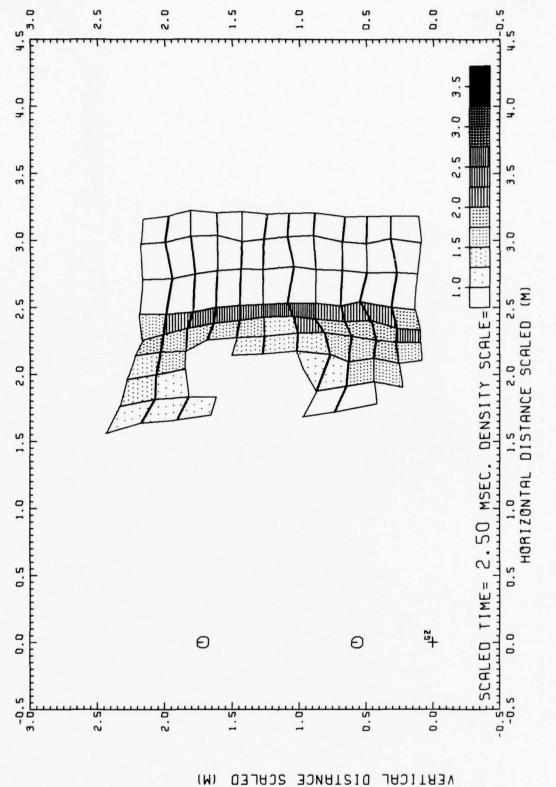


Fig. 13.12 DENSITY FIELD, DIPOLE WEST/10

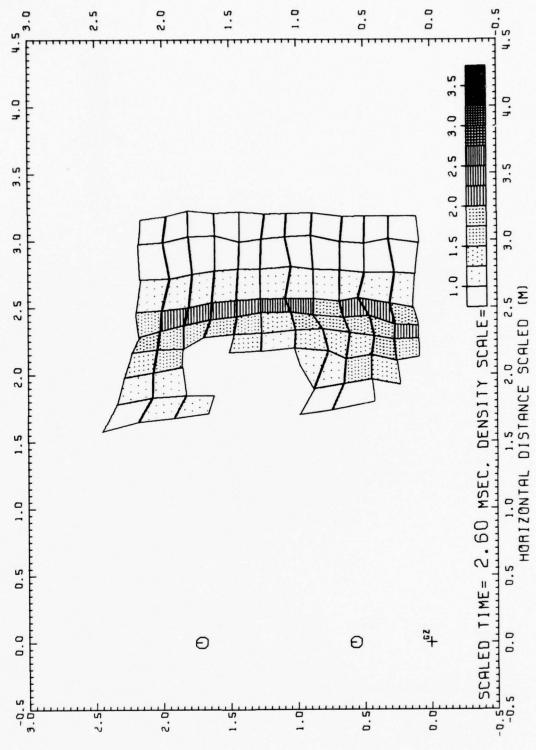


13.13 DENSITY FIELD, DIPOLE WEST/10

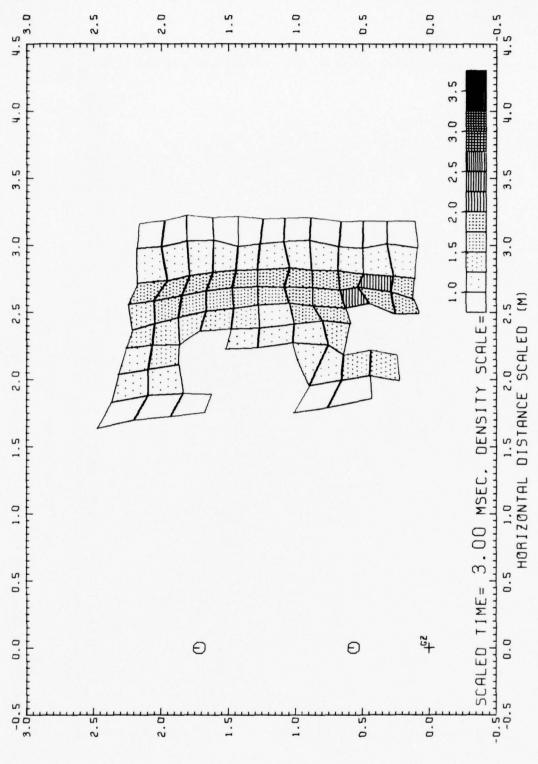
VERTICAL DISTANCE SCALED (M)



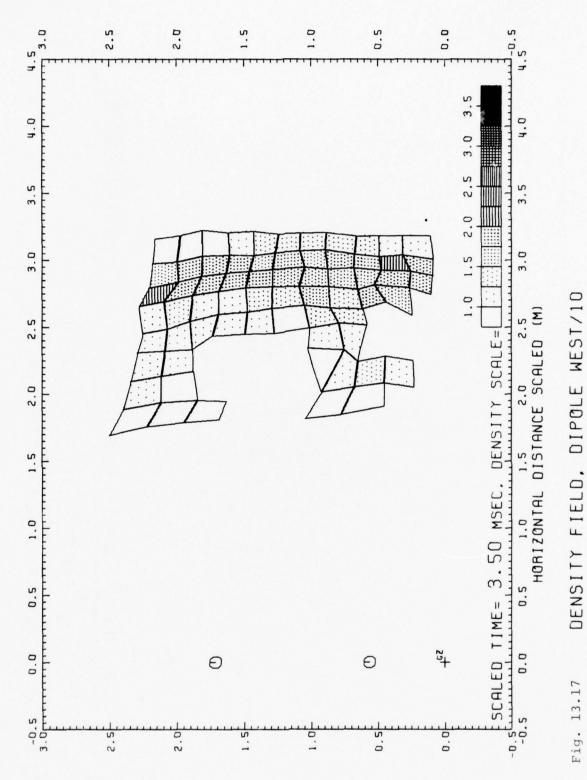
13.14 DENSITY FIELD, DIPOLE WEST/10



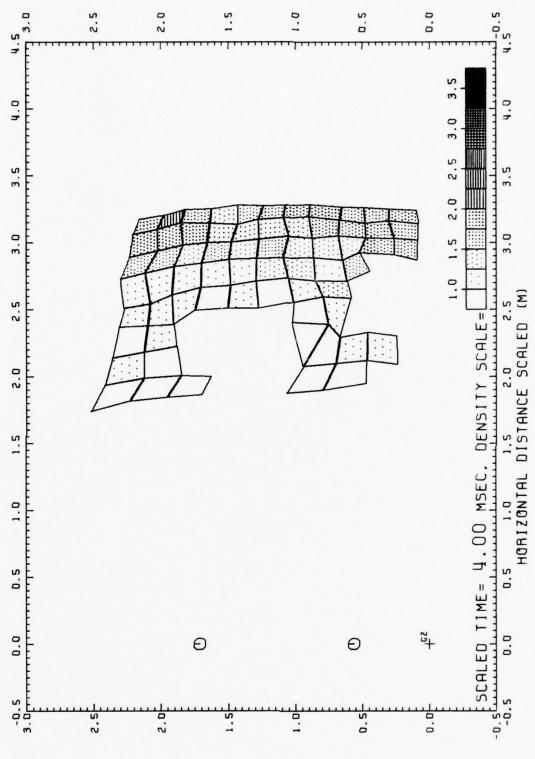
VERTICAL DISTANCE SCALED (M)



VERTICAL DISTANCE SCALED (M



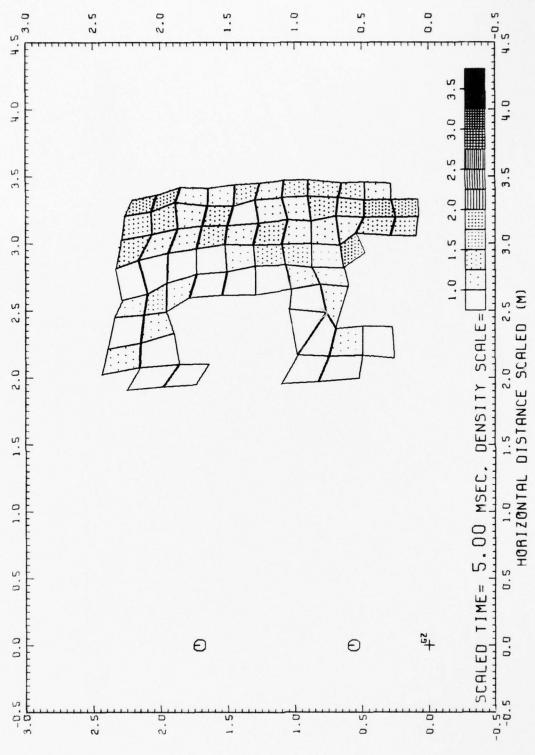
VERTICAL DISTANCE SCALED (M)



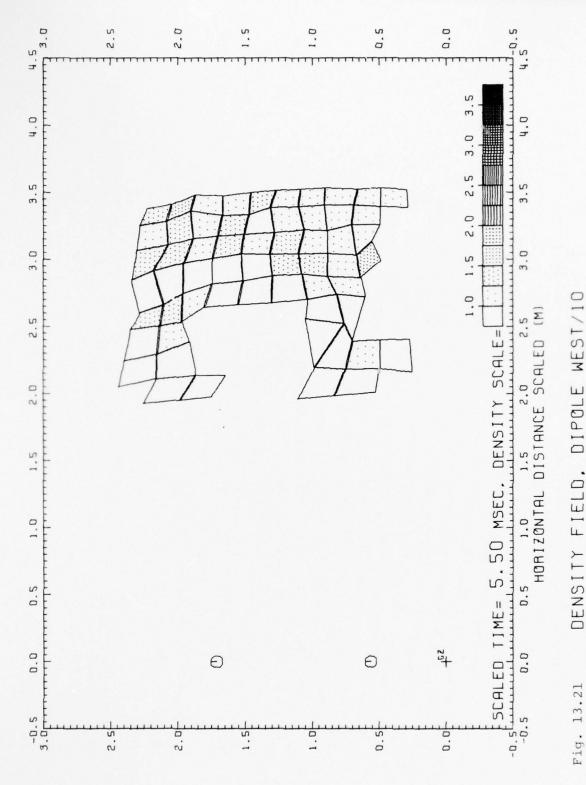
VERTICAL DISTANCE SCALED (M)

Fig. 13.19 DENSITY FIELD, DIPOLE WEST/10

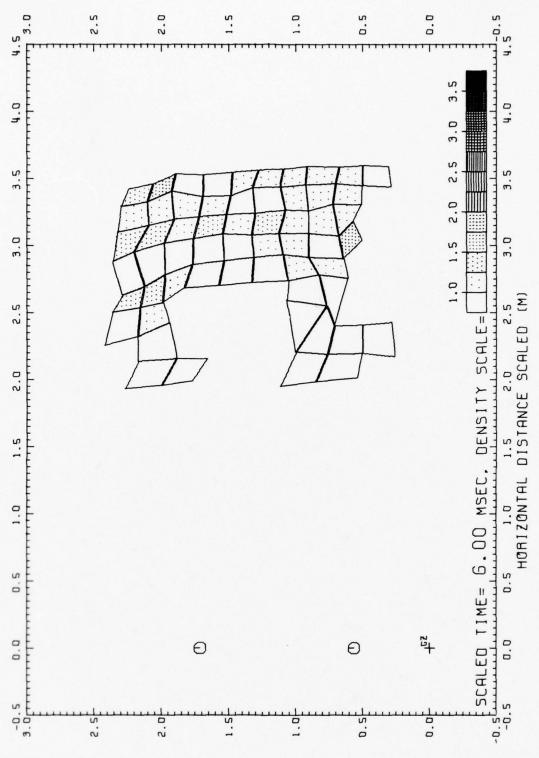
VERTICAL DISTANCE SCALED



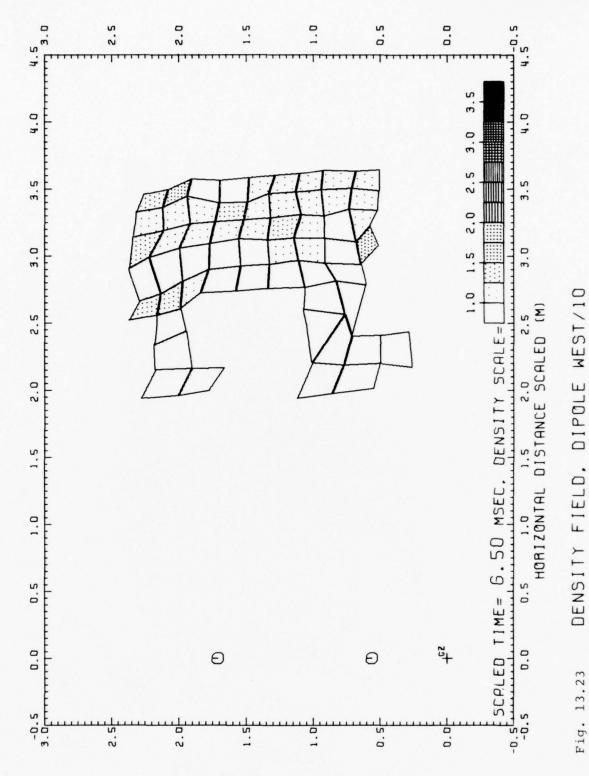
VERTICAL DISTANCE SCALED (M.



VERTICAL DISTANCE SCALED (M)

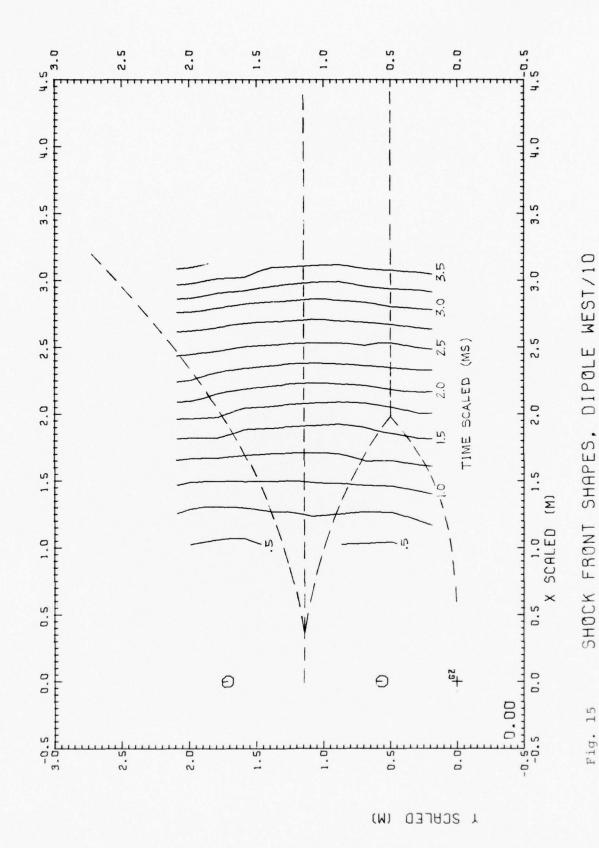


VERTICAL DISTANCE SCALED (M)



VERTICAL DISTANCE SCALED (M)

Fig. 14 Time-of-arrival surface, Dipole West/10



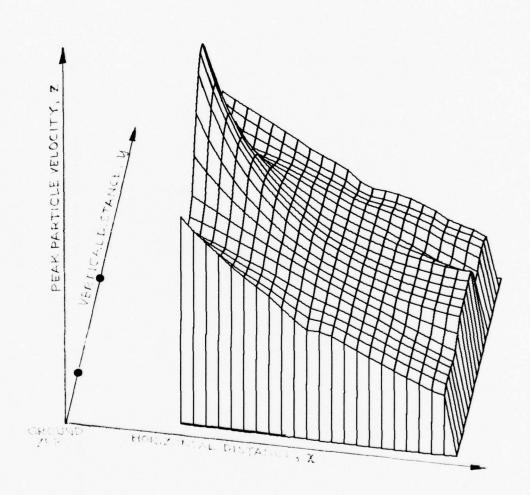
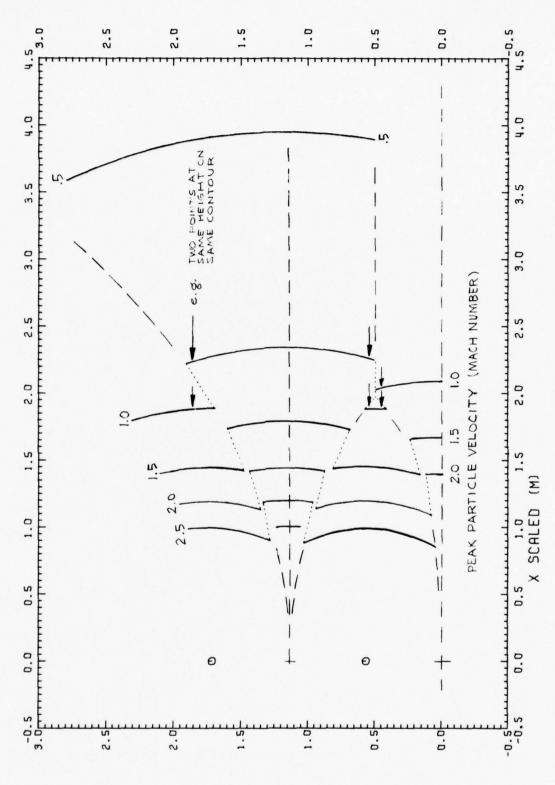


Fig. 16 A shock strength surface, Dipole West/10



L SCHEED (W)

SHOCK STRENGTH CONTOURS, DIPOLE WEST/10

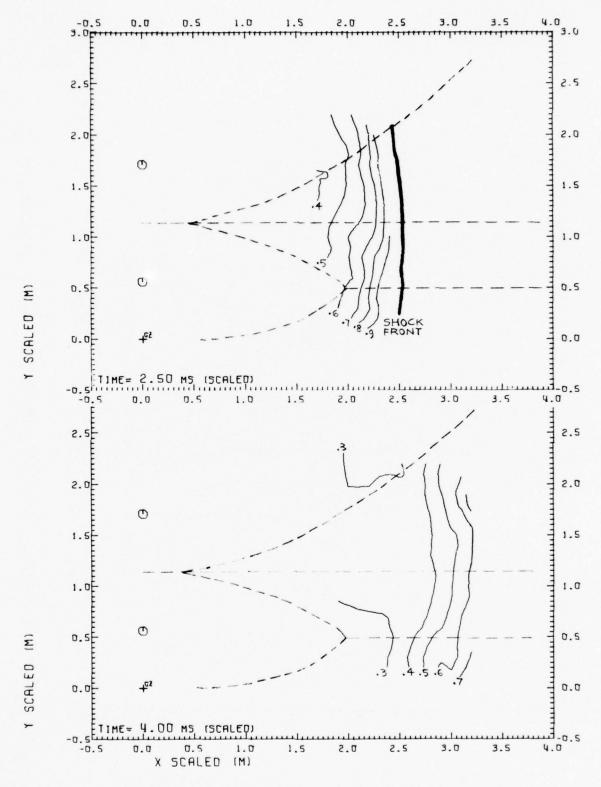


Fig. 18 DIPOLE WEST/10 PARTICLE VELOCITY
128

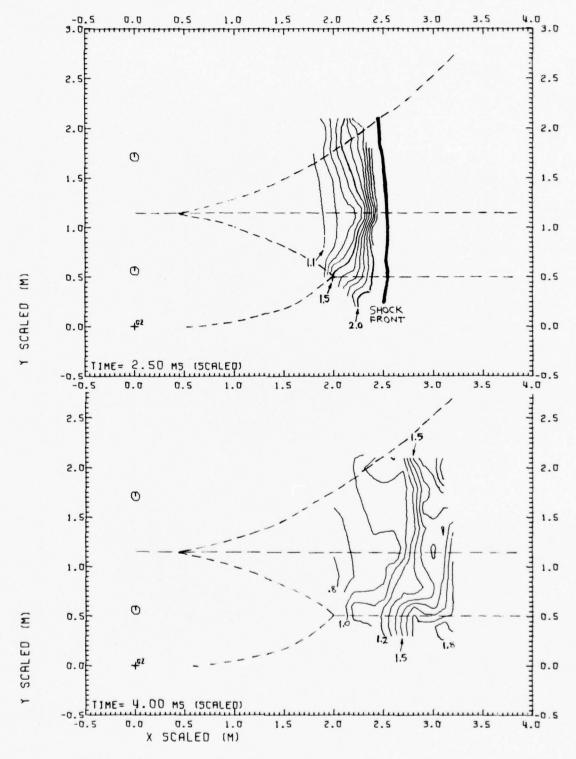


Fig. 19 DIPOLE WEST/10 DENSITY

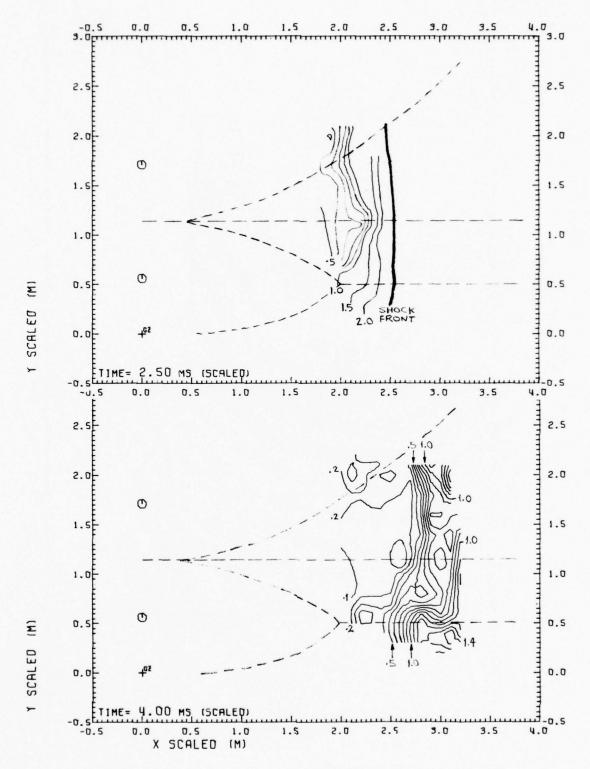


Fig. 20 DIPOLE WEST/10 HYDROSTATIC OVERPRESSURE

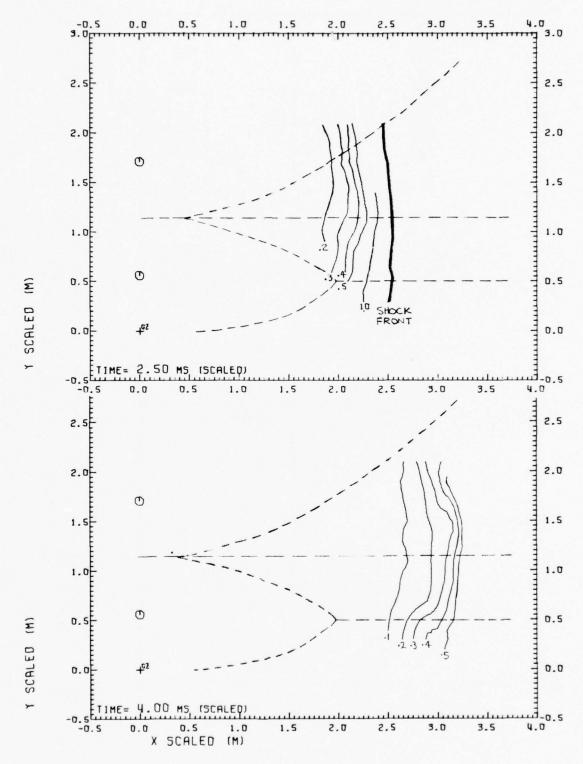
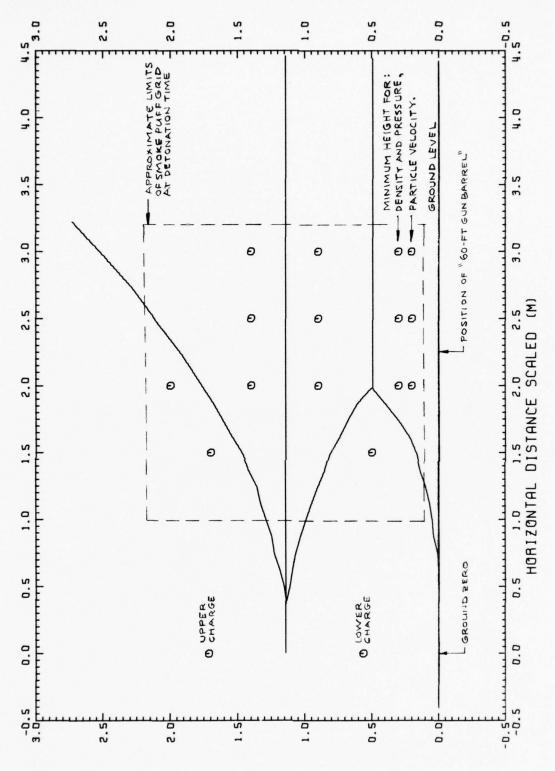


Fig. 21 DIPOLE WEST/10 DYNAMIC PRESSURE



TIME HISTORY STATIONS, DIPOLE WEST/10

Fig. 22

VERTICAL DISTANCE SCALED (M)

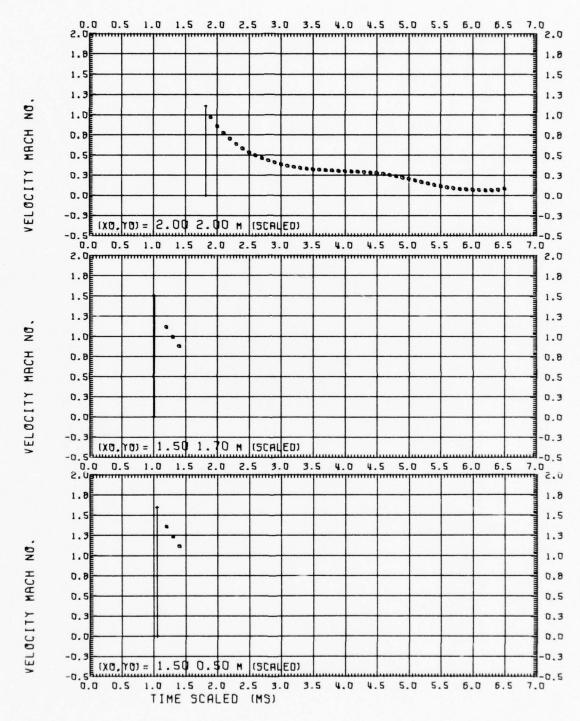


Fig. 23.1 DIPOLE WEST/10 PARTICLE VELOCITY

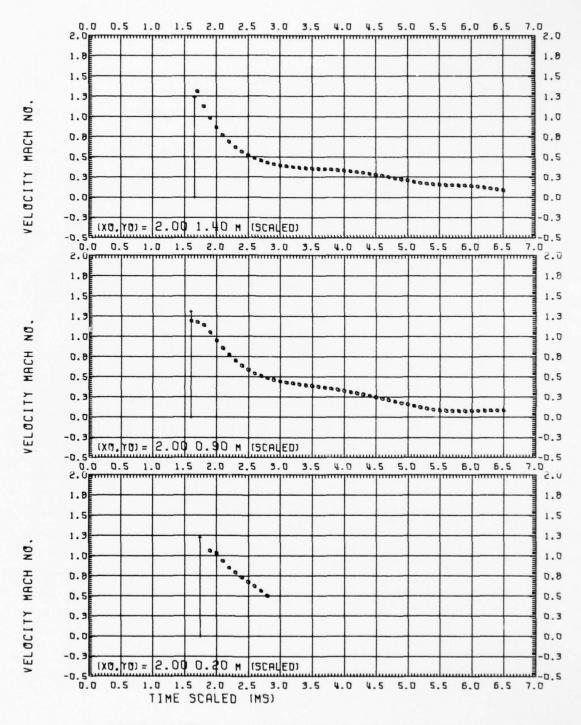


Fig. 23.2 DIPOLE WEST/10 PARTICLE VELOCITY

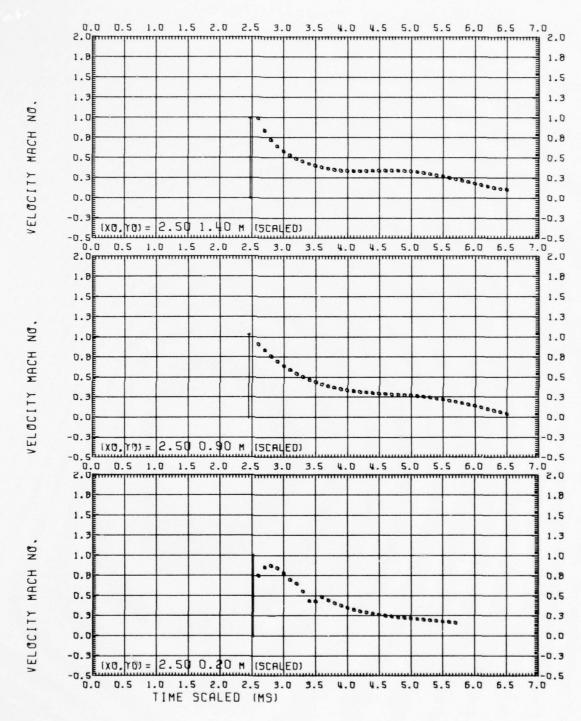


Fig. 23.3 DIPOLE WEST/10 PARTICLE VELOCITY

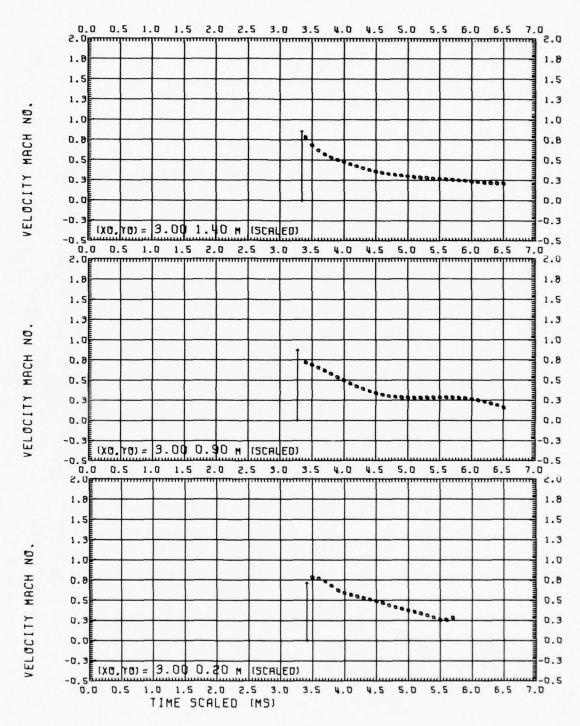


Fig. 23.4 DIPOLE WEST/10 PARTICLE VELOCITY

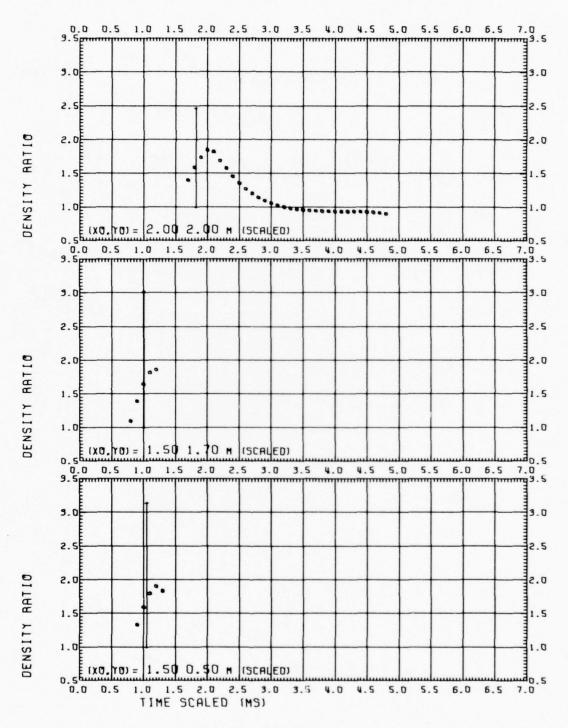


Fig. 24.1 DIPOLE WEST/10 DENSITY

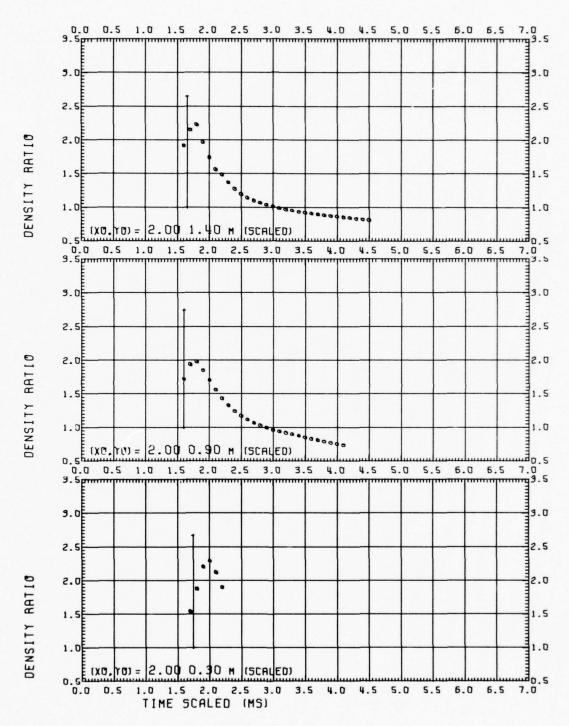


Fig. 24.2 DIPOLE WEST/10 DENSITY

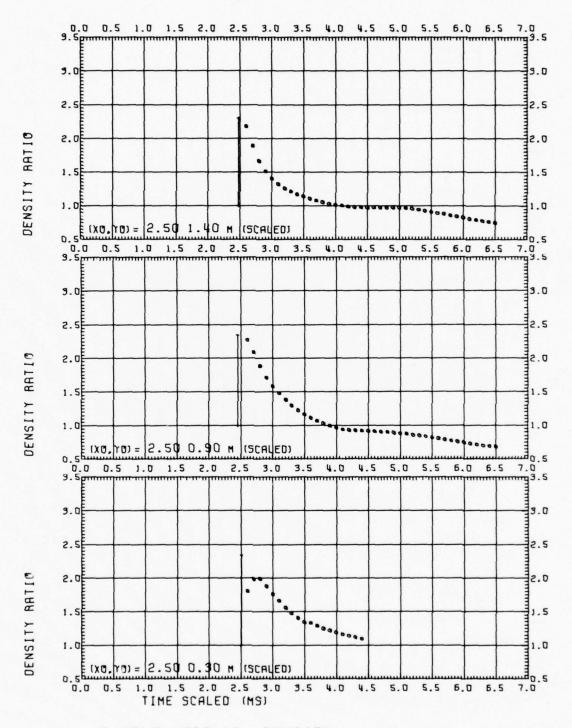


Fig. 24.3 DIPOLE WEST/10 DENSITY

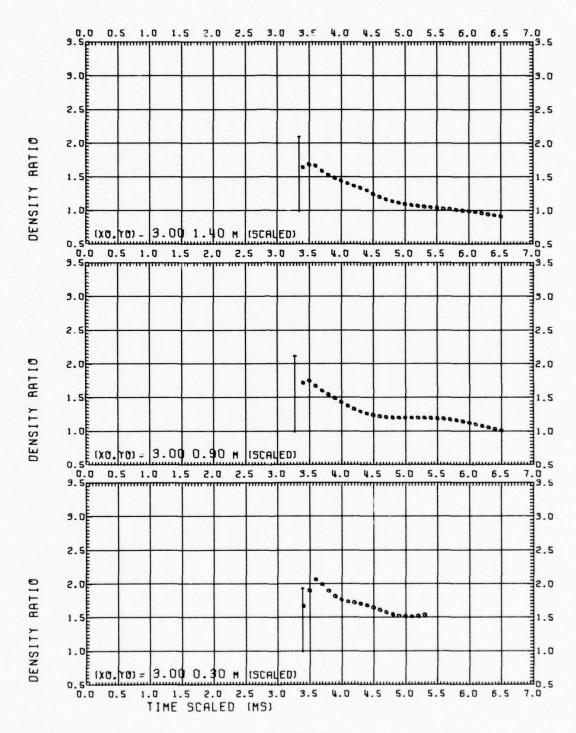


Fig. 24.4 DIPOLE WEST/10 DENSITY

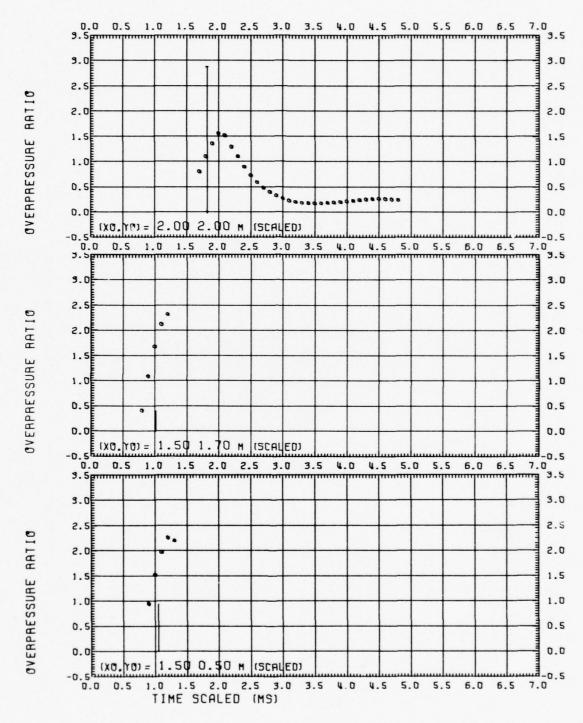


Fig. 25.1 DIPOLE WEST/10 HYDROSTATIC OVERPRESSURE

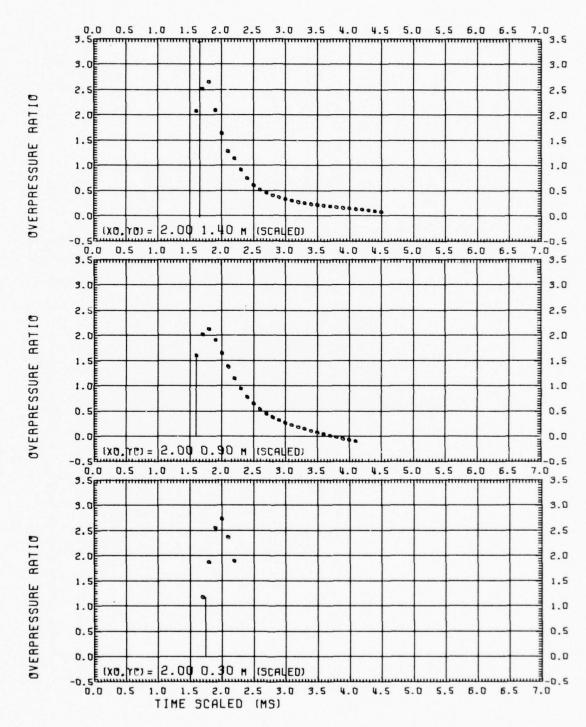


Fig. 25.2 DIPOLE WEST/10 HYDROSTATIC OVERPRESSURE

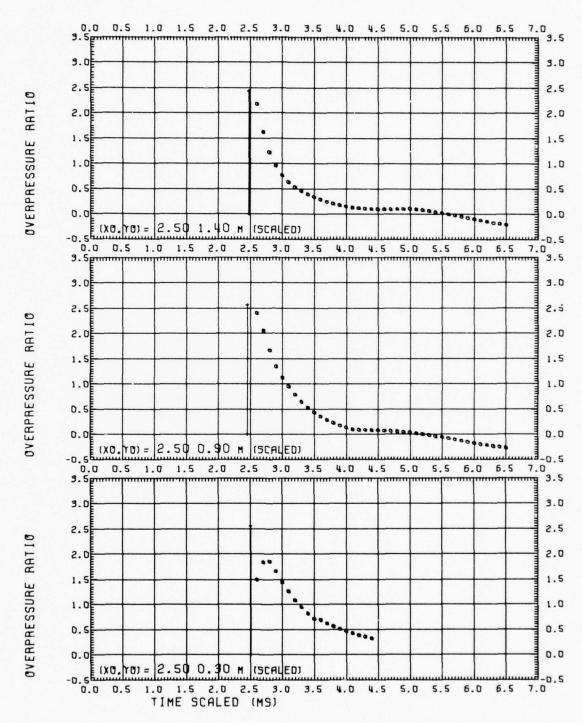


Fig. 25.3 DIPOLE WEST/10 HYDROSTATIC OVERPRESSURE

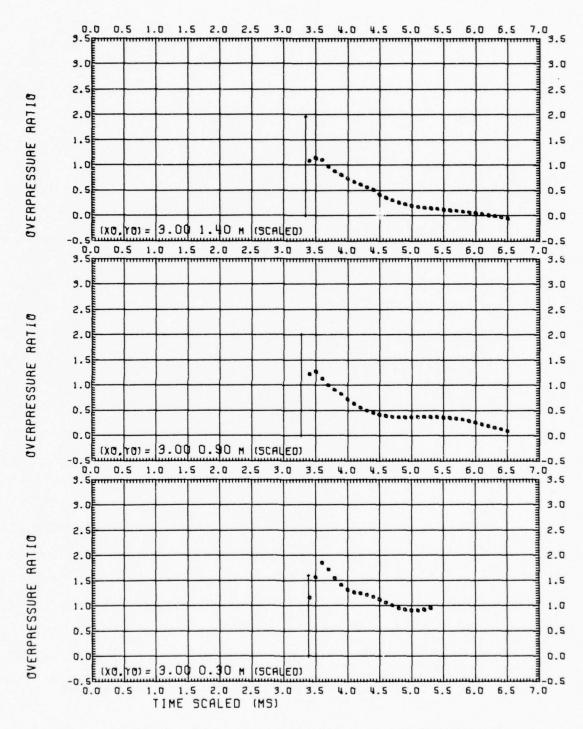


Fig. 25.4 DIPOLE WEST/10 HYDROSTATIC OVERPRESSURE

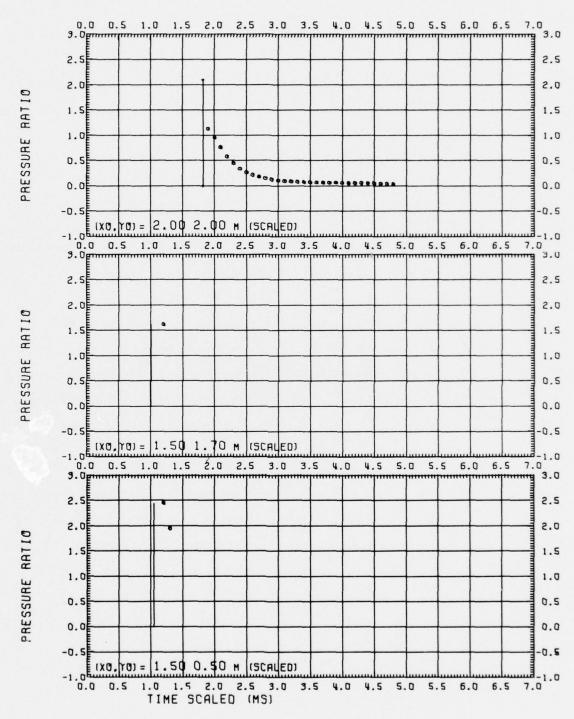


Fig. 26.1 DIPOLE WEST/10 DYNAMIC PRESSURE

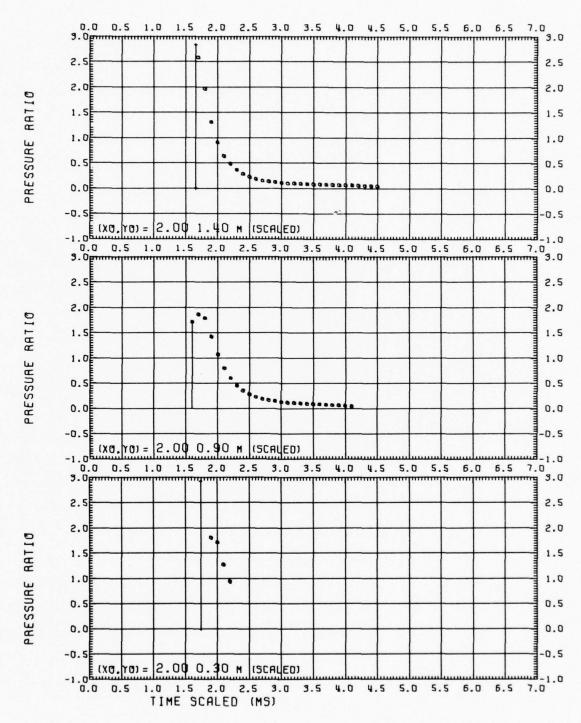


Fig. 26.2 DIPOLE WEST/10 DYNAMIC PRESSURE

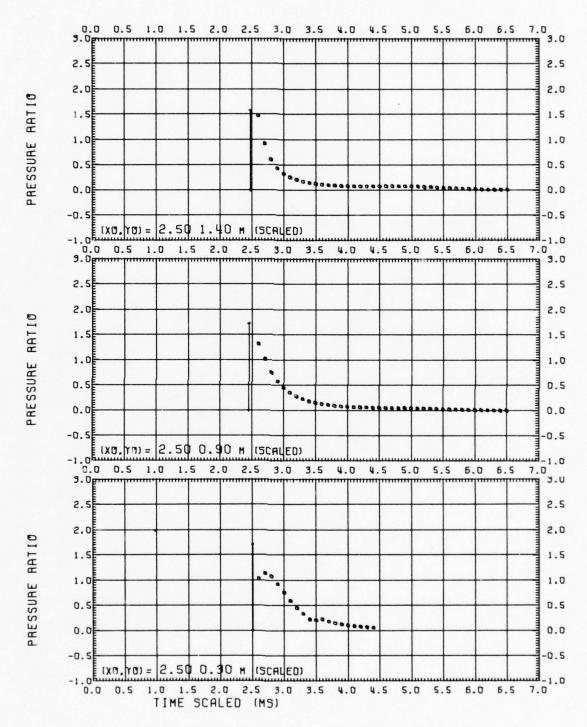


Fig. 26.3 DIPOLE WEST/10 DYNAMIC PRESSURE

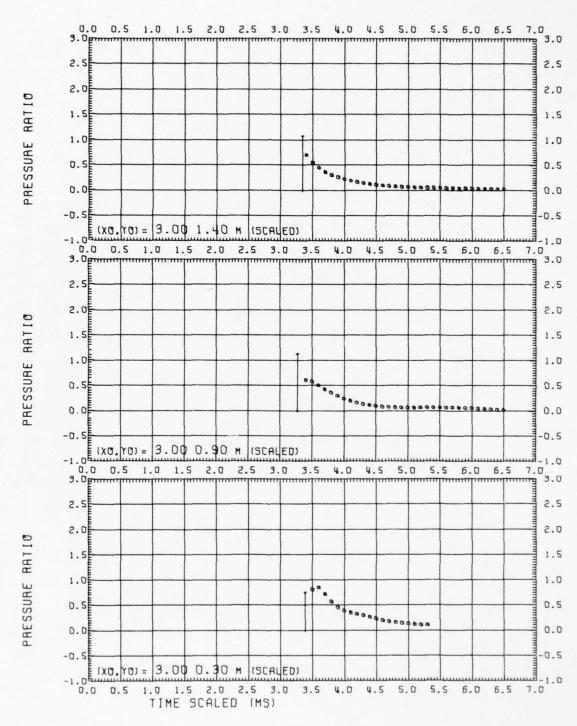
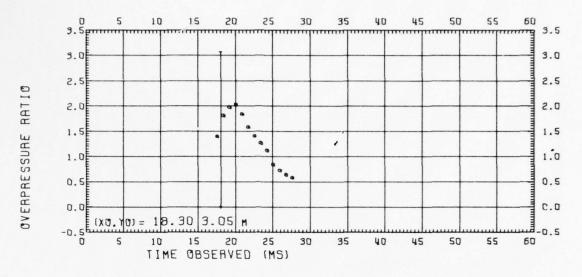
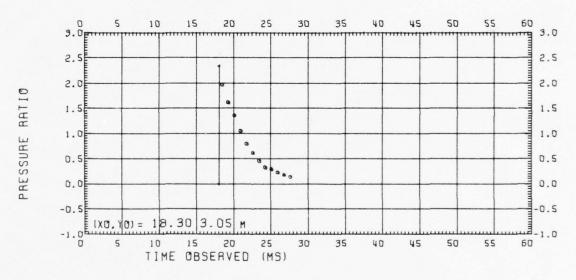


Fig. 26.4 DIPOLE WEST/10 DYNAMIC PRESSURE



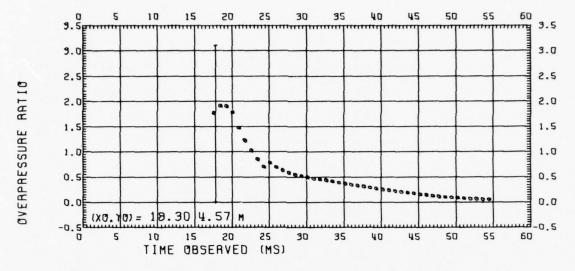
DIPOLE WEST/10 HYDROSTATIC OVERPRESSURE



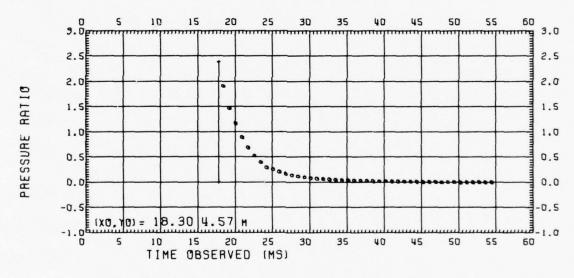
DIPOLE WEST/10 DYNAMIC PRESSURE

Fig. 27.1

PRESSURE RESULTS AT GAUGE POSITION (XO.YO) = 60 FT. 10 FT



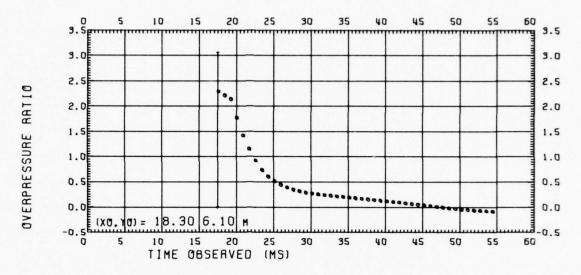
DIPOLE WEST/10 HYDROSTATIC OVERPRESSURE



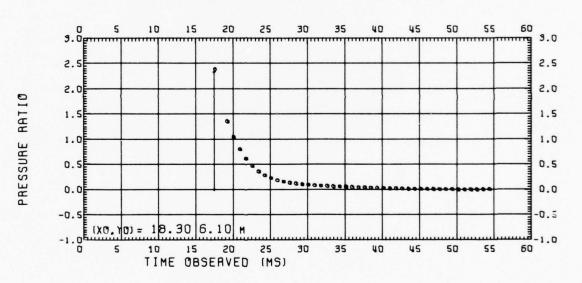
DIPOLE WEST/10 DYNAMIC PRESSURE

PRESSURE RESULTS AT GAUGE POSITION (XO.YO) = 60 FT. 15 FT

Fig. 27.2



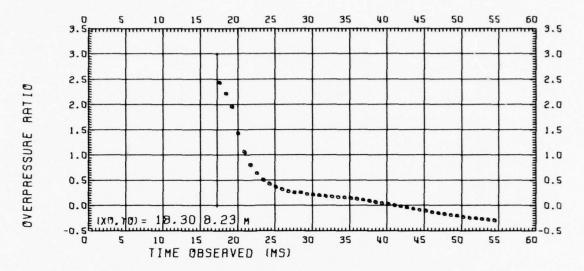
DIPOLE WEST/10 HYDROSTATIC OVERPRESSURE



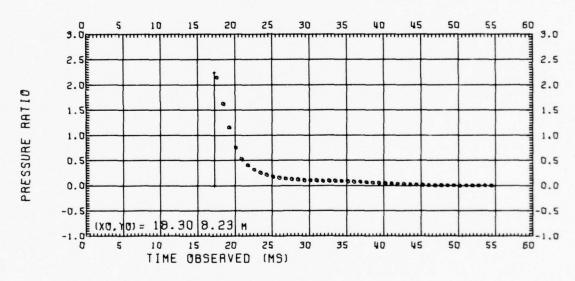
DIPOLE WEST/10 DYNAMIC PRESSURE

Fig. 27.3

PRESSURE RESULTS AT GAUGE POSITION (XO.YO) = 60 FT. 20 FT



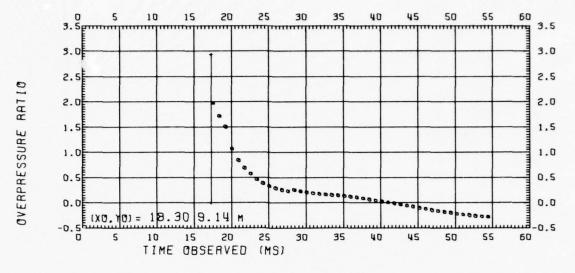
DIPOLE WEST/10 HYDROSTATIC OVERPRESSURE



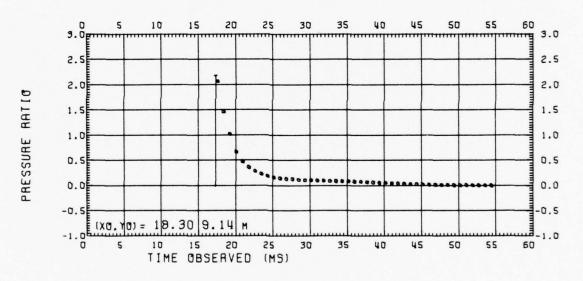
DIPOLE WEST/10 DYNAMIC PRESSURE

Fig. 27.4

PRESSURE RESULTS AT GAUGE POSITION (XO.YO) = 60 FT. 27 FT



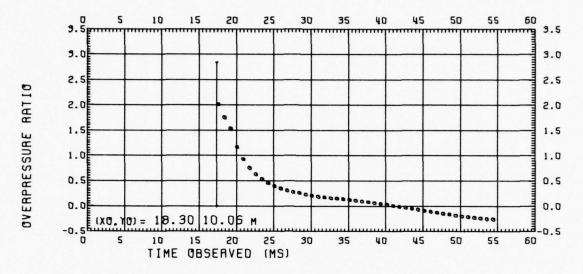
DIPOLE WEST/10 HYDROSTATIC OVERPRESSURE



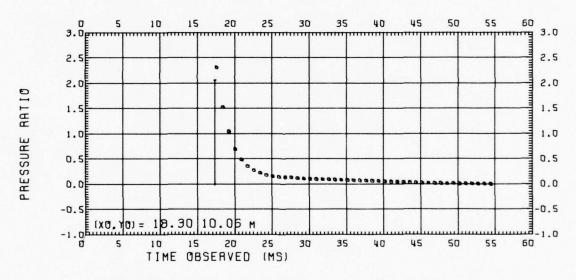
DIPOLE WEST/10 DYNAMIC PRESSURE

Fig. 27.5

PRESSURE RESULTS AT GAUGE POSITION (XO.YO) = 60 FT. 30 FT



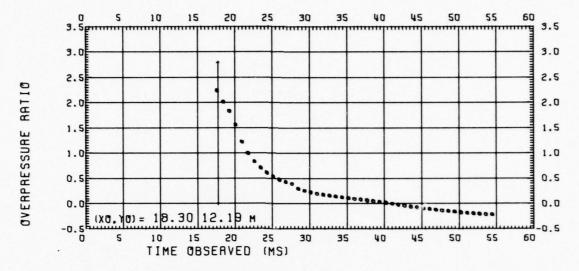
DIPOLE WEST/10 HYDROSTATIC OVERPRESSURE



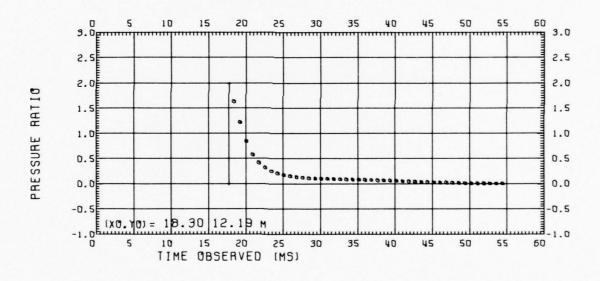
DIPOLE WEST/10 DYNAMIC PRESSURE

PRESSURE RESULTS AT GAUGE POSITION (XO.YO) = 60 FT. 33 FT

Fig. 27.6



DIPOLE WEST/10 HYDROSTATIC OVERPRESSURE



DIPOLE WEST/10 DYNAMIC PRESSURE

Fig. 27.7

PRESSURE RESULTS AT GAUGE POSITION (XO.YO) = 60 FT. 40 FT

TABLE 1

156

TOTAL NUMBER OF POINTS SURVEYED IS 32

SUNDRY DATA LIST

T = 21.3 DEG F. P = 13.7 PSI. SVP 2.95 MM HG. PH = 91.0%; W = 1030.0 LBS
SCALLING TO WO = 2.2 LBS USING FACTORS S = 8.072 AND C = 1.076 FT/MSEC
CALCULATED DISTANCE BETWEEN BOT.CRG. AND G.ZERC.B IS 14.920 FEET TO CALCULATED DISTANCE BETWEEN TOP.CRG. AND BOT.CRG. IS 30.445 FEET CALCULATED DISTANCE BETWEEN GRD.ZERO AND G.ZERO.C IS 1.193 FEET

DIPOLE WEST/10 WF5/235 PHOTOGRANMETRICS CAMEDA (LENS) POSITION IS 2002-619 FEET EAST, 1338-290 FEET NOPTH, AND 2341-720 FEET ELEVATION OPTICAL AXIS IS DRIENTED TO -6.780 DEGPEES EAST OF NORTH AND 0.575 DEGREES ELEVATION (±0.001) CBJECT PLANE IS 606-999 FEET FROM CAMERA, PEFPENDICULAR TO UPTICAL AXIS, AND INCLUDES G.ZERO.C

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CALIBRATION DATA TRANSFORMED TO THE OBJECT PLANE IN FEET

REFERENCE POINT P2 REFERENCE POINT SHIFT -0.583 -1.223 SHIFT X COORD C000PD X

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X—AXIS IS PARALLEL TO HORIZONTAL PLANE AND ORIGIN IS WHERE OPTICAL AXIS INTERSECTS OBJECT PLANE SHIFTS IN CALIGRATION DATA DEFINE THE POSITION OF POINT AS CALCULATED DIRECTLY FROM SURVEY DATA MAXIMUM CALIBRATION ERROR SCALED= 0.030 FEET MAXIMUM CAMERA ORIENTATION ERROR= 0.011 FEET

TOTAL ERRORS IN THE OBJECT PLANE = 0.091 FEET

N SURVEY DATA FUNNING DATA IS TRANSFORMED TO OBJECT PLANE USING REFERENCE POINTS VP 33 AND A CAMERA OPTICS BASED ON CALIBRATION DATA IN THE OBJECT PLANE RATHER

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SMOKE PUFF GRID 1209

DIPOLE WEST/10 WF5/295

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TABLE 4 (continued)

SCALED TIME: C3SERVED TIME MULTIPLIED BY (C/CO)/S, WHERE CO= 340, 292 METERS/SECOND
AND SCALED DISTANCE OBSERVED DISTANCE DIVIDED BY S= CUBE ROOT OF (W/WC)*(FC/P),
WHERE FC= 101.325 KILOPASCALS. (W, WO, AND PARE DEFINED ASDVE.)
SCALED EVENTE STANDARD CHARGE WO IN ATMOSPHERE WHERE CO AND PO ARE AMBIENT (TO= 15 DEGREES CELSIUS).
VELOCITY EXFDESSED IN MACH UNITS IS INVARIANT UNDER SCALING. X IS DISTANCE MEASURED HORIZONTALLY FROW GROUND ZEFO.

Y IS DISTANCE MEASURED HORIZONTALLY FROW GROUND ZERO.

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AND ARE EXPRESSED IN MACH UNITS FELATIVE TO C ABOVE.

CO BASIS OF FIRST SHOCK FROMT PASSING, CODEO USING:

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5 TABLE

SMOKE PUFF GRID 1209 FRIMARY FRONT FROM LOWER	
DIPOLE WEST/10 WF5/295	AMBJENT TEMPERATUPE T= -5.94 DEGPES CELSIUS AMBJENT PRESSURE P= 94.38 KILOPASCALS PREATIVE HUMIDITY PH= 81.0 PER CENT VARCUR PRESSURE VP= 0.32 KILOPASCALS AMBJENT SPESSURE VP= 0.32 KILOPASCALS CHANGE WEIGHT W= 489.9 KILOGAAMS CHANGE WEIGHT H= 4.95 KILOGAAMS SPENARION +2 HS= 4.56 WEIERS SPENARION +2 HS= 4.56 WEIERS SACKIS SCALING FACTOR SF 60.718
SHCCK FRCNT DATA	AMBJENT TRUPERATUP AMBJENT TRUPERSURE VERTITURE HUMBLE AMBJENT FRESSURE CHARGE FEIGHT WE SCHARGE FEIGH

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PRESSURE	865.016	804.5	804.5	752.5	70706	450.6	4 600 6	473.9	473.9	335.5	335.5	335.5
PRESSURE RATIO	9.165	8.524	8.524	7.974	7.497	5.198	D. 198	5.021	5.021	3.555	3.555	3.555
SECOCITY VELOCITY	20.776	2.682	2.882	2.799	2,725	2,336	2.336	2.303	2,303	2.0.2	2.012	2.012
R-SCAL METERS	0.961	866.0	664.00	1.033	1.05P	1 . 318	1. 318	1.347	10347	1.720	1.720	1.720
T-SCAL MSEC	0.425	0.460	09400	0.456	0.531	0.814	0.814	00840	0.849	1.343	1 . 343	1.343
DIFFERENCE	960.0-	-0.093	-0.321	0.363	0 = 321	-0.023	-0-110	00100	-0-043	25000	0.00	-0.029
AETEPS	7.759	8.055	9,055	9.341	8.620	10.640	10.540	10.074	10. A74	13.891	13.881	13.881
R-OBS METERS	7.855	9.148	8.375	7.978	0000	10.664	10.750	101700	1	2 2 2 2	13.866	13.910
T-08S	3.558	3.054	3. 3. 4	4.151	4-447	4.9.4	410	7 113	7	11.006	11. 246	11,246

DENSITY RATIO

T IS TIME-JF-ARRIVAL AND R IS RADIAL PUFF POSITION. PADIUS VALUES ARE FITTED USING RFIT=A+B*T+C*LCG(I+T). SHJCK AND PARTICLE VELOCITIES ARE EXPRESSED IN MACH UNITS. RELATIVE TO THE AMBIENT SOUND SPEED C ABOVE. PRESSURE IS PEAK OVERPRESSURE RATIO (PMAX-P)/P, AND PEAK OVERPRESSURE (PMAX-P) IN KILOPASCALS CBSERVED. WHERE P IS AMBIENT PRESSURE. DENSITY IS EXPRESSED AS A FATIO, RELATIVE TO THE AMBIENT DENSITY D. SCALED TIME= OBSERVED TIME MULTIPLIED BY (C/CG)/S, WHERE CO= 340,292 METERS/SECOND
AND SCALED DISTANCE= OBSERVED DISTANCE DIVIDED BY S= CUBE ROOT OF (W/NG)*(PO/P),
**HERE FC= 101,325 KIOPASCALS* (W, 40, AND P ARE DEFINED ABDVE.)
**ACALED EVENT= STANDARD CHARGE WOIN ATMOSPHERE WHERE CO AND PO ARE AMBIENT (TG= 15 DEGREES VELO) PRESSURE, AND DENSITY, EXPRESSED AS RATIOS, ARE INVARIANT UNDER SCALING.

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			DENSITY	60000	4.049	1000	3000	1000	20.00	001	2000	0.1.0	30100	2000	2000	00000	2.550	2.517	2.203	2.203	2.203	1.986	1.986	1.857
22			PARTICLE VELOCITY	2.426									1.571											169.0
R2 /A77032			PRESSURE	1032,356	1039-366	0000		8024100	795.033	505.602	506.802	505.802	481.141	457.845	296.924	296.924	256.924	287.556	209,358	209.358	209.358	162.216	162.216	136.744
R CHARGE			PRESSURE	950-01	000	00000	001.5	9.135	9.424	5.370	5.370	5.370	5.098	4.851	3.146	3.146	3.146	3.043	2.218	0	0000	1.710	1.710	1.449
ID 1209 FRCN UPFE		A L	SPDCK	100.5	1000	3.521	2007	2.971	2.867	2.367	2.357	2.367	2.317	2.271	1.923	1.923	1.923	1.301	1.703	1-703	1.703	0000		1.000
SMOKE PUFF GR		ES DE ARRIVAL	R-SCAL METERS	000	0.00	0.000	1.073	1.073	1.110	1.336	1.336	1.336	1.365	1.394	1.678	1.678	10678	1.700	1040	0 40	040	0000	20000	20441
n or	0 N N	ECTORY TIME	T-SCAL MSEC		0.473	0.425	0.496	0.466	0.531	0.779	0.770	0.770	414	0 0	1.037	1.037	1.037	0000	1000	7000	0.00	0.000	2.101	20101
0 WF5/295	SCALS NI SCALS SCALS NETERS/SECOND METERS/SECOND	PARTICLE TRAJECTORY TIMES	DIFFERENCE		-0.151	0.044	0.260	70100	-0-0-	8000		0000	1000	000		1000	000	000	0000	01000	100100	0.121	0.510	0.2552
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SHOCK FRONT DAT	AWBIENT TEWEEBATURE AWBIENT TEWEESURE VARCUA PRESSURE VPH AMBIENT TEWESURE AMBIENT TEWESURE CHARGE WEIGHT HE CHARGE PEIGHT HE CHARGE PEIGHT HE SACHS SCALING FACTOR SCALING TO CHARGE WEIGHT HE CHARGE PEIGHT HE C	SHOCK FRENT	1-09S		4 4 6	0 0 0	7 4	101.	0	1000	0.000	6.520	0.130	6.316	7.112	100 361	10,361	10,351	10.656	161011	14,191	14.141	19,010	.01
0) 1																								

| 16.054 | 15.054 | 15.054 | 15.057 | 1.695 | 1.695 | 1.695 | 1.703 | 2.218 | 200.358 | 0.930 | 1.703 | 1.704 | 15.057 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695 | 1.695

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R4 /A770322				PRESSURE KPA	1503.364	763.359	487.257	487.257	349.701	349.701	343.701	251.168	251.168	251.168	193.407	193.467	1930 400	140.00	000000	6200	000000000000000000000000000000000000000	0 10	240.000	000	2000	117.245	117.245	
	TICH PLANE			PRESSURE	15,592	8.083	5.163	5.163	3.705	3,705	3.705	2.661	2.651	2.000	0000	0000	0000	4000	1 704	100	0-4-	1.010	1.4.1	217	0.70	1.242	1.242	
GRID 1209	INTERACTION			SHOCK	3.835	2.816	2.329	2,329	2.043	2.043	2.043	1.6911		11001	1 0000	0 0	0000	2000	1.560	1.860	1.487	10497	1.487	1.447	1 - 4 37	1-437	1.437	
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MOKE PUFF	CH STEM		S OF AR	R-SCAL METERS	66 0	1.39	1.71	1,719	1.98	1.986	000	× 0	2000	000	20.00	20.00	2007	2.79	2.701	2070	3.03	3.03	3,03	3.035	3.214	3.214	3.214	
Σ (1)	4 5	Q Z	JECTORY TIME	T-SCAL MSEC	5	0.835	N.	NI	01	DI	00	00	20	210) U	15	10	0	0	0	m	m	m	m	10	3.684	58	
0 WF5/295		DEGREES CELSIUS CENT CENT CENT COPASCALS COPASCALS ANS ANS 31.0 KILGGPAMS	RTICLE TRA	DIFFERENCE	-0.013	0.149	0000	450.0	10000	2000	20.00	0.263	0.247	-0.309	0.329	0.219	0.165	-0.197	-0.143	10.042	-0.101	0.056	-0.110	0.063	-0.015	0.034	-0.01B	OCC DUTTO IN
DIPOLE WEST/10		32 A D X D X D X D X D X D X D X D X D X D	UTED FROM PA	*ETEPS	8.007	11.234	0,000	10000	10000	100	18.501	18.521	19.521	20.756	20.756	20.756	20.766	22.530	22.530	22.530	26.097	24.497	240497	24.497	25.943	25.943	25.943	A 10 A 0 15 B A 0 1 A
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DIPOLE WEST/10 WF5/295

/A770322

SZZYS WACH STEM ABOVE INTER	CELSIUS	S EPS/SECOND	ILOSRAMS
SHOCK FRONT DATA DIPOLE WEST/10 *F5/295	E T= -5.94 DEG9EES	RELATIVE PUBLITY RHE RI.O PER CENT VARCUT PRESSURE VP 0.32 KILOPASCALS VARCUT PRESSURE OF SOUND CE 328,004 VETERS/SECOND AVRIENT FREED OF SOUND CE 328,004 VETERS/SECOND	CHADGE FEIGHT 3= 4.64 METERS CHADGE FEIGHT 4.64 METERS SEDARATION 42 HS= 4.64 METERS SECHS ARLING FACTOR 3= 4.05 METERS SCALING TG CHARGE WEIGHT WO= 1.00 KILGGRAMS
SHOCK FRENT DATA	AMBLENT TEMPERATURE PROPERTY PROPERTY OF THE P	AFECATIVE HENTOLTY VATOUR PRESSURE VP AMBIENT SEEED OF S	CHANGE REIGHT 38 SEDAGATION #2 HSS SACHS SCALING FACT SCALING TO CHANGE

### COMPOUND ####################################			1	7 101104	RAJECTORY TIME	S OF APRIVA	AL			
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NUMBER

DENSITY RATIO

T IS TIME-CE-ARRIVAL AND B IS RADIAL PUFF BOSITION. RADIUS VALUES ARE FITTED USING RFIT=A+8*T+C*LCG(1+T).

SHOCK AND BARTICLE VELOCITIES ARE EXPRESSED IN MACH UNITS. RELATIVE TO THE AMBIENT SOUND SPEED C ABOVE.

PRESSURE IS AMALENT PRESSURE PATIC (PMAX-P)/P, AND PRESSURE (PMAX-P) IN KILOPASCALS GRSERVED.

WHERE P IS AMALENT PRESSURE DENSITY IS EXPRESSED AS A RATIO, RELATIVE TO THE AMBIENT DENSITY D.

SCALED TIME DUSSERVED TIME WULTIPLIED BY (C/CO)/S, WHERE COB 340,292 WETERS/SECOND

AND SCALED DISTANCE, GRSERVED DISTANCE DIVIDED BY SE CODE FOOT OF (W/WC)*(FO/P).

5 5 TABLE

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MACH STEM AT GROUND SURFACE SMOKE PUFF GRID 1209

WF5/295

WEST/10

DIPOLE

SHOCK FRENT DATA

	NUMBER	5.5	36	4	4	47	69	000	58	72	7.1	82	43	8.3	95	9.6	106	108	107
	DENSITY	m	m	00	in	co	10	2.464	m	0	0	0	On	13	m	PET	10	m	1.735
	PARTI CLE VELCCITY	2.099	1.634	1.336	1.336	1.336	1.109	1.109	1.092	0.930	0.930	0.795	0.785	0.785	0.674	0.674	0.617	0.604	0.604
	PRESSURE KPA	795.298	513.866	367.742	367.742	367.742	273.433	273.433	265.894	209.224	209.224	166,520	163.607	163.607	132.078	132.078	117,153	113.825	113,825
	PRESSURE RATIO	8.426	5.445	3.896	3.890	3.896	2.897	2.897	2.828	2.217	2.217	1.764	1.733	1.733	1 . 399	1.399	1.241	1.206	1.206
٧ ل	SFOCK VELOCITY	. 86	.38	800	. O.B	.09	.86	1.966	· 83	.70	070	.58	.57	.57	648	649	.43	045	1.426
ES OF APRIVAL	R-SCAL METERS	1.384	1.708	1.984	1.984	1.984	2.252	2.252	2,275	2.516	2.516	2.759	2.778	20778	3.022	3.022	3.165	3,200	3.200
CND	T-SCAL MSEC	0.390	1.343	1.695	1.695	1,695	2.081	2.091	2.116	2,501	2.501	2.919	2,954	2.954	3.406	3.406	3.584	3,753	3.753
TEWFERATURE T= -5.04 DEGREES CELSIUS FRESSURE P= 94.34 KILDPASCALS HUMINE B1.0 DES CENT FRESSURE VP= 0.32 KILDPASCALS FRESSURE VP= 0.32 KILDPASCALS SPECEN C= 10.32 KILDPASCALS SPECEN C= 10.32 KILDPASCALS FRESSURE VP= 0.45 KILDRAMS FRESSURE VP= 4.55 WETERS KILDRAMS TO CHARGE WEIGHT WO= 1.0 KILDGRAMS CNT DATA COMPUTED FROM PARTICLE TRAJECTORY	DIFFERENCE	0.000	0.213	660.0	-0.715	-0.057	0.206	0.1AB	-0.119	0.288	-0.023	-0.358	0.222	0.047	0.331	-0.060	-0.445	0.143	0.017
946 948 948 948 948 948 948 948 948 948 948	R-FIT METERS	11.175	13.788	15.014	16.014	16.014	18.178	19.178	13,364	20.308	20.308	22.267	22.424	22.424	24.393	24.393	125	25.830	(2)
AWBIENT TEVERATUDE TE -5.04 RELATIVE HURSTONE DE 94.38 KIL RECTUR PRESSURE DE 94.38 KIL AWATENT SPRES CONNO C 32 KIL CHARGE WEIGHT NE 6000 C 32 KIL CHARGE WEIGHT NO C 32 KIL SCALING TO CHARGE WEIGHT NO T 32 KIL SCALING TO CHARGE WEIGHT NO T 33 KIL SCALING TO CHARGE WEIGHT NO T 34 KIL SC	RACHES	11,124	17,575	15,915	16,730	15.072	17,973	14.091	14,403	20.050	20,331	22,624	22.202	22.377	24,062	240453	26. 493	25,681	25.813
AVBIENT TEVEERATUDE RELATIVE PRESSURE PER VARIOUT PRESSURE VARIENT PRESSUR	1-08s	4.294	11.246	161.61	14,191	14.191	17.423	17.423	17.717	20.940	20.040	24.447	24.719	24.739	28.526	28.526	30.851	31.432	31.432

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GRID 1209 FLOW INTER		R-089	24.75	25.15	25.55	26.32	26.70	27,44	27.30	29.16	2000	20.00	23.5	29.00	30.2	30.0		310	31.8	32.1	32.7	33.0	2 3 5 5	3.20	33.9	34.5	1400	35.1	30.00	35.0	36.3	VOUL OT A
MACH STEM BE		240 -1	28.472	29.310	30,147	31.822	32.659	33.497	35.172	36,009	36.947	386.312	38.521	39.359	41.034	41.871	42.708	44,343	45.221	46.058	47,733	48.570	807.64	100 CO	51.083	51.920	52.535	4 4 3 2	55.270	500 107	57.782	
WF5/295 S CELSIUS LS TERS/SECOND	KILDGRAMS	DAG	> H																													
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7	8 S = 8 B E I C H T B B	T-SCAL	A S S S S S S S S S S S S S S S S S S S	0.597	000000	0.700	400	0.932	006 0	0000	1.195	1.200	1.300	1.400	1.500	1.5664	1.700	1.725	1.000	2.000	2.020	2000	2.300	2.400	2.500	2000	2.700	2.900	2000	3.100	3.200	•
11	TNG FACTO	TICLE VELDO	XETERS OF COMPANY	6.793	8.919	059.5	50803	10.776	11.309	12.065	12.43	13.493	14.159	4 • • • • • • • • • • • • • • • • • • •	150435	15.926	16.631	16,776	17.203	16,301	18.408	16,329	0.00	20.339	20,820	2000	21.753	22.206	22.000	23.514	03,430	
PERSON DESCRIPTION OF THE PROPERTY OF THE PROP	SEDABATION SACHS SCAL SCALING TO	PEAK DART	2 9 9 9 9	0000	5,025	5,862	6.063	10000	7,537	A. 374	212.6	0.00	10.086	11.724	12.561	13,047	13, 399	14.446	15.074	16.911	16,915	17.555	19.423	20.003	20.936	21003	22.410	23.449	24.20	25.960	26.737	27.633

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5.063	5. 310	311	1.228	2000	31.000	26.207	00	4	m
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6.567	10.785	7	1.336	1.469	33.497	26.951	0	(a)	-
7.537	11.315	000	1.402	2.035	34,334	27.316	- 0	m c	~ <
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9*531 24*430 3*413 3*027 0*836 55*945 35*587 6*800 4*458 9*210 24*739 3*500 3*071 0*836 57*782 36*276 6*900 4*494	24.551	25.472	24.176	3.400	3.020	0.356	56.107	35.695	70	4.422	0.248
9*:10 24:735 3*:500 3*:071 0*:836 57*:782 36*:276 6*:900 4*:494	29.210 24.735 3.500 3.071 0.835 57.782 36.276 6.900 4.49 IS FITTED SHOCK RADIUS AT TIME T, CALCULATED FROM PARTICLE TRAJECTORY DATA. ELOCITY IS EXPRESSED IN MACH UNITS RELATIVE TO THE AMBIENT SOUND SPEED C, AND	150.52	24.430	3.413	3.027	0.836	56.945	35.587	80	4.458	0.171
	IS FITTED SHOCK RADIUS AT TIME T, CALCULATED FROM PARTICLE TRAJECTORY DATA. ELOCITY IS EXPRESSED IN MACH UNITS RELATIVE TO THE AMBIENT SOUND SPEEC C, AN	29.310	56. 735	3.500	3.071	0.836	57.782	36.276	06	4.404	0.236
THE PARTY OF THE P	S EXTRABOLATED FROM VELOCITY VERSUS DADIUS DATA AT T. USING VETTEA+B+BADADIE	SEXTR	CLATEDF	NO N	Z Z Z Z	DADIUS DATA A	ANGIENI SCON	SPEEL C	4 6		

TABLE 7.1

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GRID 1209		X-SCAL	METERS	2.005	2014	2007	1	000	200	2000	K 1 0 0	20000	2.225	2.220	2,194	2.161	0.145	2-1-7		1000	0.5.10	C+247	2.2.70	2.300	2.322	2,325	2,310	2.203	010.0		1 1 2 0 2	2.241	
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E WEST/10	ALED TIME		2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2	200	000	0.00	00	0.10	90.0	0.23	0.10	-0.01	0.22	0.16	0-11		000	000	21.0	0000-	-0.03	50.0-	0.23	0.21	a	0		0.00	50.0-	-0015	+0.0-	0.01	A 7 7 8 8 8
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		R-SCAL	METERS	2.093	2.105	2.201	2.191	2.281	2.320	2.299	2.287	2.312	2.335	2.346	2.296	2.254	2.241	2.300	2.317	2.443	2.431	2.423	2.433	2.435	2.428	2.442	2.446	2.355	2.334
		PARTICLE			0.879	0.664	0.642	0.635	0.729	0.693	0.698	0.730	0.848	0.659	0.875	0.508	0.932	0.792	0.890	0.368	0.528	0.903	0.911	646.0	0.919	0.938	1.107	0.922	0.971
/A770322		- 11	MACH NO	-0.12	-0.06	0.24	0.21	0.13	0.19	60.0	0.08	-0.03	-0.10	-0.16	-0.10	-0.11	-0.07	0.26	0.25	0.28	0.22	0.14	0.10	0.01	-0.02	-0.18	-0.10	-0.10	-0.17
		X O =	MACH	0	0.88	0.62	0.61	0.62	0.40	69.0	69.0	0.73	0.84	0.85	0.87	0.80	0.93	0.75	0.85	0.62	06.0	0.89	0.91	0.95	0.92	0.92	1.10	0.92	96.0
		A	U	5	80	CV	02	E 3	58	9 4	25	00	A C	0,609	2 4	73	80	16	00	83	53	4	26	40	67	55	5	26	460 0
FUFF GAID 1209		X-SCAL	ETE	2.062	20103	2.142	2.168	20171	2.255	2.276	2.234	2.308	20312	2.285	2.257	2.237	2.239	2,255	2.299	2,343	2.360	2.404	2.429	2.433	2.414	2, 393	2.402	2.340	2.332
PUFF																													
SMOKE		ш	CODE	~	2	N	2	2	4	4	1	2	2	0	ស	4	4	4	n	3	2	2	Ŋ	u)	S	4	4	4	m
	SE	R-SCAL	METERS	1.719	1.701	10672	1.697	1.692	1.653	1.765	1.786	1.970	1.350	1.834	1.892	2.053	1.857	10974	1.995	10917	2.042	2.030	2.045	2.174	2.156	2.173	2.183	2, 156	2.169
WF5/295	= 2.500	PARTICLE	2	0.536	0,392	0.416	0.398	0.388	00457	0.482	0.538	0.482	0.468	00 4 34	0 3 3 9 0	0.623	0.498	0.549	665 0	0.612	0.574	0.549	0.508	0.575	0.573	0.564	0.720	0.595	0.710
LE WEST/10	CALED TIME	V=DY/DT	ACH	0.24	0.15	0.12	0.10	0.05	0.24	0.12	0.04	0.21	0.14	0.10	90.0	0.03	0.14	0.03	00 00	-0.05	0.22	61 0	0.15	0.12	0.03	-0.01	-0010	0.01	-0.15
DIPOLE	ARTICLE VELOCITIES AT SCALED	U=DX/DT	U	0.49	0.36	0 \$ 0	0.39	0.38	0 4 0	0.47	0.04	44.0	0.45	0.42	0.39	0,63	0.40	0.55	0.60	0.61	0.53	0.51	0 0 0	95.0	0.57	99.0	0.71	0.60	0.69
FIELD	VELOCI	Y-SCAL	E	2.438	2.176	1.911	1.652	1.365	0.572	0.734	0.418	2,334	2. 072	1.921	1.612	645.0	0.858	0.643	0.440	0.225	2.280	2.064	1.0 040	1.502	1.259	1.016	C. 764	0.631	0.444
VELOCITY FIEL	PARTICLE	X-SCAL	w	1.559	1.636	1,560	1.696	10675	1.685	1.719	1.780	1.763	1.815	1.830	1. F 3.2	2.005	1.873	1.912	1.945	1.904	1.561	2.000	2.041	20144	2.152	2.175	2.151	2.035	2.123

DBSERVED DISTANCE VALUES= 8.0718 TIMES SCALED VALUES AND CBSERVED TIME VALUE = 8.374? TIMES SCALED VALUE. VELDCITY VALUES AS SHOWN ARE INVARIANT UNDER SCALING.

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TABLE

/A770322

WF5/295 000

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VALUES VALUE. 8.3742 TIMES SCALED 8.3742 TIMES SCALED ARE INVARIANT UNDER SC OBSERVED DISTANCE VALUES= AND CBSERVED TIME VALUE = VELOCITY VALUES AS SHOWN A

VELOCITY FIEL	DO OIPOL	E WEST/1	0 WF5/295		SMOKE PUFF	GRID 1209			/A77032	8			
PARTICLE VELD	CITIES AT	SCALED TIME	E= 3.500	S									
Y-3CA	L U=DX/	V=DY/	PARTICLE	R-SCAL	REGN	X-SCAL	Y-SCAL	FOXXO=U	0 % N	PART ICLE	MF TERS	CODE	
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OHAGEDVED DIS	TANCE VALUE	S= 8,071	6 TIMES SCA	ALED VALU	ES								

OBSERVED DISTANCE VALUES= 8.0718 TIMES SCALED VALUES AND DBSERVED TIME VALUE = 8.3742 TIMES SCALED VALUE. VELOCITY VALUES AS SHOWN APE INVARIANT UNDER SCALING.

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	12		•		2.221	2	3.051	9.	10			d	
200			•		2.346	2	3.034	· t	10	0.0		0	
0			•		2.390	ω·	3.031	8	10	0		04	
200	175				204.00	4 .	3.014	0	.0	0		01	
20	200				224.00	4	3,058	0	0	n.		3	
333	46				2. 375	\$ 17	2005	•	0	-		22	
.30	. 24				2.317		32145	•	17.56	-4 (N C	
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· FB	44				2.702) v	2000	10	0.11	30	0,	71	
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WEST/10 TIME ED.

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 $\overset{\circ}{\times} \overset{\circ}{\circ}$

B.071P TIMES SCALED VALUES B.3742 TIMES SCALED VALUE. ARE INVARIANT UNDER SCALING. CBSERVED DISTANCE VALUES= AND GBSERVED TIME VALUE = VELOCITY VALUES AS SHOWN A

183

	60,000								
0 .	STO / G18		1 V O F O	0120					
IME	2,0000	M.S							
TOXY	PASTICLE VELOCITY	R-SCA	REGN	X-SCAL METERS	Y-SCAL METEGS	VACH NO	WACH NO	VELOCITY	METERS
	0.1	1.0		(L)	1.53	0.0	0.11	0.0	
	0.173	1.94		UU	1.57	000	4 O	0.0	3.0
	000	2000		D. C.		0.5	0	0.28	2.9
	0.125	000		10	1.09	0.0	0	0.27	200
	00 157	1.33		6.	0.67	2.00	0	000	3.
	0.139	2.0		a r	0.00	0.00	20	20.00	
	0.212	2.14		- 4	000	000	$\supset C$	00	200
	0.238	2.12		UU	0.0		20	0.4	3.0
	0,237	2.10			200	000	2	0.32	30
	0.245	201		4 .7	000	000	. 0	0.27	3.5
	0.302	**		-	0 0	0	-	0.29	3.1
	00.214	200		4.1	1,66		.0	0.29	3.1
	0.1.0	000		111	10	0	0	0.3	3.
	0.120	7.00		1.65	2	0	0	0 30	3
	1000				0	0	0	0.28	3.
	0.00	0.0			0	0	0	0.30	3
	£12.0	0.00			0.0	0	0	0.36	3.
	0000	0		0	0.48	0.0	~	0.40	10
	0.293	2.5		6	0.24	0	0,	0.3	2.
	0.242	2.5		~	0	0	-	30	
	0.190	204		٠.	200		-	200	
	0.124	2.4						0	3
	00101	203					-	0.3	3
	00.204	0.0			100	0		0.3	3
	0000				2	0	-	0.3	3
	000				0	0	-	0 3	3.
	0000	000			0	0		0.3	3.
	00000	0 0		-	0.6	0		4.0	3
	010.0	10			0.4	0	-	4.0	3
	0.071	200			002	0	-	0 . 3	0
	1 70 0	2.7			0.0	• 0		0.2	2
	0.100	2		-	200	• 0		0.3	2
	0.236	2.7		~	2.0	•		0.3	20
	0.264	0.00		-	1.8	•	_	000	0,1
	0.205	2.0		-	1.6	•		0.3	0,1
	0.317	203		-	1.4	0		0.0	0,1
	0.299	200			1.2	0		4.00	,,
	0,299	200		-	1.0	•		0.0	2.
	0.304	2		+	C.B	0		4.0	
	0.276	0		-	0.6	0		4.0	
	0.378	0.0			0.4	0		4.0	2.
	0.223	3.0		-	0.2	•0		4.0	3
	0.272	3.0							

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00000000000000000000000000000000000000	1746 1746 1717 1740	3.00 3.00 3.00 3.00 3.00 3.00 3.00 3.00	- m 10 0 a		-0	22	
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TABLE 7.15

		REGN	0	מי	0	S	S	S	4	4	4	4	m	S	S	5	ď) u	o cr) 4	td	1 1	4	· u	n un	u u	ď) li) (C	4	4	4	4	3	2	S	S	S	S	S	4	4	4	4	m	
		U	METERS	-		-	-	0	0	-	-	OJ.	-	M	CV	2	0	10	10	10	10	M	J. Le	.4	4	· ct	-4	m	1.4	ct	3	ct	10	ct	in	10	15	10	10	in	10	10	1.0	10	0	
21		ICL	VELOCITY	000	.27	. 26	.26	. 24	3	. 28	. 22	. 27	. 25	. 23	027	027	25	2	2	10	100	20	20	2	25	.26	2	0	5	30	200	52.	.28	. 27	. 25	620	.30	0 CB	.25	. 25	.27	32	32	32	39	
/A770322		V=DY/DT	I C	0.13	0.12	0.07	90.0	90.0	0.13	0.08	0.02	0.07	-0.01	0.10	0014	0.10	60.0	0.03	0.11	0	0.10	0.09	0000	0.10	0.11	0.10	0.10	0.05	0.05	60.0	0.10	0.08	90.0	0.05	60.0	0.11	0.11	0.10	90.0	0.04	0.04	0.07	0.07	0.05	0.05	
		0	ACI		2	0				2	0	20			2.	2	2	2	0	0	N	2	2	2	2	2.	-	-	2	2	2.	2	2.	0 5	20	0	20	0		. 2	0	(1)		5	3	
		001	D . T	0	5			•	-	0	01	00			0	5	1.	4)	3	0		1	5	3	~	0	0	4	7	-	0	1.	4	2	2	0	0)	0	4	2	0	0	90	4	.3	
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0 WF5/295	E= 6,000	PARTICLE	0.0	0		0	•		0.	•			•	•			0		-	0,2	20		-		0.				0		0		·				0	•	0			2.	0	2		
E WEST/1	ALED TIM	V=DY/DT	0.0	0	0.	- (20	2 0	0 0	0	20	20) (0	0	0	0	~	0	-	0	0	0	0	0	0	0	~	0		Q,	-	00	20	0	00	3	0.	4	0	0	-	0	0	0	0
DIPOL	TIES AT SC	MACH NOT	0	0	06	20	20	36	> -	. 0	0	>-	. 5	20	20	20	2	0	-	-	N.	~	-	0	0	-	~	~	V	NI	VI	v.	40	20	v.	٠,		40	46	v.	-	NI	VI	NI	N	- 2011 1417 3
FIELD	VELOCI	Y-SCAL MFTFFS	2.25	0	000			13	117		14	40	00	144				. 42		5.			07.	0	. 47						10.	900	0 0	10	30			76			. (200	4 0	. 50	CTATATO
VELOCITY	PARTICLE	X-SCAL FFFFS	1.932	611	7. 16	3 16	a		. 1	10	16		0	5 0				00	110	350		2 1	010	36	03	110	3 1	47 1	40	tir	30	1 6	* a	160	1.0	30	- 2	00	3.0	0 -	- 1	n.	-1	20	**	00000000

OBSERVED DISTANCE VALUES= 8.0718 TIMES SCALED VALUES AND CBSERVED TIME VALUE = 8.3742 TIMES SCALED VALUE. VELOCITY VALUES AS SHOWN ARE INVARIANT UNDER SCALING.

VELOCITY FIELD

PARTICLE VELOCITIES AT SCALED TIME= 6.500 MS

DIPCLE WEST/10 WF5/295

/A770322

SMOKE PUFF GRID 1209

SCODE																																												
Z H V C A	3.15	0	1	4	u (10	I Le	4	7	2	2	1	100	1 14	3	0	10	UC	0 0	3	37	6	M	0	M	2 4	-	. 1	1	- 0	10	4	1	1	20	14	20	0 0	200	- 0	4 4	1 4	0	0
VELOCITY	0.00	1 5	. 21	25.	24	26	200	. 21	1 7	000	221	100	020	0	25	000	1	000		. 23	. 22	.17	. 24	224	503	1	000	10	10	10	10	200	.07	533	120	2	10	000	10	10	10	10		. 24
MACH NO	0	0.07	0.07	90.0	0.07	60.0	0.08	0.03	-0.03	0.01	-0.01	0.13	0.11	0.06	90.0	0.01		000		0.00	90.0	0.05	60.0	0.10	0.06	0.01	50.0	000	000	0.06	0.07	0.03	00.00	0.10	60.0	6000	40.0	0.04	0.07	2000	000	50.0	30	0.03
MACH NO	0.0	-	O	0		10	N	C	-	0	CV	0	CV	0	N	0	10	0 0			0	-	2	2	C		0	0	0	a	0	~	0	~	~	2	-		10	10	C	10	9 (4
NE TEACAL	37	22	17	16	5.5	3.6	110	0 01	11	111	20	34		0	73		12.	10		1	.73	130	31	.12	E C	00	0 4		7	0.3	71	50	2 8	.2€	0.07	050	0	47	20	0	0	1 4	0 0	000
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NACH	0.11	0000	0.01	0.08	0.03	0.04	0.08	0.02	0.12	0.02	0.04	0.07	0.05	-0.01	0.03	0.03	-0.01	0-14		0.00	10.01	0.01	0.02	0.08	0.01	-0000	0.03	0.03	0.02	40.0	0,01	-0.00	-0.02	0.03	0.03	-0.02	0.05	0.06	0.07	0.05	0.02	-0.03	000	20.0
MACH	0	0	0	0	0	0	0.0	0	0	0	0	000	0	0.0	0.0	0	0	0) (0 0	•	C	0.0	0	0	-	-	C	-	0	-		0	S	-	-	N	-	-	-	-	-		-
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TERS	5	00	0	0	40	15	0 1	(1)	16	10	0	50	T	0	17	33	4 4	Cu	L	2 4	0 0	9	4.1	43	23	56	60	72	73	10	70	0	73	66	72	10	06	25	63	96	40	5	10	5

TABLE 8.1

		REGN CODE		CODE 1		SEGN CODE		CODE SPEC SPEC	v	80 80 80 80 80 80 80 80 80 80 80 80 80 8	
		METERS 1.163		METERS 1.200 1.232		METS A P P P P P P P P P P P P P P P P P P		A	7	A A A A A A A A A A A A A A A A A A A	
		DENSITY RATIO 1.0044		DENSITY RATIO 1.2296 1.261		DENSITY RATIO 1.6643 1.6621 1.3955		DENS RATIC 1100000000000000000000000000000000000	5	D	
70322		Y-SCAL METERS 0 • 2394 0 • 44		X X X X X X X X X X X X X X X X X X X		ME-SCAL 00.1892 00.1883		MY)	A 4 4 4 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5	
/A7		X-SCAL METERS 1.151		XX METSCAL 11.1899 11.1883 17.50		XX SCAL 1 1 2 2 3 0 1 4 2 1 0 0 1 4 2 1 0		AET SC T SC		METTER 11.1.1.1.1.1.1.1.1.1.1.1.1.1.1.1.1.1.1	SMOKE DIEES
		я В В В В В В В В В В В В В В В В В В В		ж о о о о о о о о о о о о о о о о о о о		გე მე გე გეგ გ. 4 - 4		C C C C C C C C C C C C C C C C C C C		80 80 80 80 80 80 80 80 80 80 80 80 80 8	BULKING
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PUFF GRID		DENSITY RATIO 10007 10037 10111		DENSITY 2ATIO 10.330 10.191 10.327		DE VSTTY R VSTTY 100000000000000000000000000000000000		DE 2000 S S S S S S S S S S S S S S S S S		EN 2000 EN 200	A CELL DF
SMCKE		Y-SCAL METERS 1.328 0.954 0.778		Y-SCAL METERS 1.325 1.143 0.950		Y-SCAL 10.326 10.326 00.968		METERS 1010176 000176 000176 000176		MY WE	FILLI IS
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DESCRIPTION OF VALUES.

AND OBSERVED TIME VALUES.

DENSITY VALUES AS SHOWN ARE INVARIANT UNDER SCALING.

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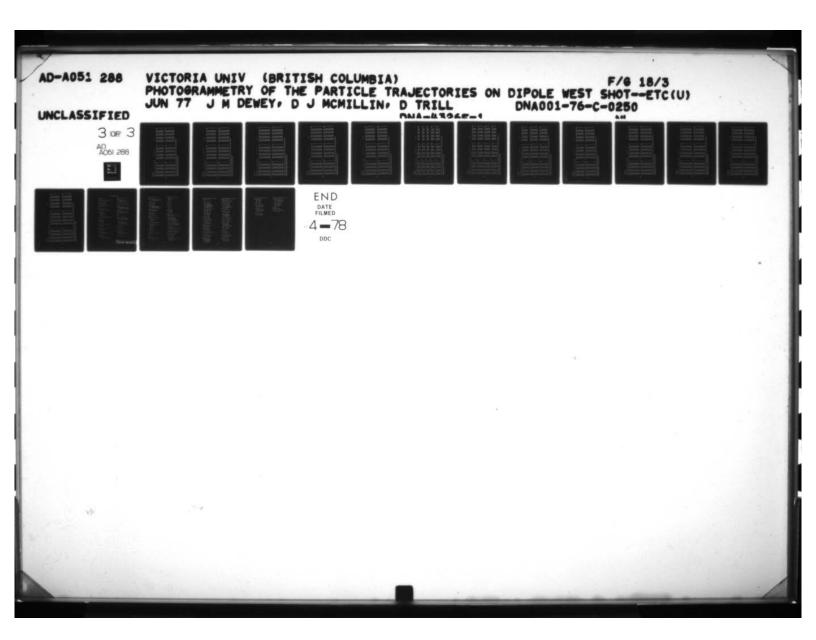
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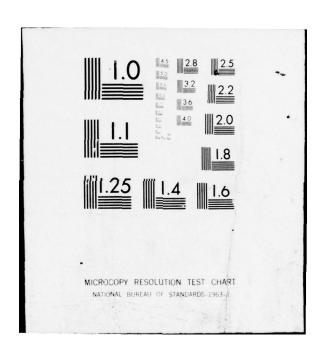
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DENSITY IS AVERAGED OVER THE AREA OF THE CELL AND IS EXPRESSED AS A RATIO TO THE AMBIENT DENSITY.

DHSEPVED DISTANCE VALUES = 8.3742 TIMES SCALED VALUES

DENSITY VALUES AS SHOWN ARE INVARIANT UNDER SCALING.

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X AND Y LOCATE THE CENTER OF A PLANE OUADRILATERAL WHICH IS A CELL OF 4 NEIGHBOURING SMOKE PUFFS.

DENSITY IS AVERAGED OVER THE AREA OF THE CELL AND IS FXPRESSED AS A RATIO TO THE AMBIENT DENSITY.

THE SERVED DISTANCE VALUES B.0719 TIMES SCALED VALUE.

AND CRSEEVEL BY VALUES B.3742 TIMES SCALED VALUE.

DENSITY VALUES AS SHOWN ARE INVARIANT UNDER SCALING.

TABLE 8.8

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X AND Y LOCATE THE CENTER OF A PLANE CUADRILATERAL WHICH IS A CELL OF 4 NEIGHBOURING SWOKE PUFFS. Density is averaged over the area of the cell and is expressed as a ratio to the avutent density.

DBSEDVED DISTANCE VALUES= 8.0718 TIMES SCALED VALUES
AND CHSERVED TIME VALUE = 8.3742 TIMES SCALED VALUE.
DENSITY VALUES AS SHOWN ARE INVARIANT UNDER SCALING.

TABLE 9.1

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A770322	Y-SCAL METERS 0.394 0.204	Y-SCAL METERS 0.591 0.1988	METERS 0 • • 0 9 9 2 0 0 0 1 8 6 6 1 8 6 6 1 8 6 6 1 8 6 8 8 8 8 8	7-SCAL 0-178 0-178 1-868 1-868	METS 10 TS 10
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X AND Y LOCATE THE CENTER OF A PLANE CUADPILATERAL WHICH IS A CELL OF 4 NEIGHBOURING SMOKE DUFFS.

OVERPRESSOUR! IS AVERACED OVER THE AREA OF THE CELL AND IS EXPRESSED AS A RATIC TO THE AMBIENT PRESSURE.
OBSERVED DISTANCE VALUES= 6.0716 TIMES SCALED VALUES
AND CREEVED DIME = A.3742 TIMES SCALED VALUE.
PRESSURE VALUES AS SHOWN ARE INVARIANT UNDER SCALING.

TABLE 9.2

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A AND Y LICEATE THE CENTER OF A PLANE GUADDILATERAL WHICH IS A CELL OF 4 NEIGHBOURING SMOKE PUFFS.

OVERPRESSURE IS AVERAGED OVER THE AFEA OF THE CELL AND IS EXPRESSED AS A RATIO TO THE AMBIENT PRESSURE.

AND OBSERVED TIME VALUES = 8.374.2 TIMES SCALED VALUES

FRESSUGE VALUES AS SHOWN ARE INVARIANT UNDER SCALING.

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X AND Y LOCATE THE CENTER OF A PLANE QUADRILATERAL WHICH IS A CELL OF 4 NEIGHBOURING SMOKE PUFFS.

OVEFPRESSURE IS AVERAGED OVER THE AREA OF THE CELL AND IS EXPRESSED AS A RATIO TO THE AMBIENT PRESSURE.

OBSERVED DISTANCE VALUES= 8.0718 TIMES SCALED VALUES
AND CHSERVEC TIME VALUE = 6.3742 TIMES SCALED VALUE.
PRESSURE VALUES AS SHOWN ARE INVARIANT UNDER SCALING.

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X AND Y LOCATE THE CENTER CF A PLANE GUADRILATERAL WHICH IS A CELL OF 4 NEIGHBOURING SMOKE PUFFS.

OVERPRESSURE IS AVERAGED OVER THE APEA OF THE CELL AND IS EXPRESSED AS A RATIO TO THE AMBIENT PRESSURE.

CREENUE LISTANCE VALUES: 8.0718 TIMES SCALED VALUES

AND OBSERVED TIME VALUE: 8.3742 TIMES SCALED VALUE.

PRESSURE VALUES AS SHOWN ARE INVARIANT UNDER SCALING.

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OBSERVED DISTANCE VALUES: N. 9.0718 TIMES SCALED VALUES

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PRESSURE VALUES AS SHOWN ARE INVARIANT UNDER SCALING.

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OVERPRESSURF IS AVERAGED DV6A THE AREA OF THE CELL AND IS EXPRESSED AS A RATIO TO THE AMBIENT PRESSURE.

OPSERVED DISTANCE VALUES = 8.071A TIMES SCALED VALUES

AND CBSEAVED TIME VALUE = 3.3742 TIMES SCALED VALUE.

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